

- introduction
- particle physics
- synthesis of elements
- big bang nucleosynthesis
- observations

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where do all the elements come from?

hydrogen 1 H 1.0079																	helium 2 He 4.0026				
lithium 3 Li 6.941	beryllium 4 Be 9.0122															boron 5 B 10.811	carbon 6 C 12.011	nitrogen 7 N 14.007	oxygen 8 O 15.999	fluorine 9 F 18.998	neon 10 Ne 20.180
sodium 11 Na 22.990	magnesium 12 Mg 24.305															aluminium 13 Al 26.982	silicon 14 Si 28.086	phosphorus 15 P 30.974	sulfur 16 S 32.065	chlorine 17 Cl 35.453	argon 18 Ar 39.948
potassium 19 K 39.098	calcium 20 Ca 40.078	scandium 21 Sc 44.956	titanium 22 Ti 47.867	vanadium 23 V 50.942	chromium 24 Cr 51.996	manganese 25 Mn 54.938	iron 26 Fe 55.845	cobalt 27 Co 58.933	nickel 28 Ni 58.693	copper 29 Cu 63.546	zinc 30 Zn 65.39	gallium 31 Ga 69.723	germanium 32 Ge 72.61	arsenic 33 As 74.922	selenium 34 Se 78.96	bromine 35 Br 79.904	krypton 36 Kr 83.80				
rubidium 37 Rb 85.468	strontium 38 Sr 87.62	yttrium 39 Y 88.906	zirconium 40 Zr 91.224	niobium 41 Nb 92.906	molybdenum 42 Mo 95.94	technetium 43 Tc [98]	ruthenium 44 Ru 101.07	rhodium 45 Rh 102.91	palladium 46 Pd 106.42	silver 47 Ag 107.87	cadmium 48 Cd 112.41	indium 49 In 114.82	tin 50 Sn 118.71	antimony 51 Sb 121.76	tellurium 52 Te 127.60	iodine 53 I 126.90	xenon 54 Xe 131.29				
caesium 55 Cs 132.91	barium 56 Ba 137.33	57-70 *	lutetium 71 Lu 174.97	hafnium 72 Hf 178.49	tantalum 73 Ta 180.95	tungsten 74 W 183.84	rhenium 75 Re 186.21	osmium 76 Os 190.23	iridium 77 Ir 192.22	platinum 78 Pt 195.08	gold 79 Au 196.97	mercury 80 Hg 200.59	thallium 81 Tl 204.38	lead 82 Pb 207.2	bismuth 83 Bi 208.98	polonium 84 Po [209]	astatine 85 At [210]	radon 86 Rn [222]			
francium 87 Fr [223]	radium 88 Ra [226]	89-102 **	lawrencium 103 Lr [262]	rutherfordium 104 Rf [261]	dubnium 105 Db [262]	seaborgium 106 Sg [266]	bohrium 107 Bh [264]	hassium 108 Hs [269]	meitnerium 109 Mt [268]	ununilium 110 Uun [271]	unununium 111 Uuu [272]	ununbium 112 Uub [277]	ununquadium 114 Uuq [289]								

* Lanthanide series

** Actinide series

lanthanum 57 La 138.91	cerium 58 Ce 140.12	praseodymium 59 Pr 140.91	neodymium 60 Nd 144.24	promethium 61 Pm [145]	samarium 62 Sm 150.36	europium 63 Eu 151.96	gadolinium 64 Gd 157.25	terbium 65 Tb 158.93	dysprosium 66 Dy 162.50	holmium 67 Ho 164.93	erbium 68 Er 167.26	thulium 69 Tm 168.93	ytterbium 70 Yb 173.04
actinium 89 Ac [227]	thorium 90 Th 232.04	protactinium 91 Pa 231.04	uranium 92 U 238.03	neptunium 93 Np [237]	plutonium 94 Pu [244]	americium 95 Am [243]	curium 96 Cm [247]	berkelium 97 Bk [247]	californium 98 Cf [251]	einsteinium 99 Es [252]	fermium 100 Fm [257]	mendelevium 101 Md [258]	nobelium 102 No [259]

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stellar fusion reactions!
(stellar evolution work in 1920/30's...)

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- cannot account for observed ${}^4\text{He}$ abundance

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→ if MW were to “shine” for 10^{10} years, it would generate*

4% ${}^4\text{He}$

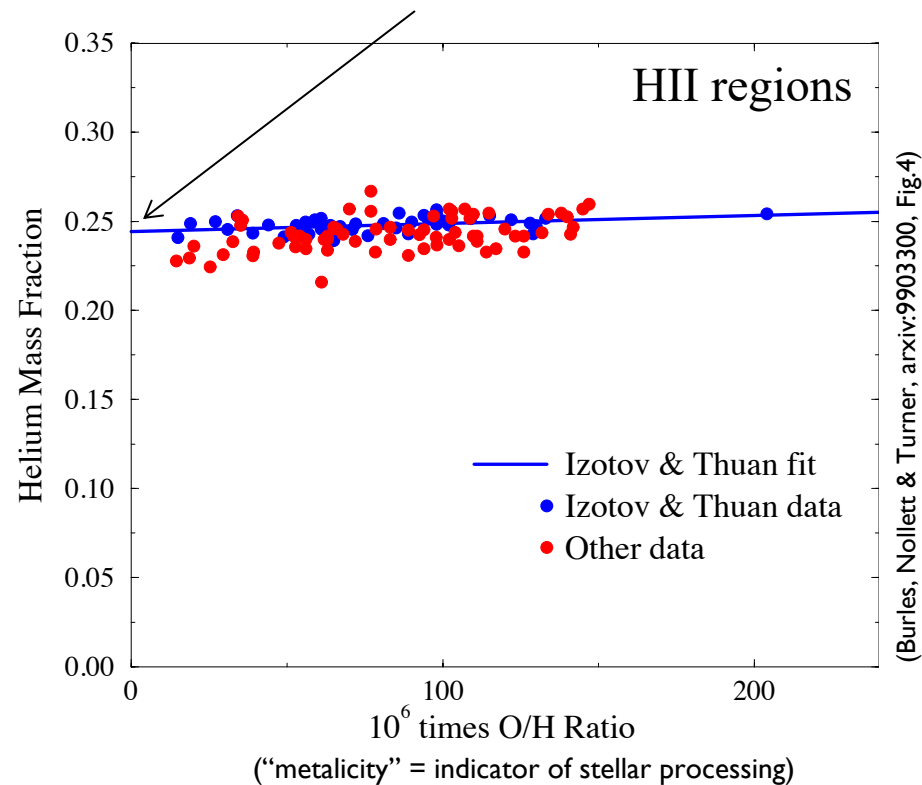
*(Burbidge, 1958, PASP, 70, 83)

▪ stellar fusion...

- cannot account for observed ${}^4\text{He}$ abundance

→ if MW were to “shine” for 10^{10} years, it would generate

4% ${}^4\text{He}$ vs. 24% ${}^4\text{He}$ as observed!



▪ stellar fusion...

- cannot account for observed ${}^4\text{He}$ abundance
- cannot create D^* but only burn it:

fusion reactions in the sun are consuming D faster than generating it

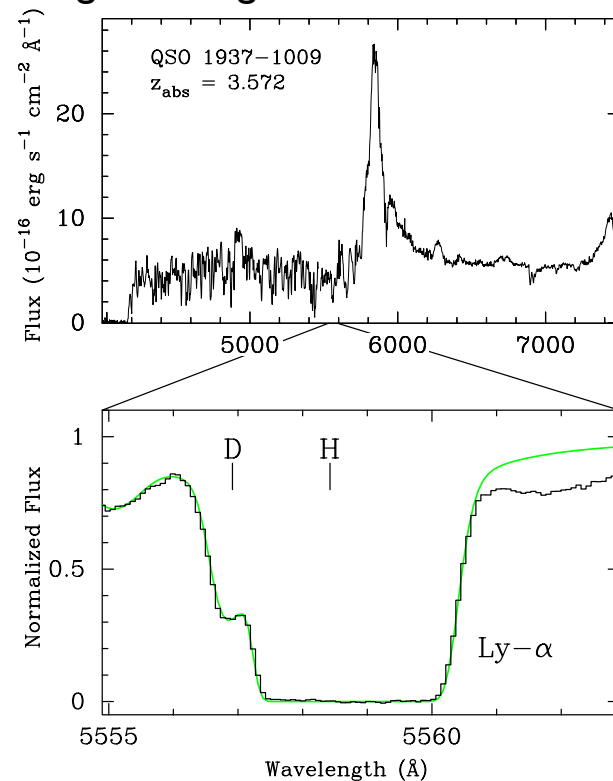
* $\text{D} = {}^2\text{H} = \text{p} + \text{n}$

▪ stellar fusion...

- cannot account for observed ${}^4\text{He}$ abundance
- cannot create D but only burn it:

fusion reactions in the sun are consuming D faster than generating it

...but we do observe D
in the interstellar medium:



▪ stellar fusion...

- cannot account for observed ${}^4\text{He}$ abundance
- cannot create D but only burn it

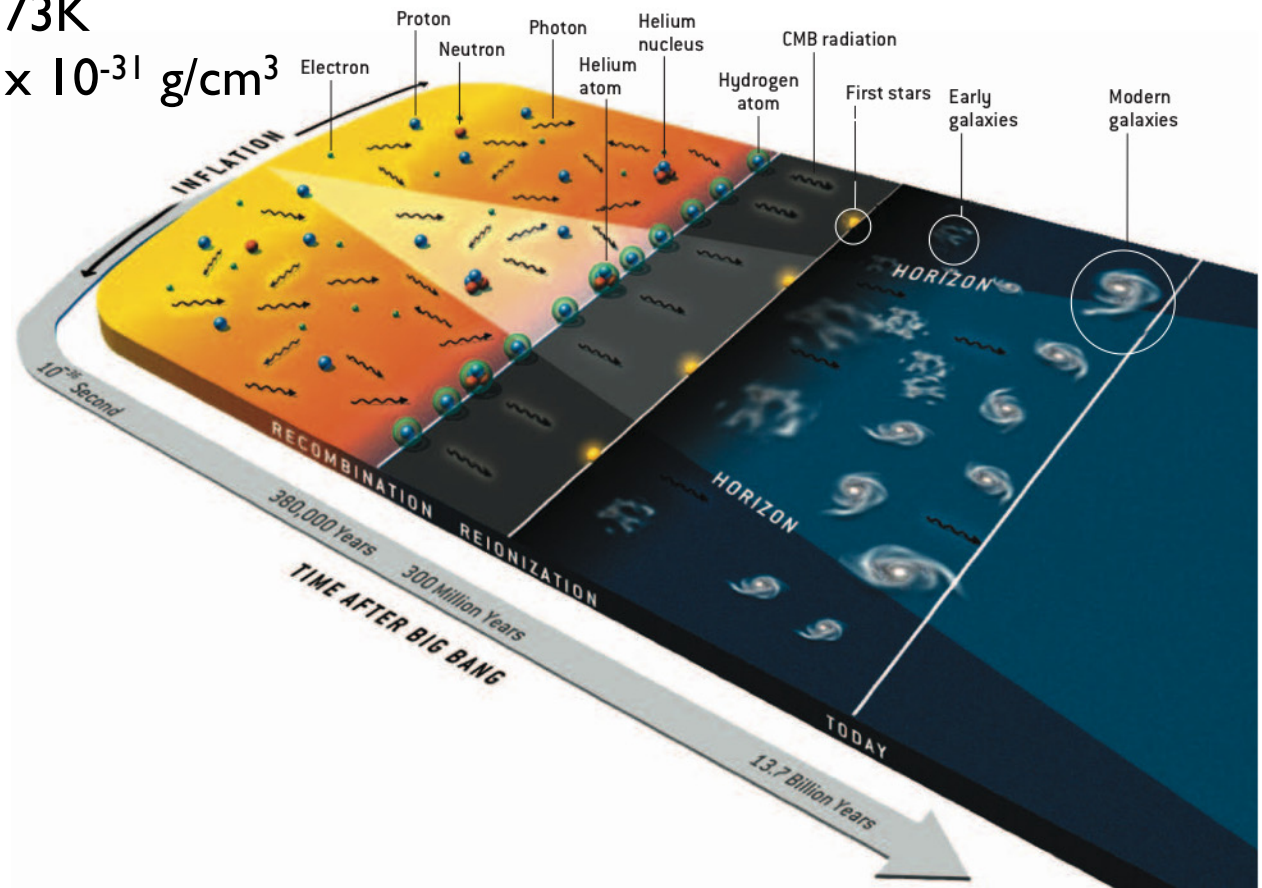
fusion events other than stellar must have happened during the ca. 13.7 Gyrs of the Universe's existence...

■ cosmological considerations

- $z = 0$ (today):

$$T_{CMB} = 2.73K$$

$$\rho_b = 5 \times 10^{-31} \text{ g/cm}^3$$

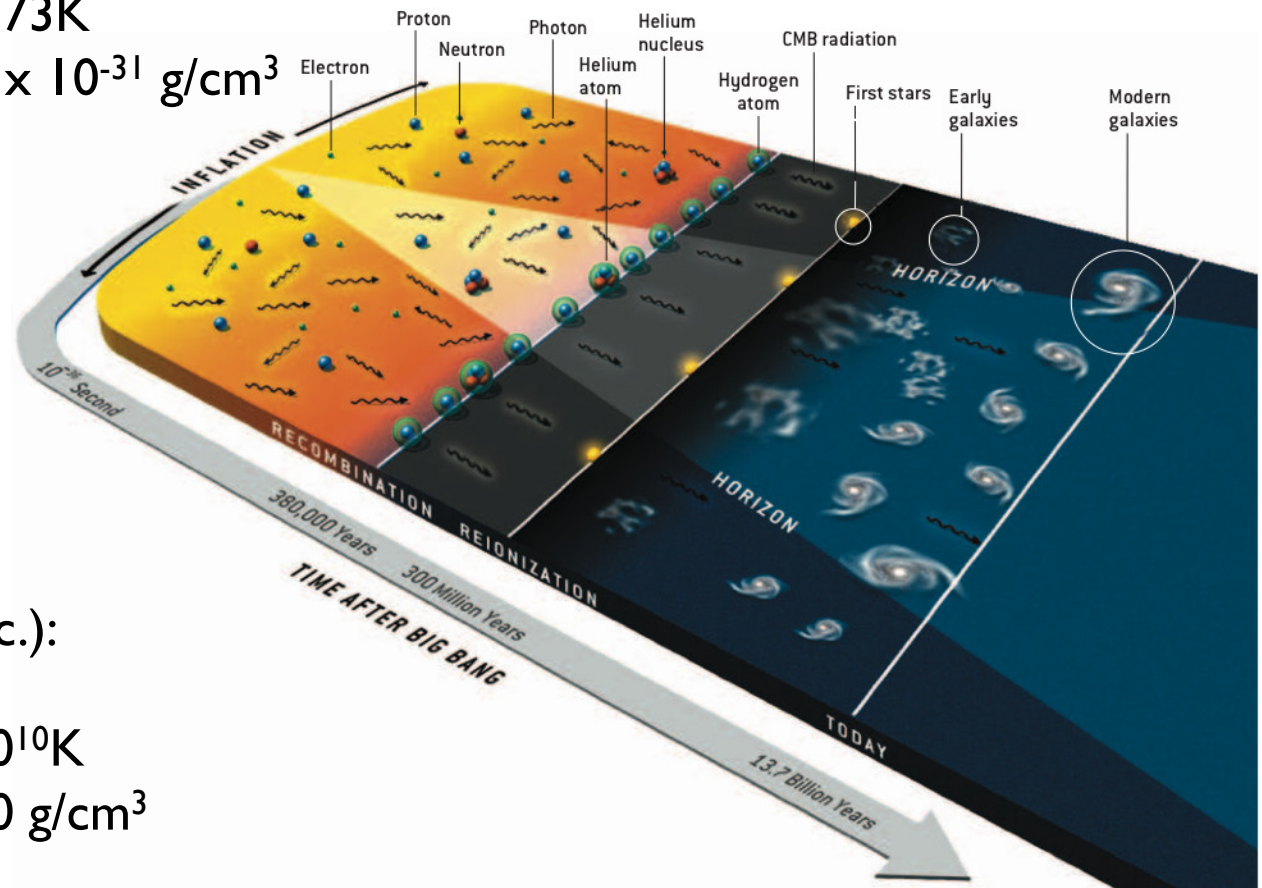


■ cosmological considerations

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- $z \approx 10^{10}$ ($t \approx 1 \text{ sec.}$):

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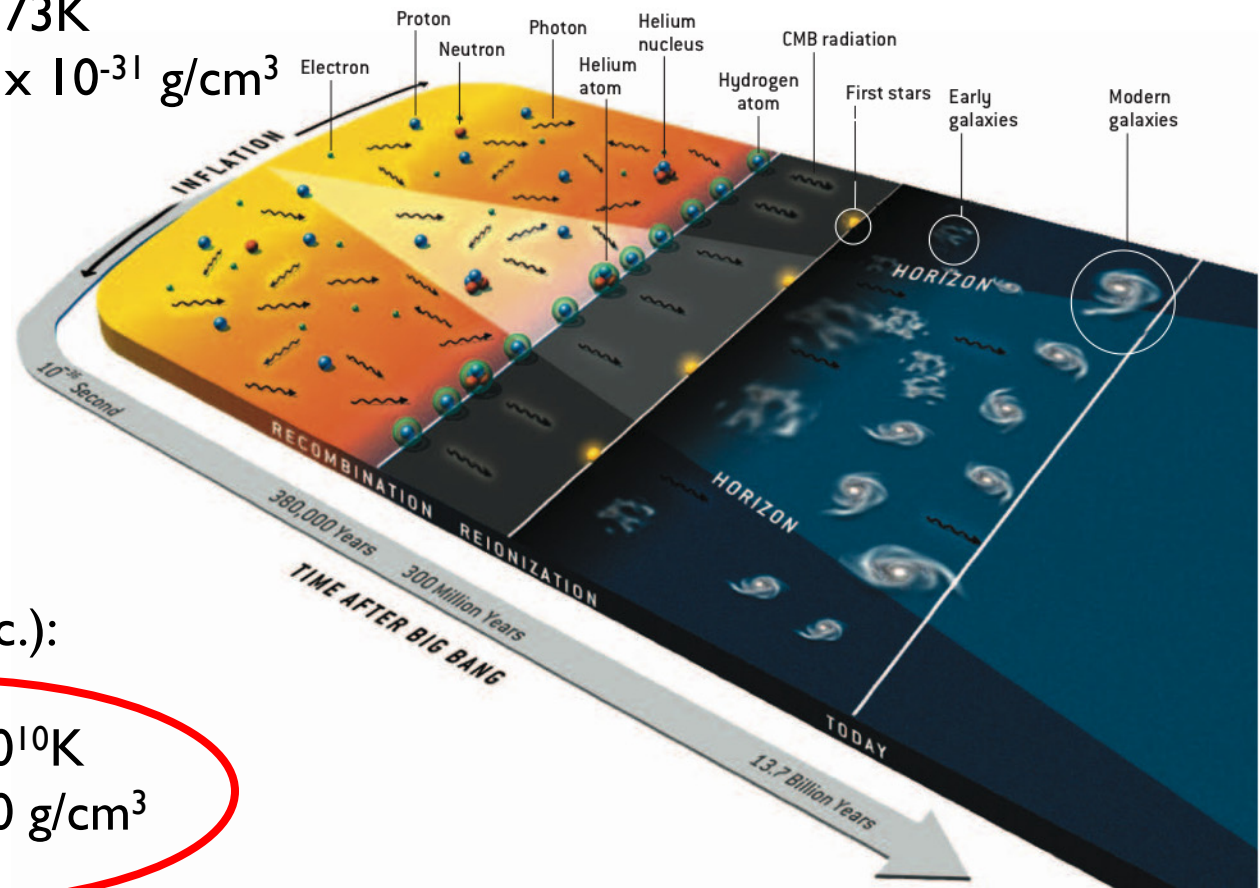
$$\rho_b \approx 80 \text{ g/cm}^3$$

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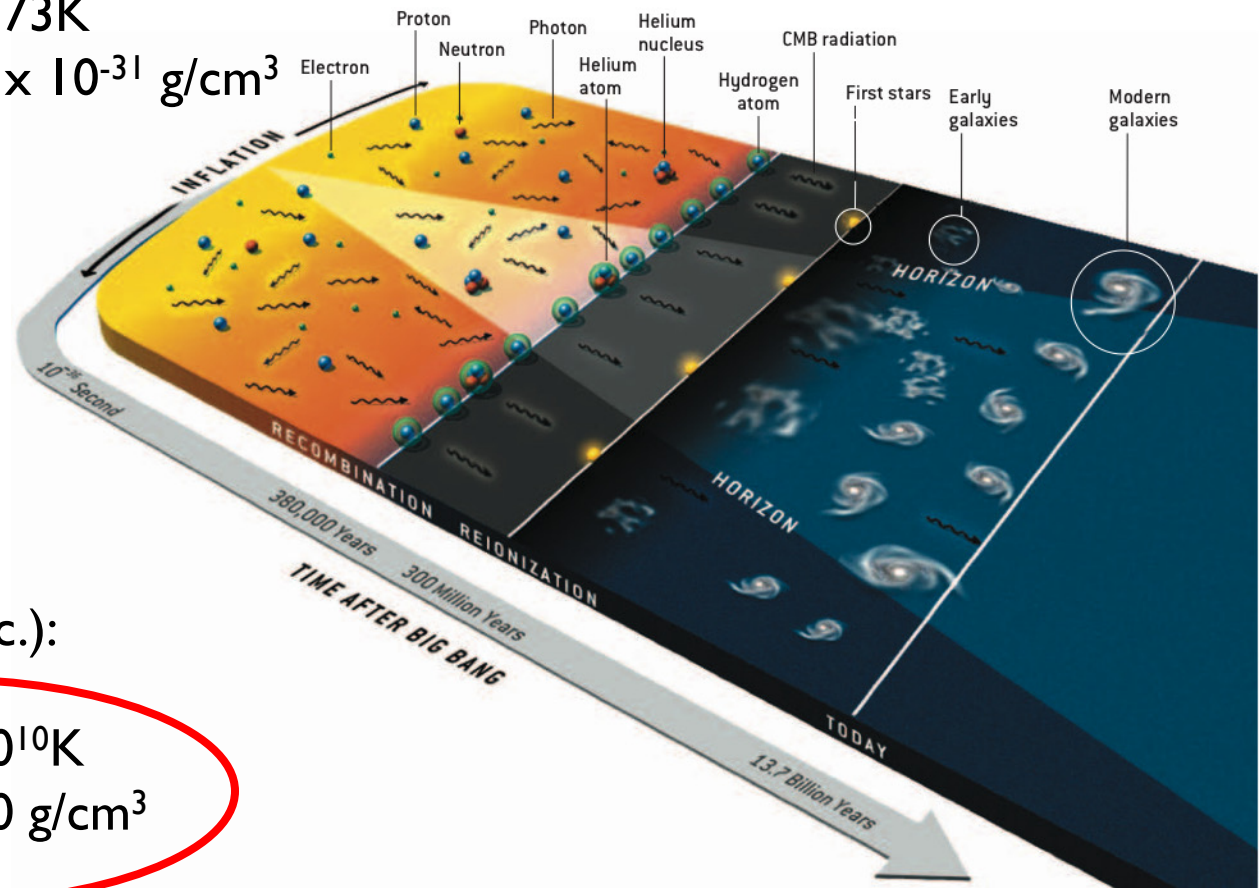
sufficient for nuclear fusion!

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sufficient for nuclear fusion → “Big Bang Nucleosynthesis” !?

■ **BBN:**

Alpher, Bethe & Gamow (1948) “The Origin of Chemical Elements”

PHYSICAL REVIEW VOLUME 73, NUMBER 7 APRIL 1, 1948

Letters to the Editor

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The Origin of Chemical Elements

R. A. ALPHER*
Applied Physics Laboratory, The Johns Hopkins University,
Silver Spring, Maryland

AND
H. BETHE
Cornell University, Ithaca, New York

AND
G. GAMOW
The George Washington University, Washington, D. C.
February 18, 1948

As pointed out by one of us,¹ various nuclear species must have originated not as the result of an equilibrium corresponding to a certain temperature and density, but rather as a consequence of a continuous building-up process arrested by a rapid expansion and cooling of the primordial matter. According to this picture, we must imagine the early stage of matter as a highly compressed neutron gas (overheated neutral nuclear fluid) which started decaying into protons and electrons when the gas pressure fell down as the result of universal expansion. The radiative capture of the still remaining neutrons by the newly formed protons must have led first to the formation of deuterium nuclei, and the subsequent neutron captures resulted in the building up of heavier and heavier nuclei. It must be remembered that, due to the comparatively short time allowed for this process,² the building up of heavier nuclei must have proceeded just above the upper fringe of the stable elements (short-lived Fermi elements), and the present frequency distribution of various atomic species was attained only somewhat later as the result of adjustment of their electric charges by β -decay.

Thus the observed slope of the abundance curve must not be related to the temperature of the original neutron gas, but rather to the time period permitted by the expansion process. Also, the individual abundances of various nuclear species must depend not so much on their intrinsic stabilities (mass defects) as on the values of their neutron capture cross sections. The equations governing such a building-up process apparently can be written in the form:

$$\frac{dn_i}{dt} = f(t)(\sigma_{i-1}n_{i-1} - \sigma_i n_i) \quad i = 1, 2, \dots, 238, \quad (1)$$

where n_i and σ_i are the relative numbers and capture cross sections for the nuclei of atomic weight i , and where $f(t)$ is a factor characterizing the decrease of the density with time.

We may remark at first that the building-up process was apparently completed when the temperature of the neutron gas was still rather high, since otherwise the observed abundances would have been strongly affected by the resonances in the region of the slow neutrons. According to Hughes,³ the neutron capture cross sections of various elements (for neutron energies of about 1 Mev) increase exponentially with atomic number halfway up the periodic system, remaining approximately constant for heavier elements.

Using these cross sections, one finds by integrating Eqs. (1) as shown in Fig. 1 that the relative abundances of various nuclear species decrease rapidly for the lighter elements and remain approximately constant for the elements heavier than silver. In order to fit the calculated curve with the observed abundances⁴ it is necessary to assume the integral of ρdt during the building-up period is equal to 5×10^8 g sec./cm³.

On the other hand, according to the relativistic theory of the expanding universe⁵ the density dependence on time is given by $\rho \approx 10^9/t^3$. Since the integral of this expression diverges at $t=0$, it is necessary to assume that the building-up process began at a certain time t_0 , satisfying the relation:

$$\int_{t_0}^{\infty} (10^9/t^3) dt \approx 5 \times 10^8, \quad (2)$$

which gives us $t_0 \approx 20$ sec. and $\rho_0 \approx 2.5 \times 10^8$ g sec./cm³. This result may have two meanings: (a) for the higher densities existing prior to that time the temperature of the neutron gas was so high that no aggregation was taking place, (b) the density of the universe never exceeded the value 2.5×10^8 g sec./cm³ which can possibly be understood if we

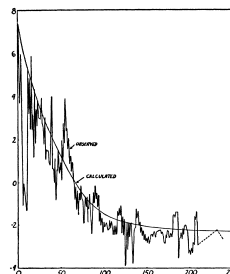


FIG. 1.
Log of relative abundance
Atomic weight

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$$\left\langle \frac{dn}{dt} + 3Hn = -\langle \sigma v \rangle (n^2 - n_{eq}^2) \right\rangle \longleftrightarrow \frac{dn_i}{dt} = f(t)(\sigma_{i-1}n_{i-1} - \sigma_i n_i) \quad i = 1, 2, \dots, 238, \quad (1)$$

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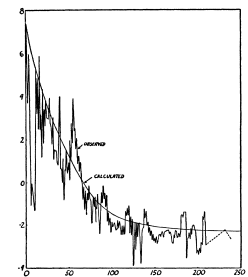


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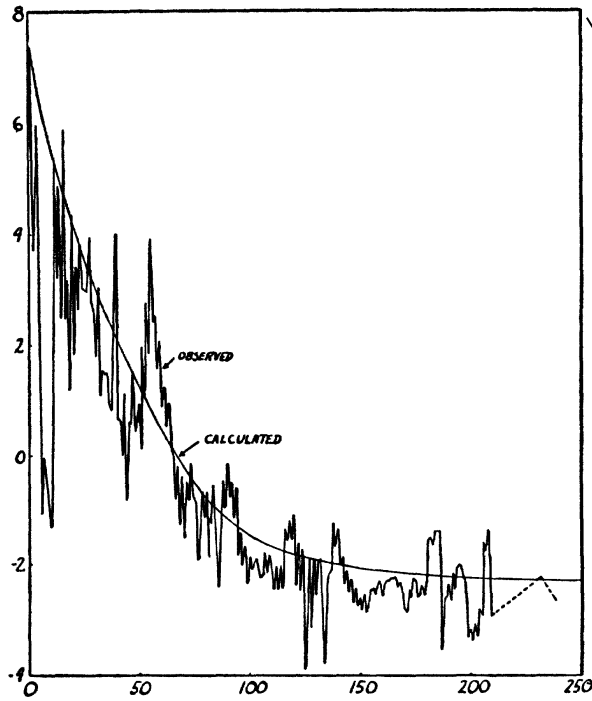


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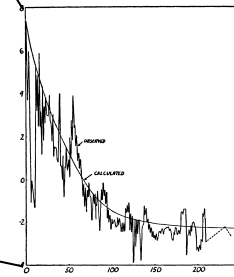


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all elements are produced in BBN !?

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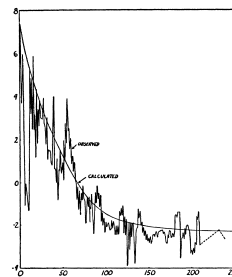


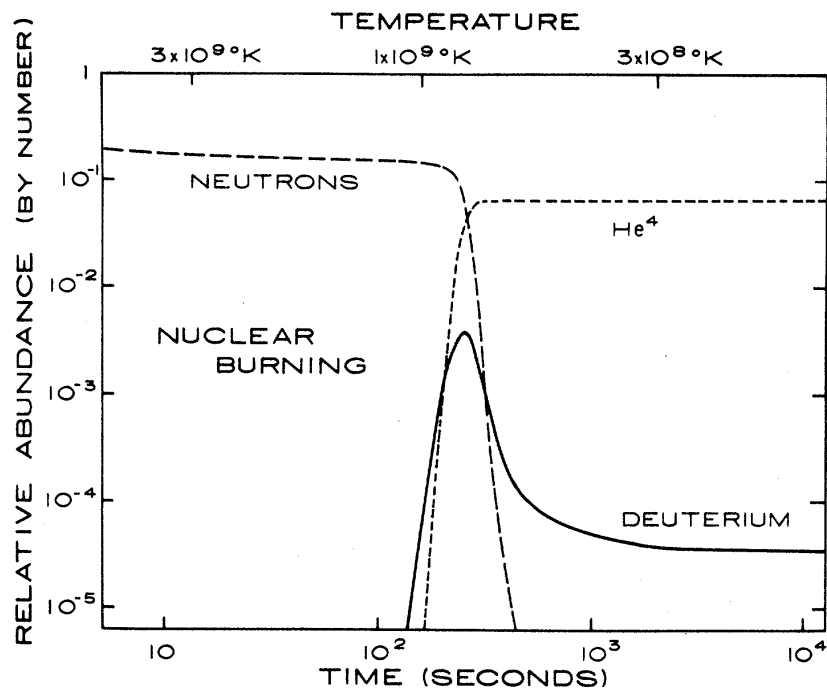
FIG. 1.
Log of relative abundance
Atomic weight

this is the whole publication!

- **BBN:**

Peebles (1966)

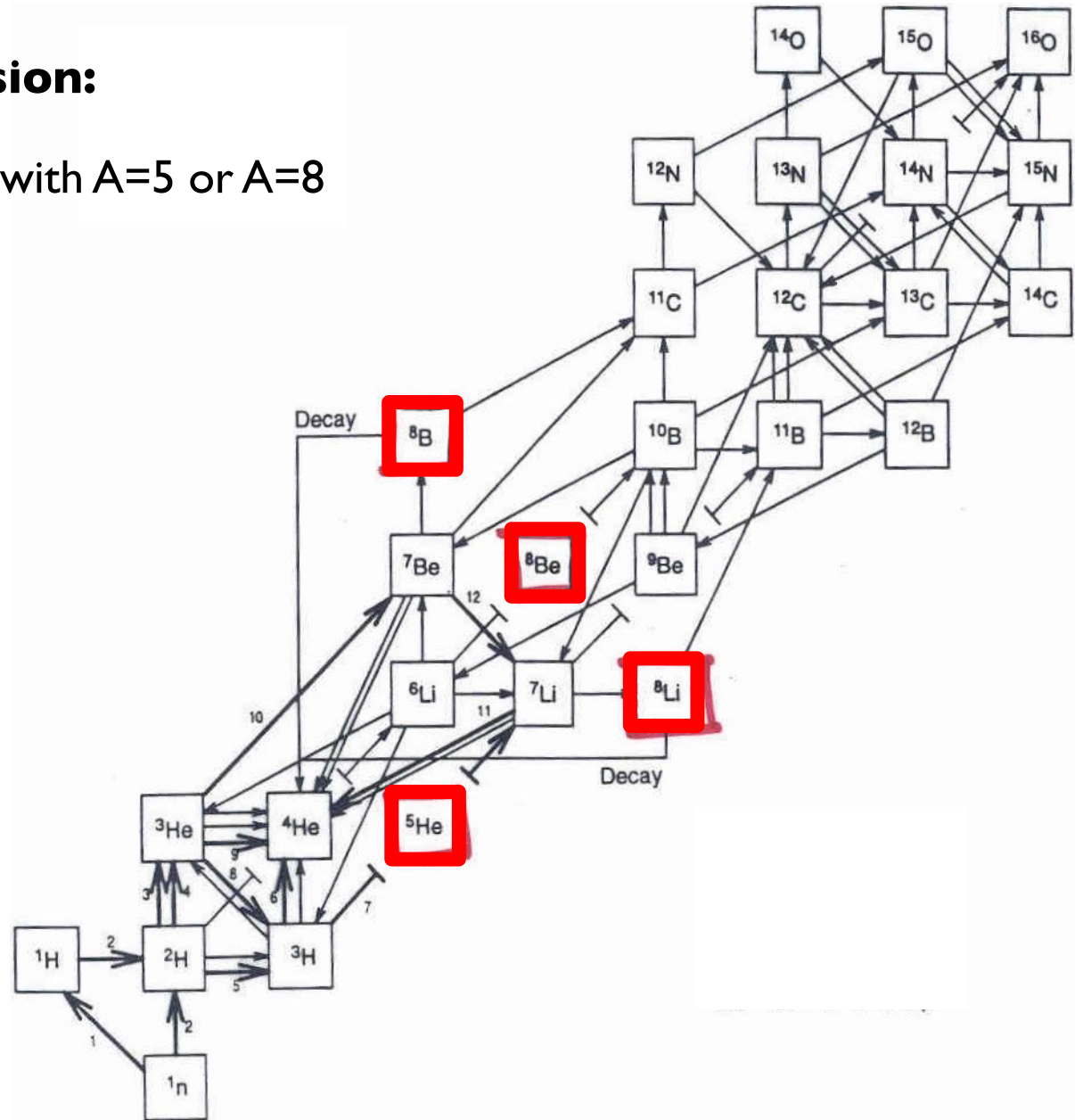
“Primeval Helium abundance and the primeval Fireball”



$\frac{dn}{dt} + 3Hn = -\langle\sigma v\rangle(n^2 - n_{eq}^2)$: detailed calculation of the temporal evolution of element abundances...

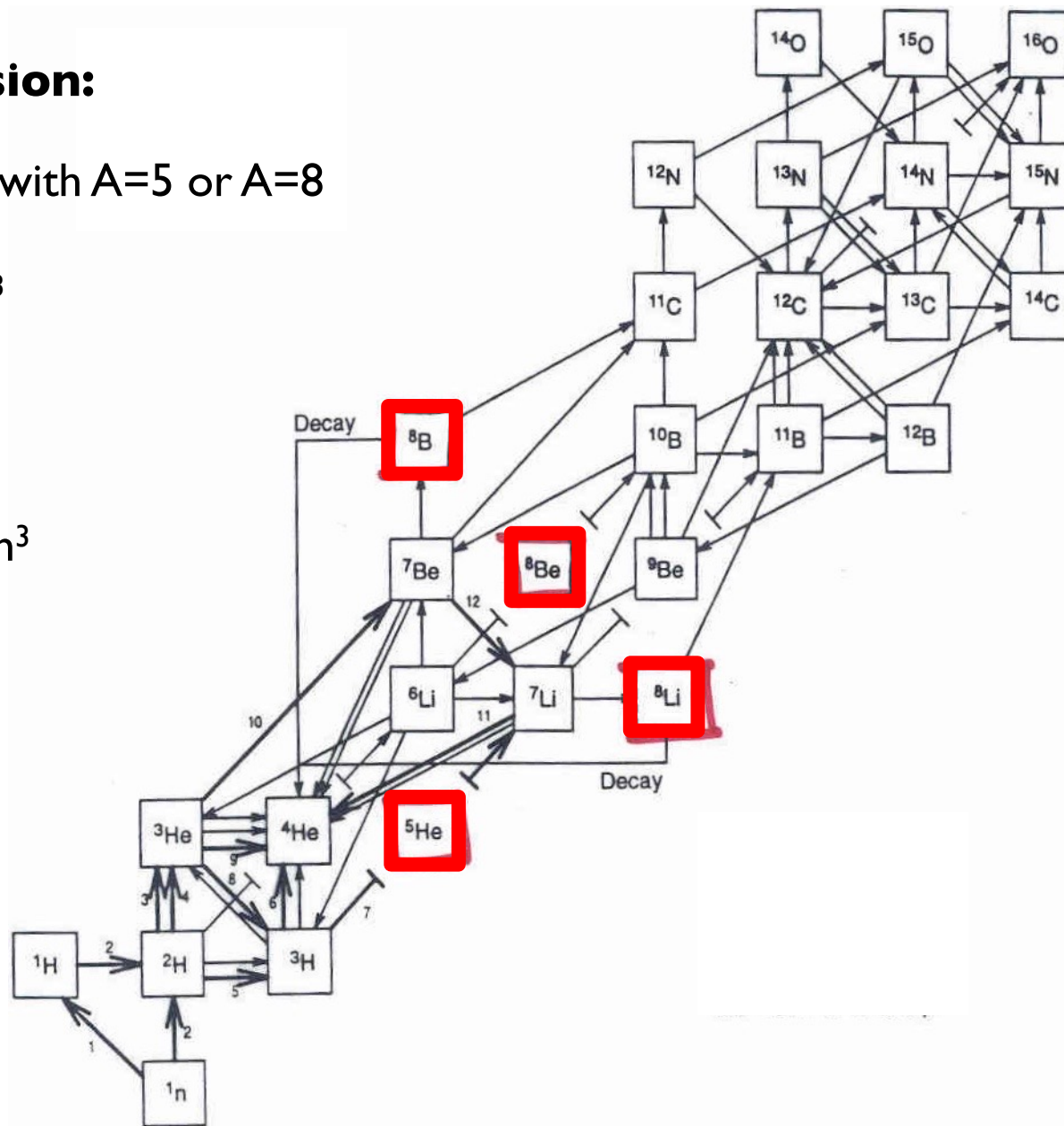
- **BBN vs. stellar fusion:**

- no stable elements with $A=5$ or $A=8$



▪ BBN vs. stellar fusion:

- no stable elements with $A=5$ or $A=8$
- BBN: $\rho_b \approx 80 \text{ g/cm}^3$
- sun: $\rho_b \approx 150 \text{ g/cm}^3$



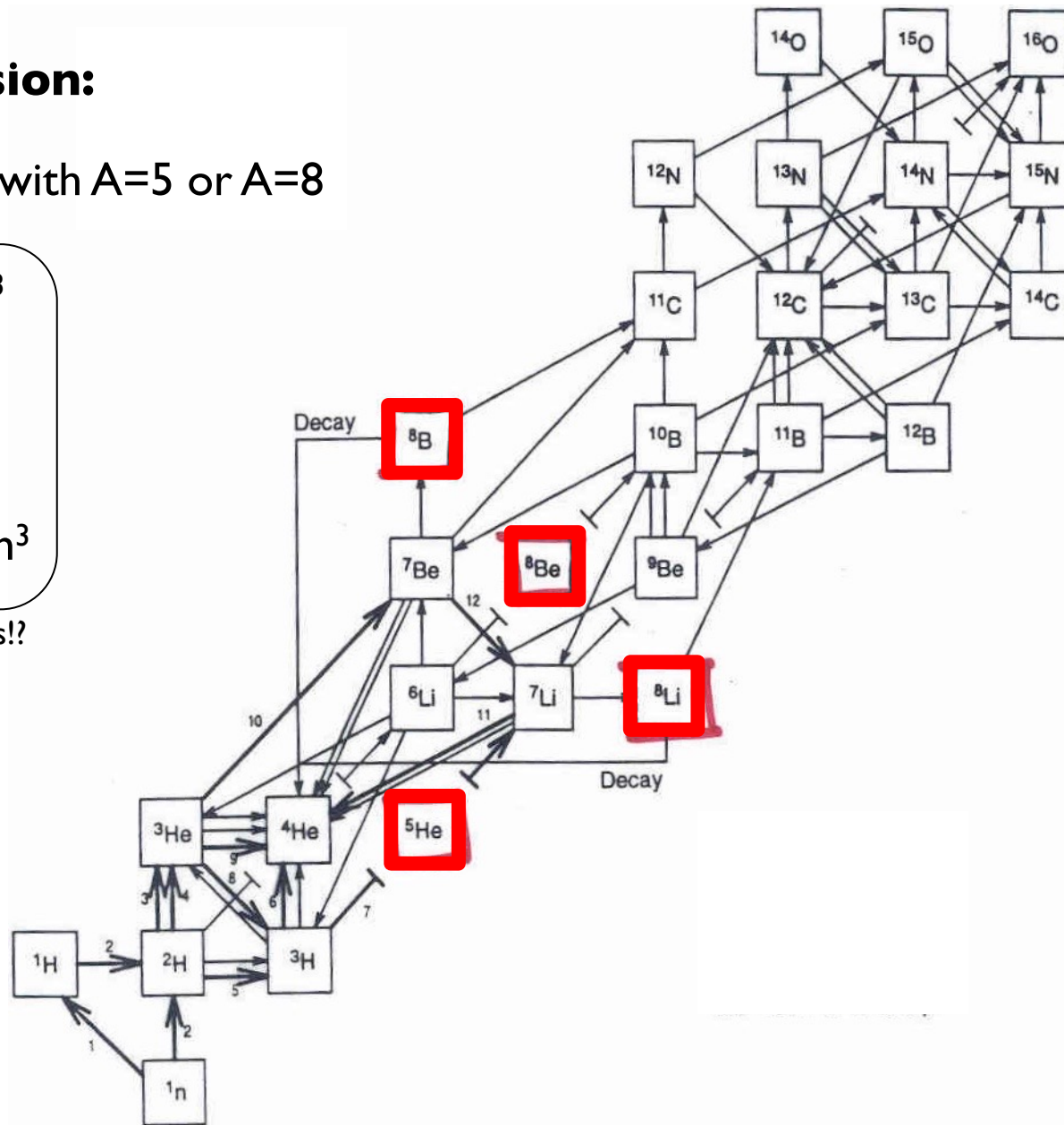
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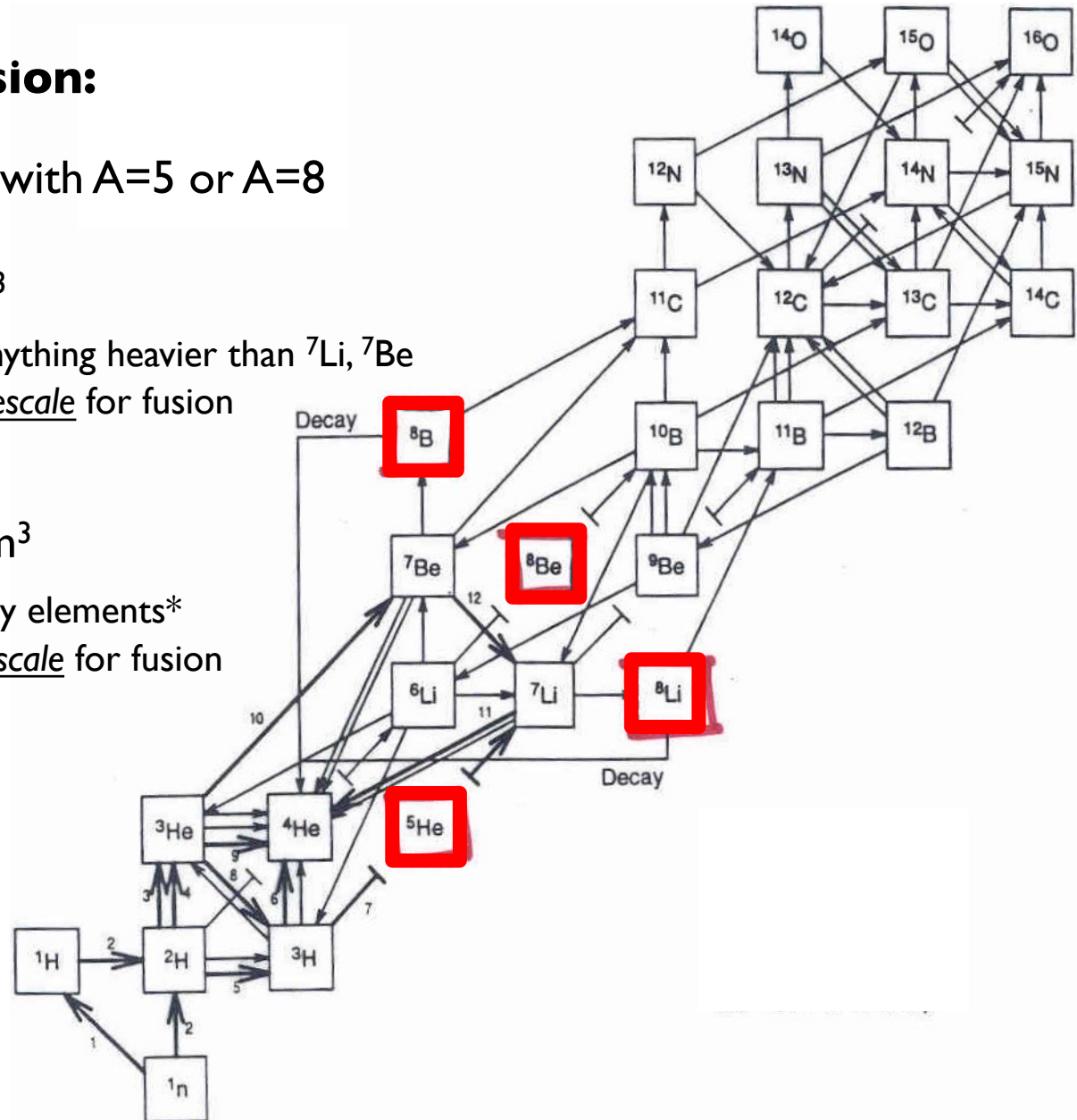
- sun: $\rho_b \approx 150 \text{ g/cm}^3$

similar conditions!?



▪ BBN vs. stellar fusion:

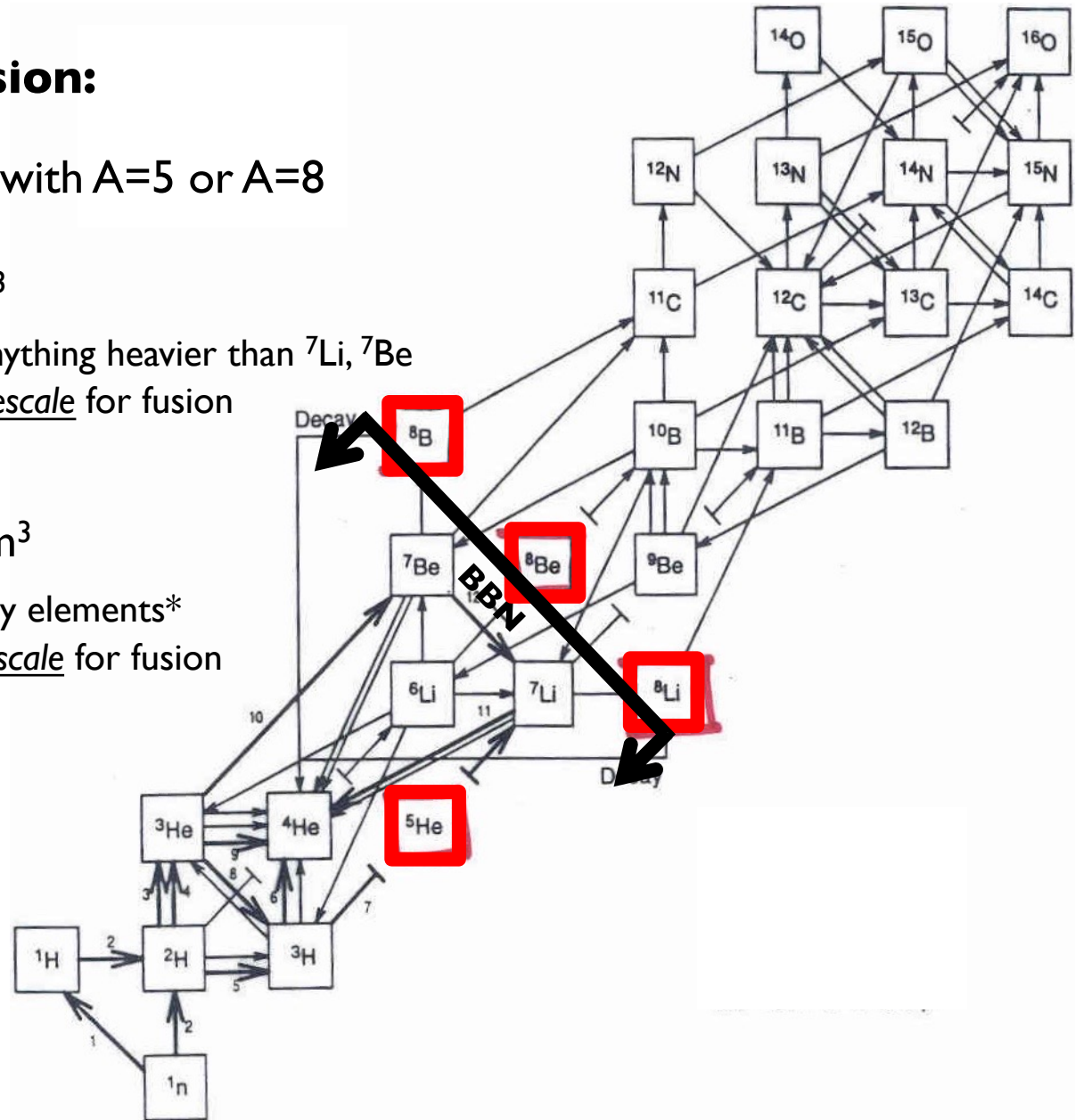
- no stable elements with $A=5$ or $A=8$
- BBN: $\rho_b \approx 80 \text{ g/cm}^3$
→ unable to process anything heavier than ${}^7\text{Li}$, ${}^7\text{Be}$
due to very short timescale for fusion
- sun: $\rho_b \approx 150 \text{ g/cm}^3$
→ able to process heavy elements*
due to extended timescale for fusion



*via unstable ${}^8\text{Be}$ (“triple- α process”)

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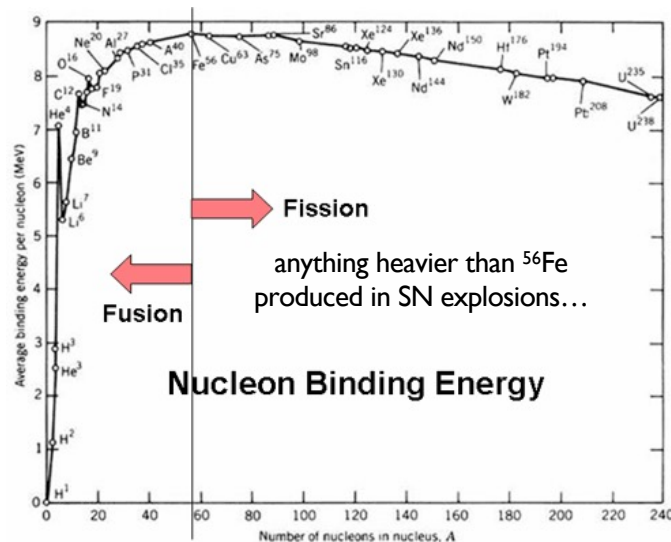
- **BBN summary:**

- production of H , ${}^2\text{H}$, ${}^3\text{H}$, ${}^3\text{He}$, ${}^4\text{He}$, ${}^7\text{Be}$, ${}^7\text{Li}$

${}^2\text{H}$: deuterium D

${}^3\text{H}$: tritium T

- **all other elements are in fact produced in stars...**



▪ BBN summary:

- production of H , ${}^2\text{H}$, ${}^3\text{H}$, ${}^3\text{He}$, ${}^4\text{He}$, ${}^7\text{Be}$, ${}^7\text{Li}$

${}^2\text{H}$: deuterium D

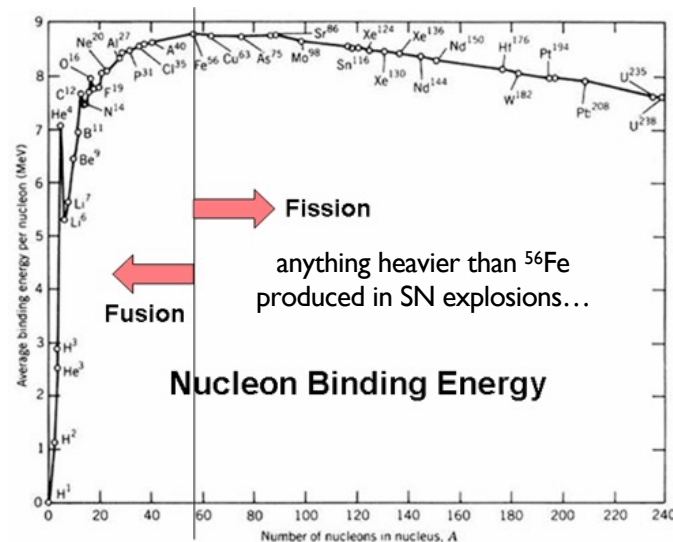
${}^3\text{H}$: tritium T

- *prediction* for mass fractions:

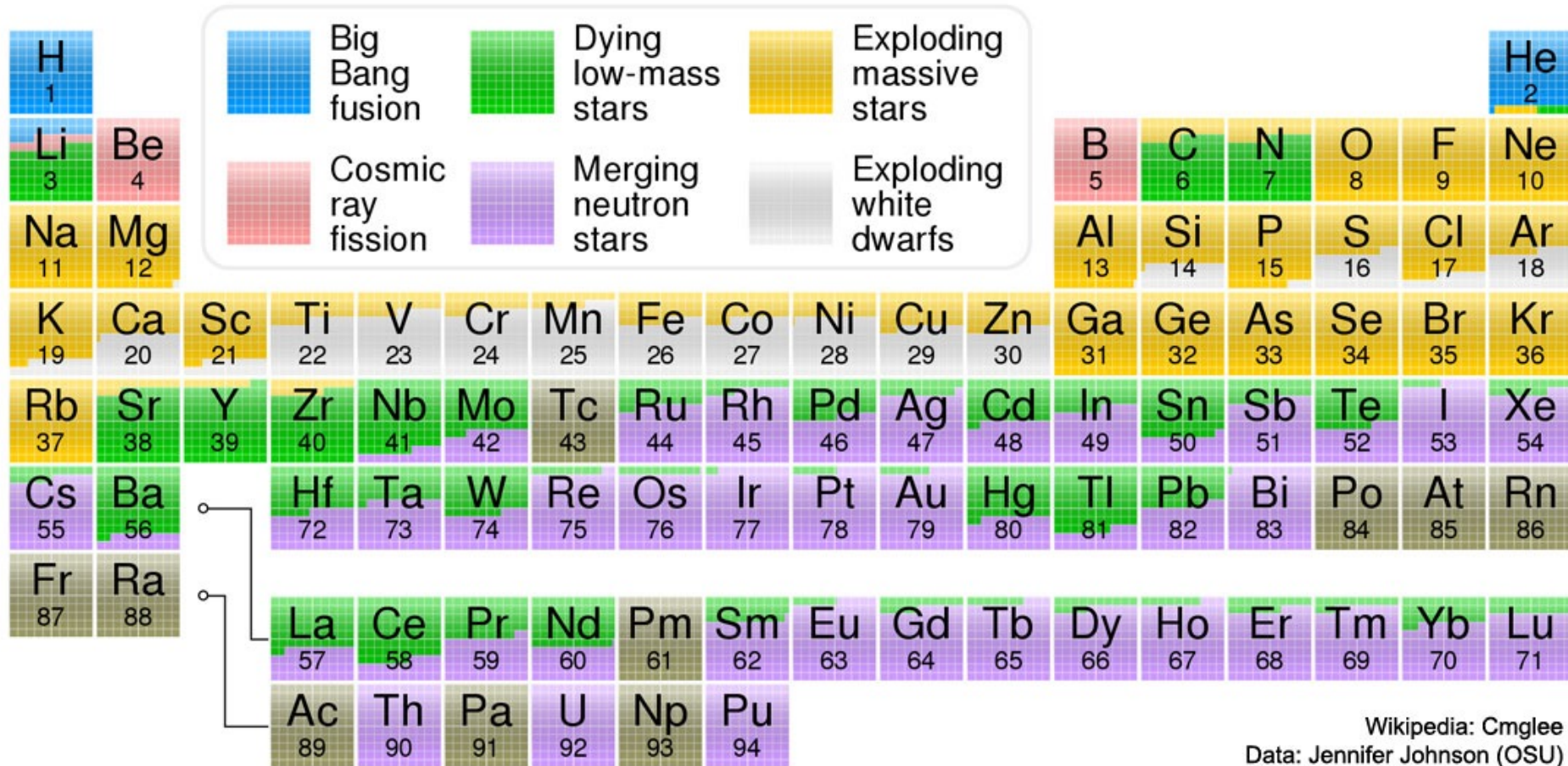
$$Y_{\text{He}} \approx \mathbf{0.24}$$

$$Y_{\text{H}} \approx \mathbf{0.73}$$

▪ all other elements are in fact produced in stars...



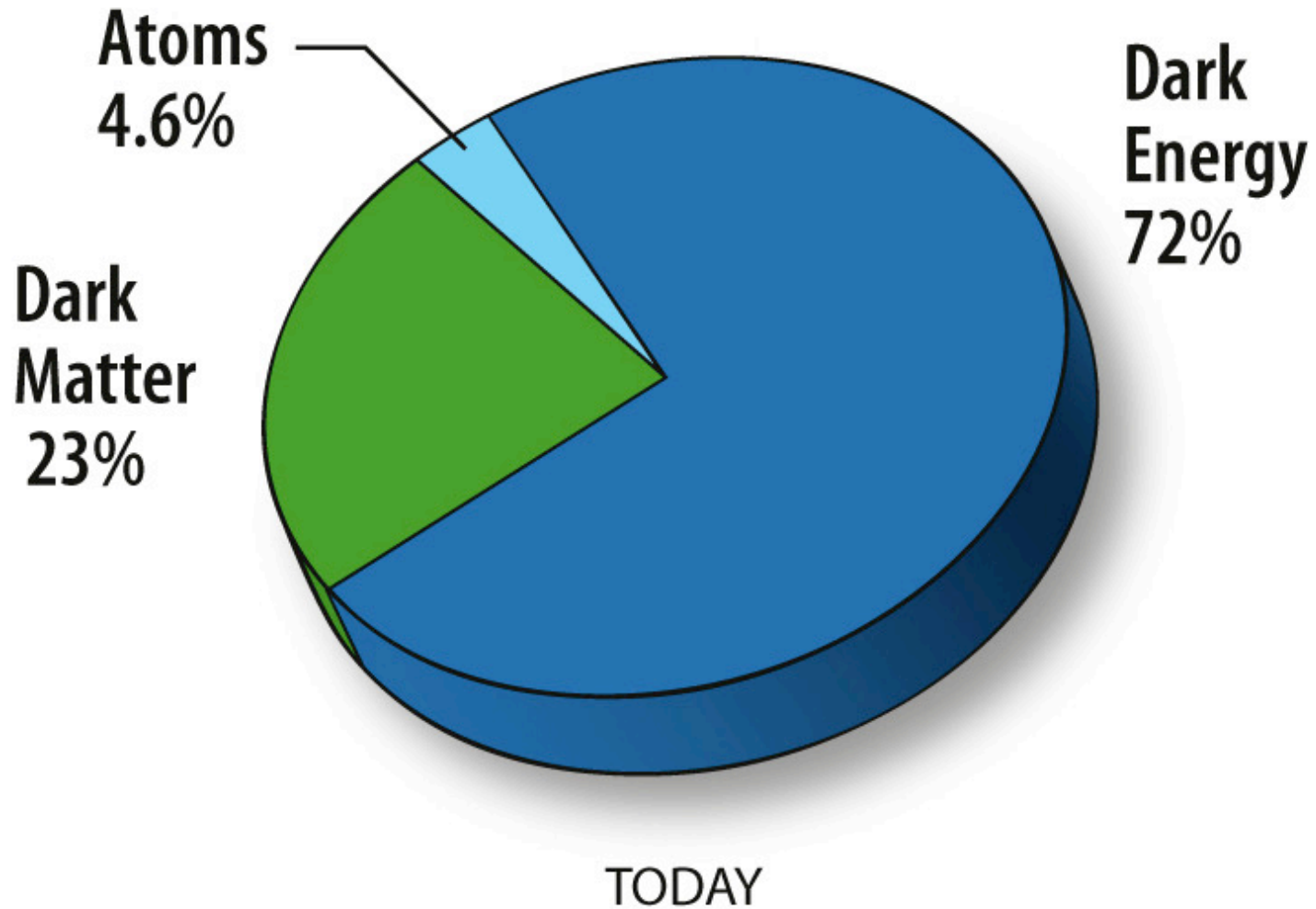
where do all the elements come from?



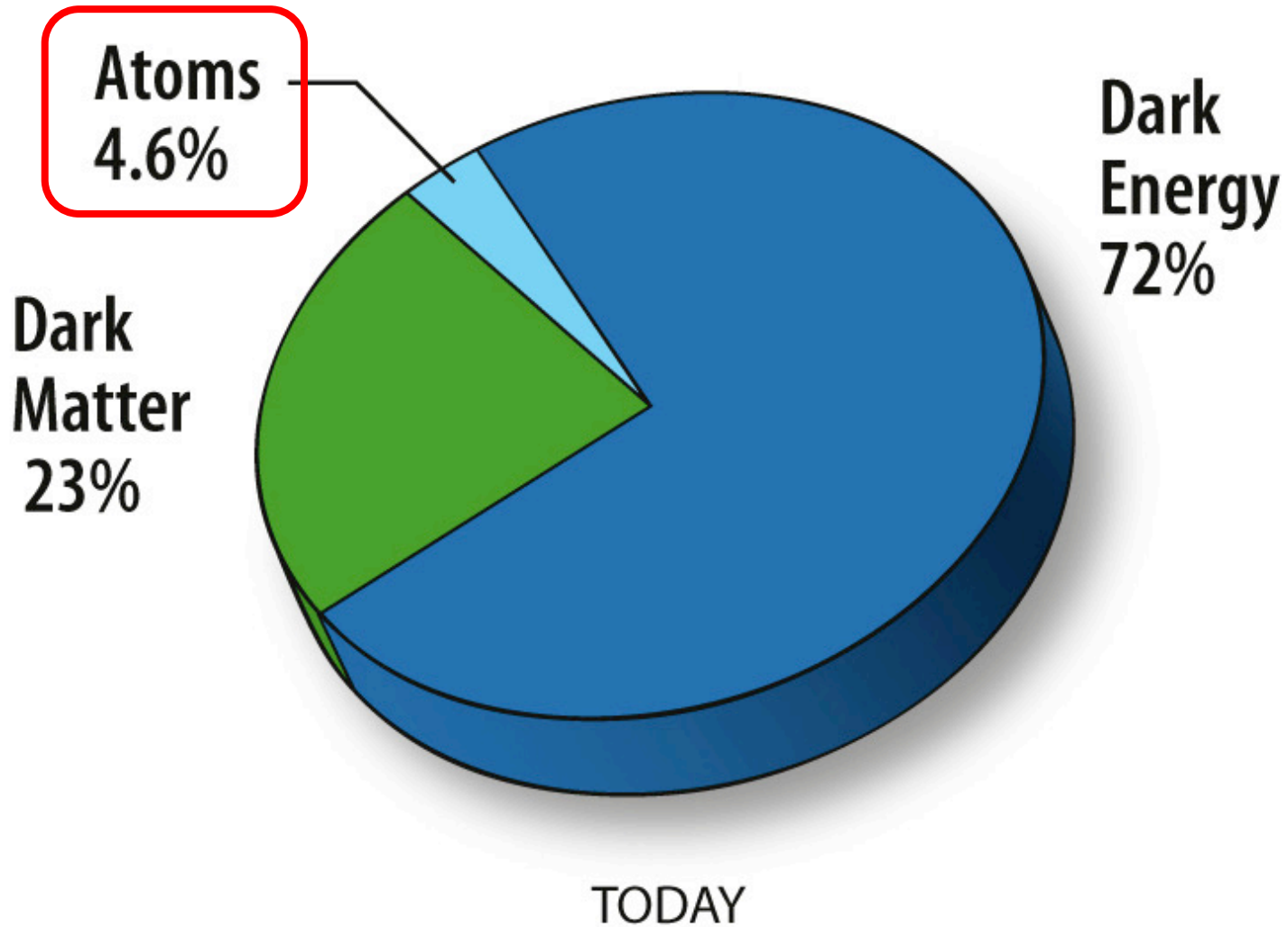
Wikipedia: Cmglee
Data: Jennifer Johnson (OSU)

(APOD 24/10/2017)

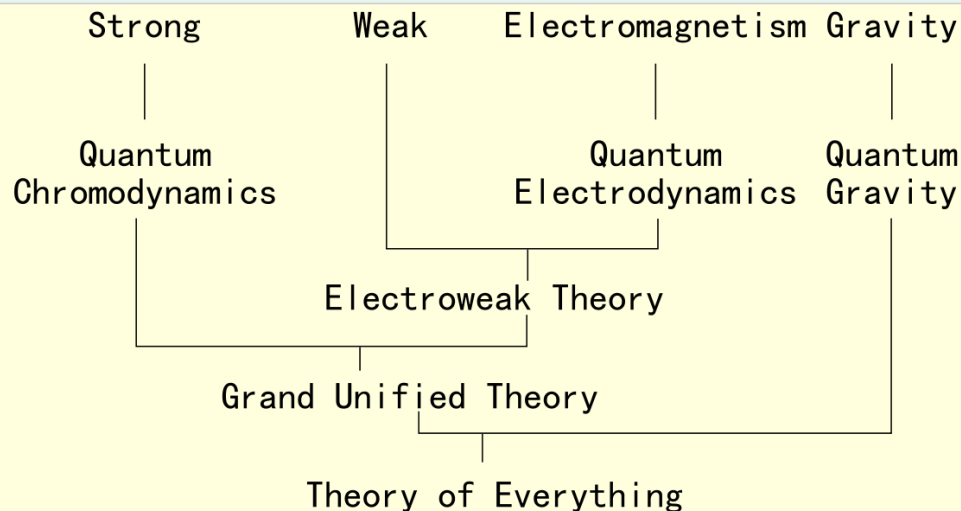
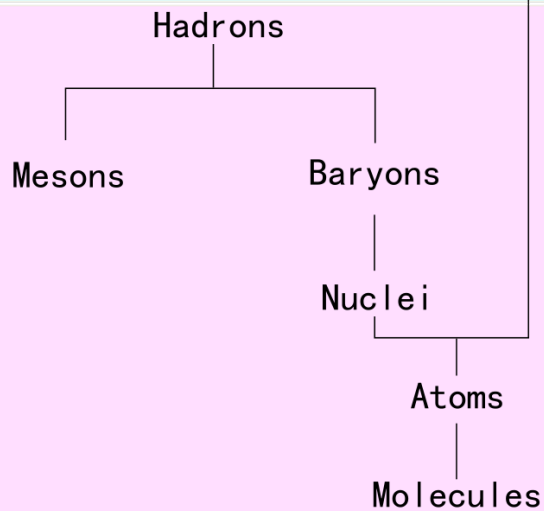
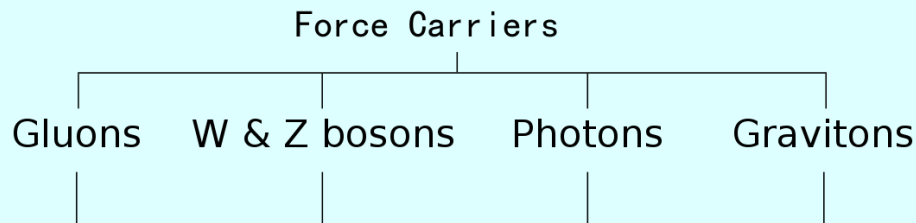
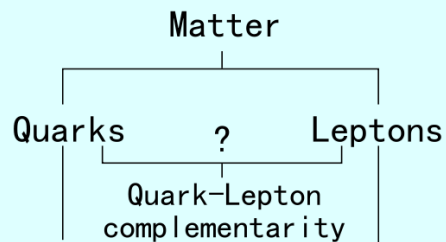
- introduction
- ***particle physics***
- synthesis of elements
- big bang nucleosynthesis
- observations



we are only talking about their nuclei here...



Elementary Particles

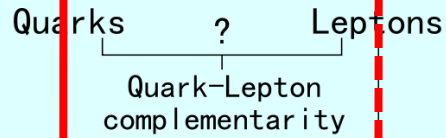


Composite Particles

Forces

Elementary Particles

Matter



Hadrons

Mesons

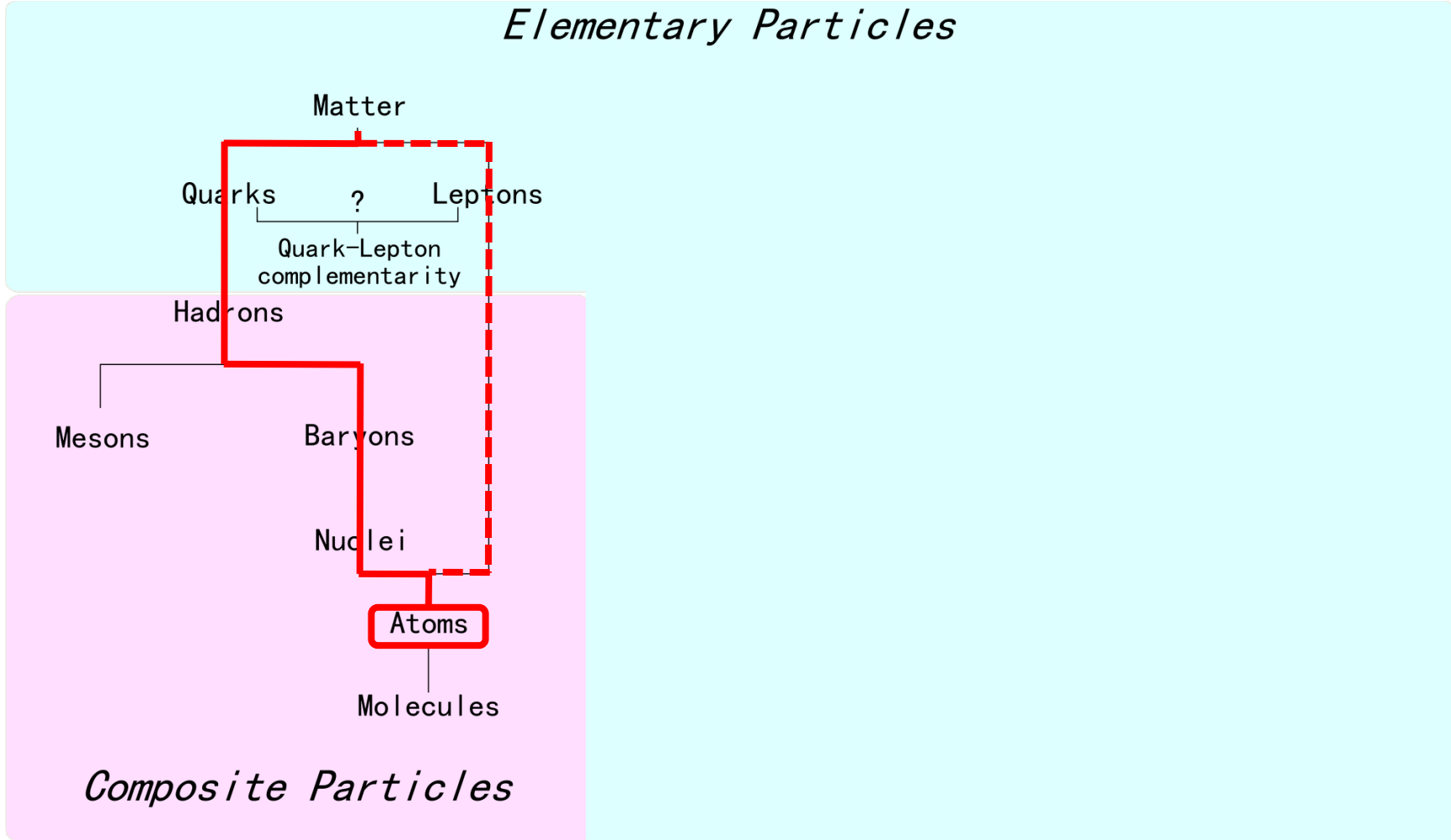
Baryons

Nuclei

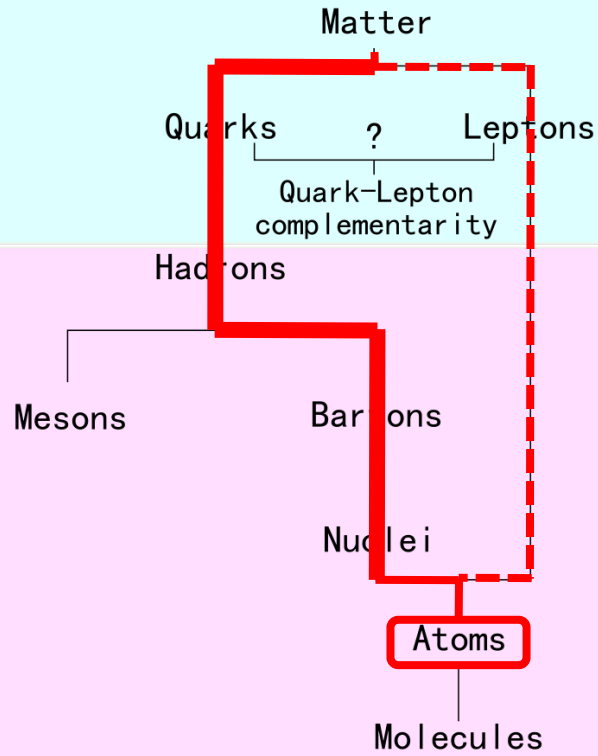
Atoms

Molecules

Composite Particles



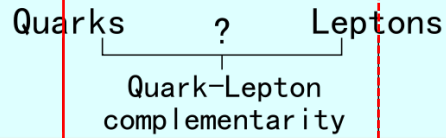
Elementary Particles



Composite Particles

Elementary Particles

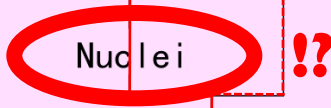
Matter



Hadrons

Mesons

Baryons

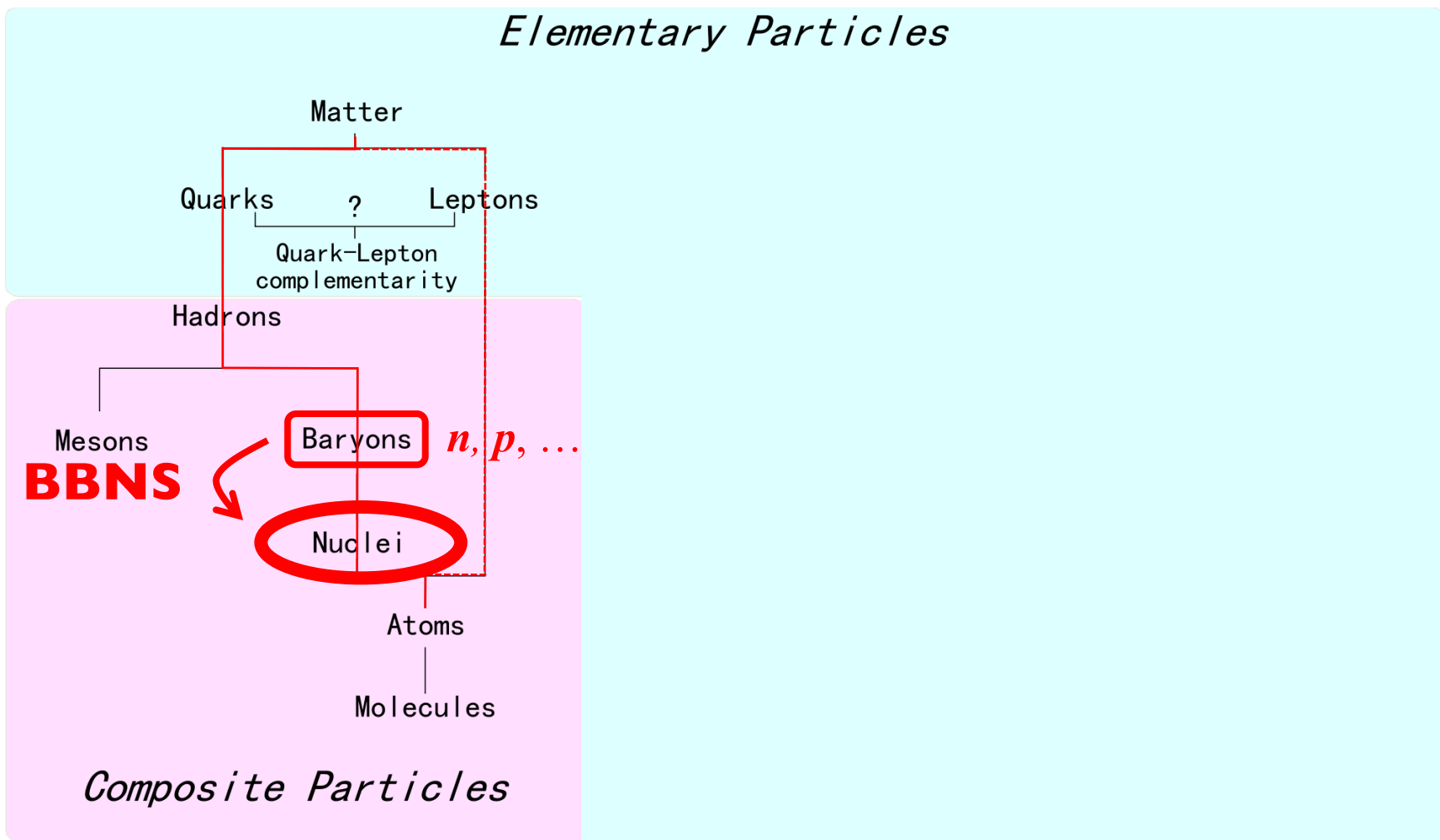


Atoms

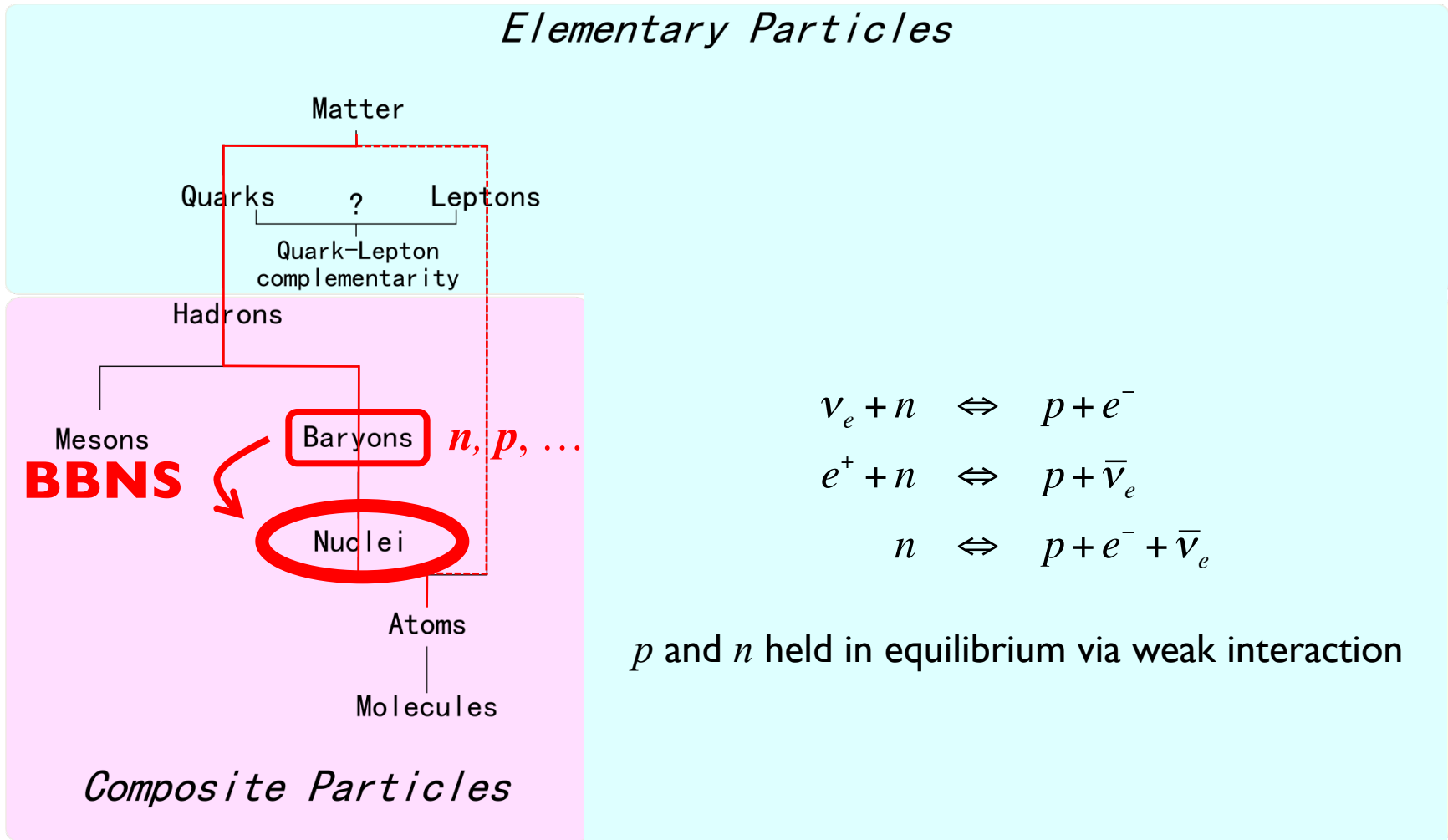
Molecules

Composite Particles

Elementary Particles



Elementary Particles



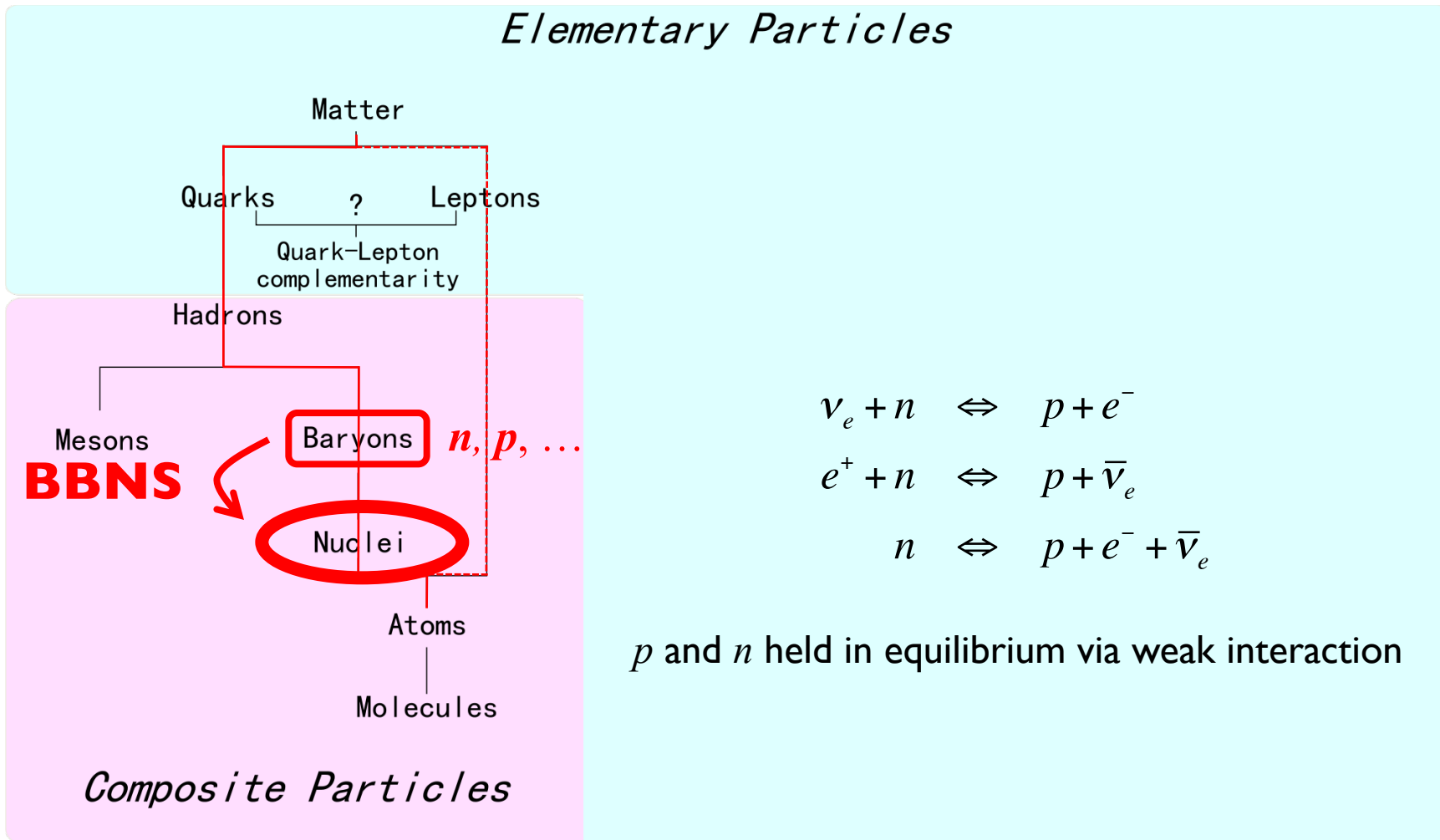
$$\nu_e + n \Leftrightarrow p + e^-$$

$$e^+ + n \Leftrightarrow p + \bar{\nu}_e$$

$$n \Leftrightarrow p + e^- + \bar{\nu}_e$$

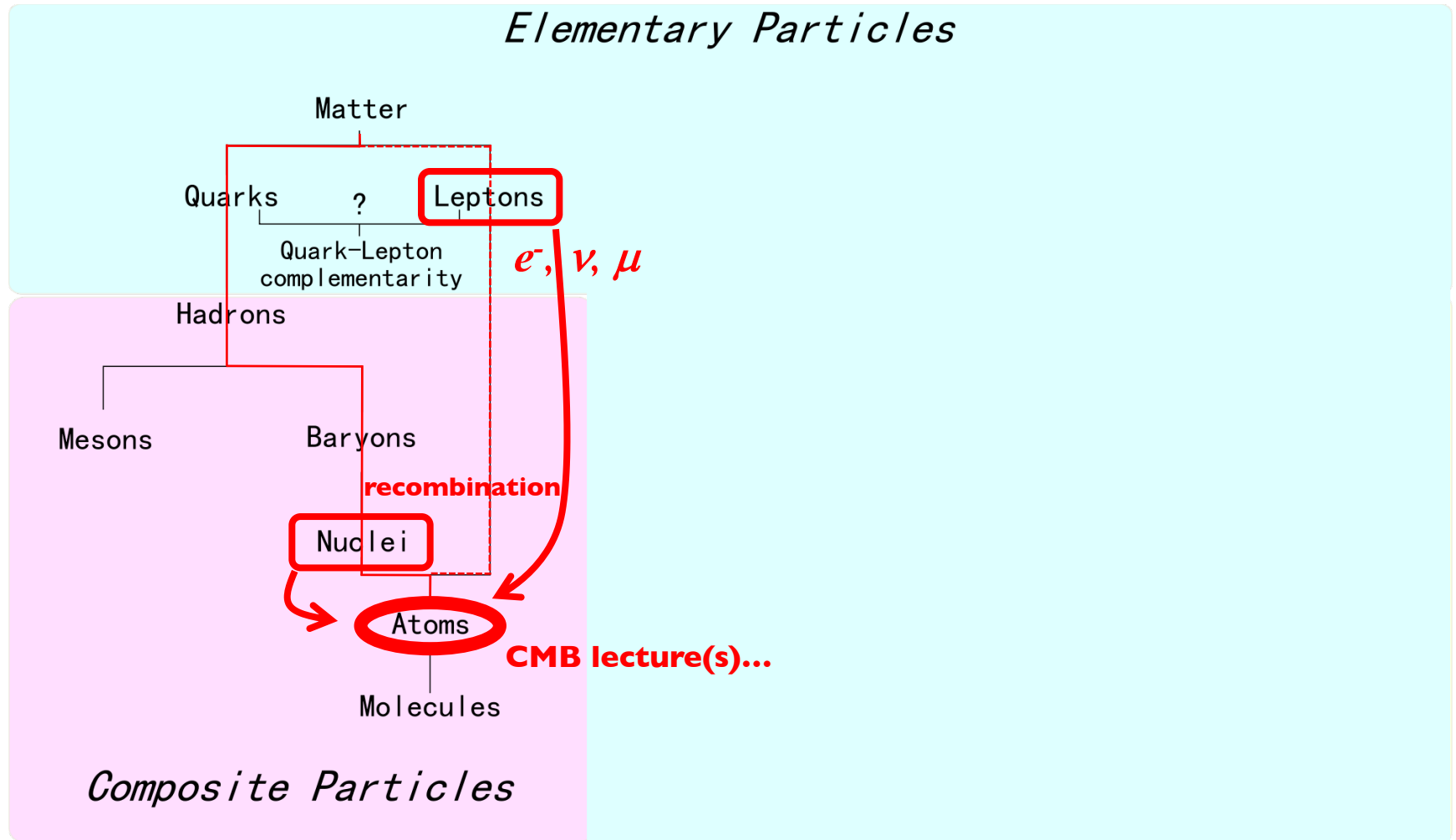
p and n held in equilibrium via weak interaction

Elementary Particles



BBNS can only start after neutrino decoupling!

Elementary Particles



Event	time t	redshift z	temperature T
Inflation	10^{-34} s (?)	–	–
Baryogenesis	?	?	?
EW phase transition	20 ps	10^{15}	100 GeV
QCD phase transition	20 μ s	10^{12}	150 MeV
Dark matter freeze-out	?	?	?
Neutrino decoupling	1 s	6×10^9	1 MeV
Electron-positron annihilation	6 s	2×10^9	500 keV
Big Bang nucleosynthesis	3 min	4×10^8	100 keV
Matter-radiation equality	60 kyr	3400	0.75 eV
Recombination	260–380 kyr	1100–1400	0.26–0.33 eV
Photon decoupling	380 kyr	1000–1200	0.23–0.28 eV
Reionization	100–400 Myr	11–30	2.6–7.0 meV
Dark energy-matter equality	9 Gyr	0.4	0.33 meV
Present	13.8 Gyr	0	0.24 meV



- introduction
- particle physics
- ***synthesis of elements***
- big bang nucleosynthesis
- observations

H, ^2H , ^3H , ^3He , ^4He , ^7Be , ^7Li

H, ²H, ³H, ³He, ⁴He, ⁷Be, ⁷Li



H, ²H, ³H, ³He, ⁴He, ⁷Be, ⁷Li



this requires free electrons,
but they are still coupled to the photons...

H, ²H, ³H, ³He, ⁴He, ⁷Be, ⁷Li

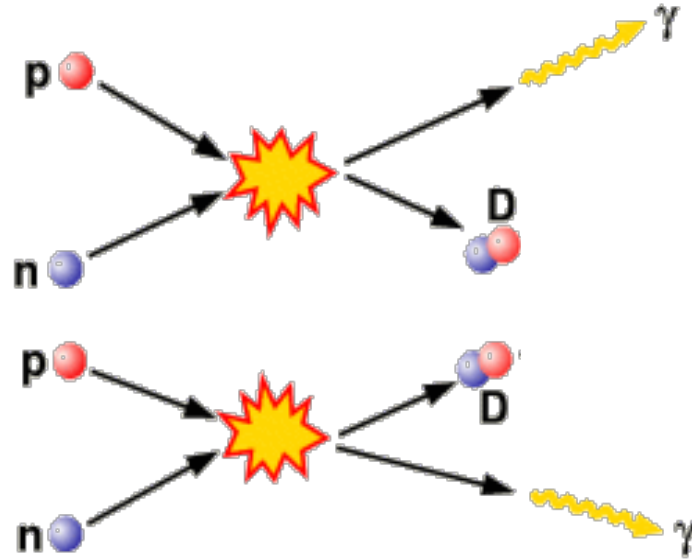


this requires free electrons,
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Note: in BBNS we are forming *nuclei* and not atoms!

$H, {}^2H, {}^3H, {}^3He, {}^4He, {}^7Be, {}^7Li$

- Hydrogen: $(p + e^- \rightarrow H + \gamma)$
- Deuterium: $p + n \rightarrow D + \gamma$



two-body processes only,
as probability for others too low
(‘bottom-up’ fusion of He^4)

H, ²H, ³H, ³He, ⁴He, ⁷Be, ⁷Li

- Hydrogen: $(p + e^- \rightarrow H + \gamma)$
- Deuterium: $p + n \rightarrow D + \gamma$
- all $A > 2$ nuclei require D and n for synthesis

H, ²H, ³H, ³He, ⁴He, ⁷Be, ⁷Li

- Hydrogen: $(p + e^- \rightarrow H + \gamma)$
- Deuterium: $p + n \rightarrow D + \gamma$
 $E_b \approx 2\text{MeV}$
- all $A > 2$ nuclei require D and n for synthesis

$H, {}^2H, {}^3H, {}^3He, {}^4He, {}^7Be, {}^7Li$ 

$$E_b \approx 2MeV ; kT_\nu \approx 0.8MeV (T @ \text{neutrino decoupling})$$

- all $A > 2$ nuclei require D and n for synthesis

H, ²H, ³H, ³He, ⁴He, ⁷Be, ⁷Li

$$E_b \approx 2\text{MeV} ; \quad kT_\nu \approx 0.8\text{MeV} \quad (T @ \text{neutrino decoupling})$$

- all $A > 2$ nuclei require D and n for synthesis

but: D easily photo-dissociated by γ until $kT_D \approx 0.086\text{MeV}$ (ca. $t \approx 100s^*$)

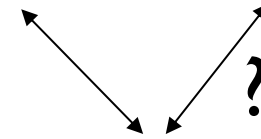
* $\frac{T}{1\text{MeV}} \approx 1.5g_*^{-1/4} \left(\frac{1s}{t}\right)^{1/2}$

H, ²H, ³H, ³He, ⁴He, ⁷Be, ⁷Li

$$E_b \approx 2\text{MeV} ; kT_\nu \approx 0.8\text{MeV} \text{ (} T @ \text{ neutrino decoupling)}$$

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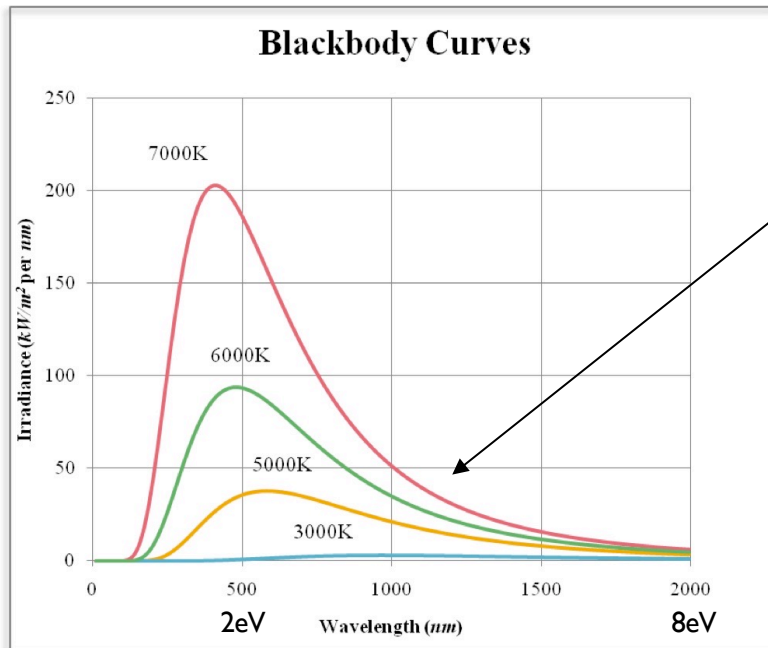
H, ²H, ³H, ³He, ⁴He, ⁷Be, ⁷Li



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for temperatures $kT_D < E_b$
there will still be lots of higher energy photons in the Planck distribution

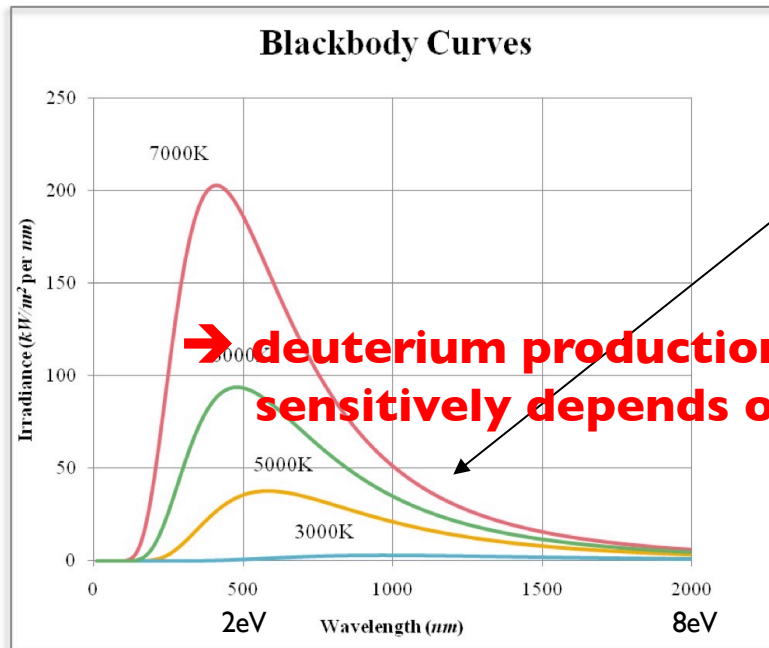
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for temperatures $kT_D < E_b$

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H, ^2H , ^3H , ^3He , ^4He , ^7Be , ^7Li 

$$E_b \approx 2\text{MeV} ; kT_\nu \approx 0.8\text{MeV} \text{ (} T @ \text{ neutrino decoupling)}$$

- all $A > 2$ nuclei require D and n for synthesis

but: D easily photo-dissociated by γ until $kT_D \cong 0.086\text{MeV}$ (ca. $t \approx 100\text{s}$)

→ ‘Deuterium bottleneck’:

- too few D → important fusion agent is missing
- too much D → locks up neutrons for further synthesis

→ deuterium production (and all successive nuclei) sensitively depends on baryon-to-photon ratio!

H, ^2H , ^3H , ^3He , ^4He , ^7Be , ^7Li 

$$E_b \approx 2\text{MeV} ; \quad kT_\nu \approx 0.8\text{MeV} \quad (T @ \text{neutrino decoupling})$$

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→ ‘Deuterium bottleneck’:

- too few D → important fusion agent is missing
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- photon-to-baryon ratio* (frozen in at BBNS):

$$\eta = \frac{n_b}{n_\gamma} = 10^{-10} \eta_{10} = 10^{-10} \cdot 274 \Omega_b h^2$$

H, ²H, ³H, ³He, ⁴He, ⁷Be, ⁷Li

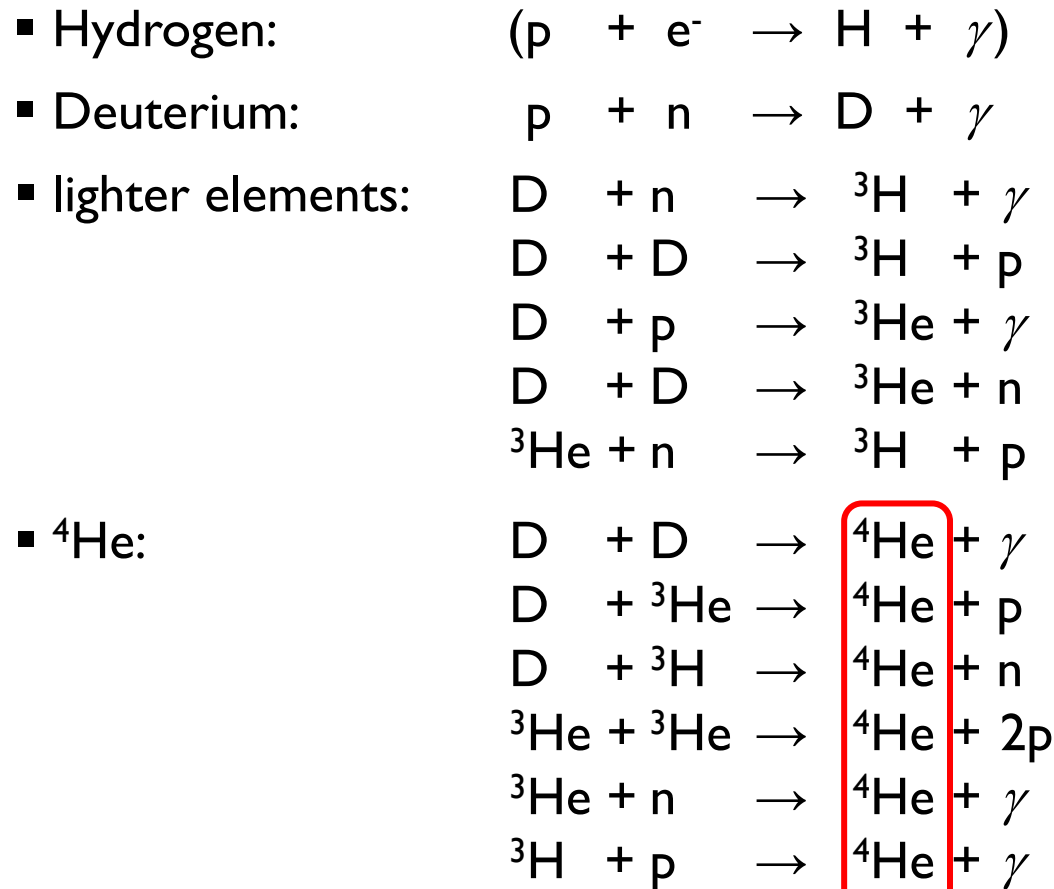
- Hydrogen: $(p + e^- \rightarrow H + \gamma)$
- Deuterium: $p + n \rightarrow D + \gamma$

$H, {}^2H, {}^3H, {}^3He, {}^4He, {}^7Be, {}^7Li$

- Hydrogen: $(p + e^- \rightarrow H + \gamma)$
- Deuterium: $p + n \rightarrow D + \gamma$
- lighter elements:
 - $D + n \rightarrow {}^3H + \gamma$
 - $D + D \rightarrow {}^3H + p$
 - $D + p \rightarrow {}^3He + \gamma$
 - $D + D \rightarrow {}^3He + n$
 - ${}^3He + n \rightarrow {}^3H + p$

H, ²H, ³H, ³He, ⁴He, ⁷Be, ⁷Li

- Hydrogen: $(p + e^- \rightarrow H + \gamma)$
- Deuterium: $p + n \rightarrow D + \gamma$
- lighter elements:
 - $D + n \rightarrow {}^3\text{H} + \gamma$
 - $D + D \rightarrow {}^3\text{H} + p$
 - $D + p \rightarrow {}^3\text{He} + \gamma$
 - $D + D \rightarrow {}^3\text{He} + n$
 - ${}^3\text{He} + n \rightarrow {}^3\text{H} + p$
- ⁴He:
 - $D + D \rightarrow {}^4\text{He} + \gamma$
 - $D + {}^3\text{He} \rightarrow {}^4\text{He} + p$
 - $D + {}^3\text{H} \rightarrow {}^4\text{He} + n$
 - ${}^3\text{He} + {}^3\text{He} \rightarrow {}^4\text{He} + 2p$
 - ${}^3\text{He} + n \rightarrow {}^4\text{He} + \gamma$
 - ${}^3\text{H} + p \rightarrow {}^4\text{He} + \gamma$

H, ²H, ³H, ³He, ⁴He, ⁷Be, ⁷Li

$$E_{4\text{He}} \approx 28.3\text{MeV}$$

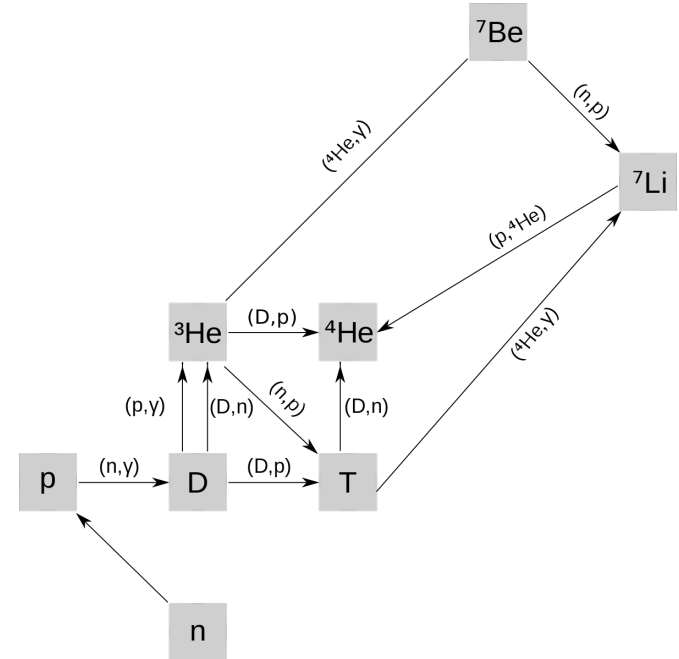
⇒ **safe from photo-dissociation**

H, ^2H , ^3H , ^3He , ^4He , ^7Be , ^7Li

- Hydrogen: $(p + e^- \rightarrow \text{H} + \gamma)$
- Deuterium: $p + n \rightarrow \text{D} + \gamma$
- lighter elements:
 - $\text{D} + n \rightarrow ^3\text{H} + \gamma$
 - $\text{D} + \text{D} \rightarrow ^3\text{H} + p$
 - $\text{D} + p \rightarrow ^3\text{He} + \gamma$
 - $\text{D} + \text{D} \rightarrow ^3\text{He} + n$
 - $^3\text{He} + n \rightarrow ^3\text{H} + p$
- ^4He :
 - $\text{D} + \text{D} \rightarrow ^4\text{He} + \gamma$
 - $\text{D} + ^3\text{He} \rightarrow ^4\text{He} + p$
 - $\text{D} + ^3\text{H} \rightarrow ^4\text{He} + n$
 - $^3\text{He} + ^3\text{He} \rightarrow ^4\text{He} + 2p$
 - $^3\text{He} + n \rightarrow ^4\text{He} + \gamma$
 - $^3\text{H} + p \rightarrow ^4\text{He} + \gamma$
- ^7Be , ^7Li :
 - $^3\text{He} + ^4\text{He} \rightarrow ^7\text{Be} + \gamma$
 - $^7\text{Be} \rightarrow ^7\text{Li} + e^+ + \nu_e$
 - $^3\text{H} + ^4\text{He} \rightarrow ^7\text{Li} + e^+ + \nu_e$

H, ²H, ³H, ³He, ⁴He, ⁷Be, ⁷Li

- Hydrogen: $(p + e^- \rightarrow H + \gamma)$
- Deuterium: $p + n \rightarrow D + \gamma$
- lighter elements:
 - $D + n \rightarrow {}^3H + \gamma$
 - $D + D \rightarrow {}^3H + p$
 - $D + p \rightarrow {}^3He + \gamma$
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 - ${}^3He + n \rightarrow {}^3H + p$
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 - $D + D \rightarrow {}^4He + \gamma$
 - $D + {}^3He \rightarrow {}^4He + p$
 - $D + {}^3H \rightarrow {}^4He + n$
 - ${}^3He + {}^3He \rightarrow {}^4He + 2p$
 - ${}^3He + n \rightarrow {}^4He + \gamma$
 - ${}^3H + p \rightarrow {}^4He + \gamma$
- ⁷Be, ⁷Li:
 - ${}^3He + {}^4He \rightarrow {}^7Be + \gamma$
 - ${}^7Be \rightarrow {}^7Li + e^+ + \nu_e$
 - ${}^3H + {}^4He \rightarrow {}^7Li + e^+ + \nu_e$

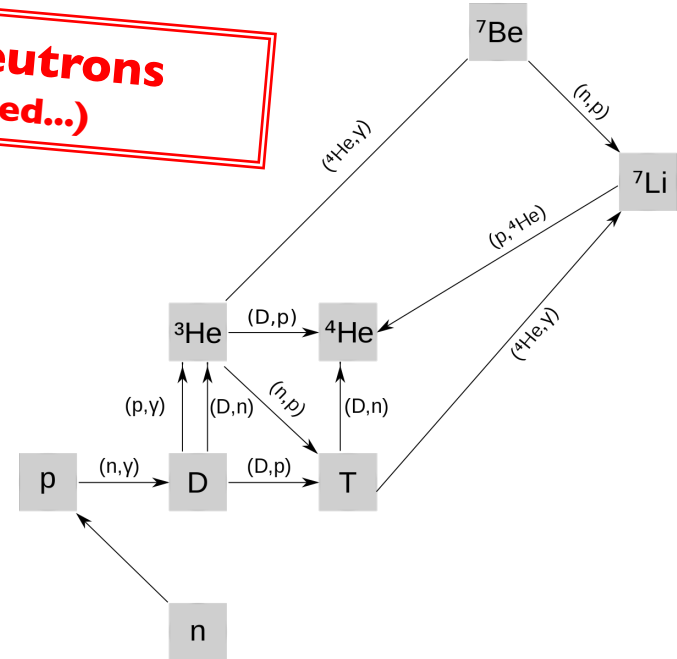


H, ²H, ³H, ³He, ⁴He, ⁷Be, ⁷Li

- Hydrogen: $(p + e^- \rightarrow H + \gamma)$
- Deuterium: $(p + n \rightarrow D + \gamma)$
- lighter elements: $(D + D \rightarrow {}^3H + \gamma)$

**we require free protons and neutrons
(weak-interaction freeze-out revisited...)**

- ⁴He:
 - $D + D \rightarrow {}^4He + \gamma$
 - $D + {}^3He \rightarrow {}^4He + p$
 - $D + {}^3H \rightarrow {}^4He + n$
 - ${}^3He + {}^3He \rightarrow {}^4He + 2p$
 - ${}^3He + n \rightarrow {}^4He + \gamma$
 - ${}^3H + p \rightarrow {}^4He + \gamma$



- ⁷Be, ⁷Li:
 - ${}^3He + {}^4He \rightarrow {}^7Be + \gamma$
 - ${}^7Be \rightarrow {}^7Li + e^+ + \nu_e$
 - ${}^3H + {}^4He \rightarrow {}^7Li + e^+ + \nu_e$

gap at A=8 prohibits production of heavier isotopes!

- introduction
- particle physics
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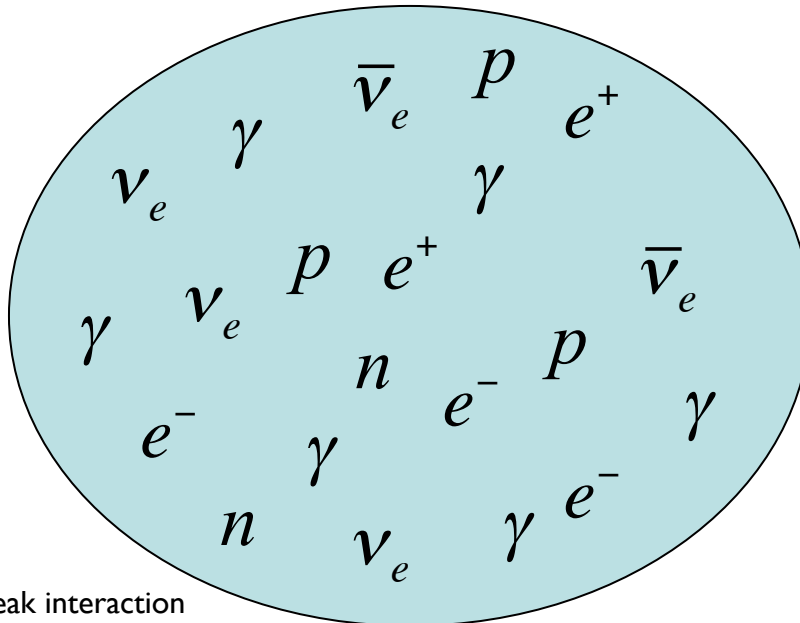
- thermal equilibrium of neutrons & protons:

$$\nu_e + n \rightleftharpoons p + e^-$$

$$e^+ + n \rightleftharpoons p + \bar{\nu}_e$$

$$n \rightleftharpoons p + e^- + \bar{\nu}_e$$

} weak interaction*



*only ν_e contributes to weak interaction

- thermal equilibrium of neutrons & protons:

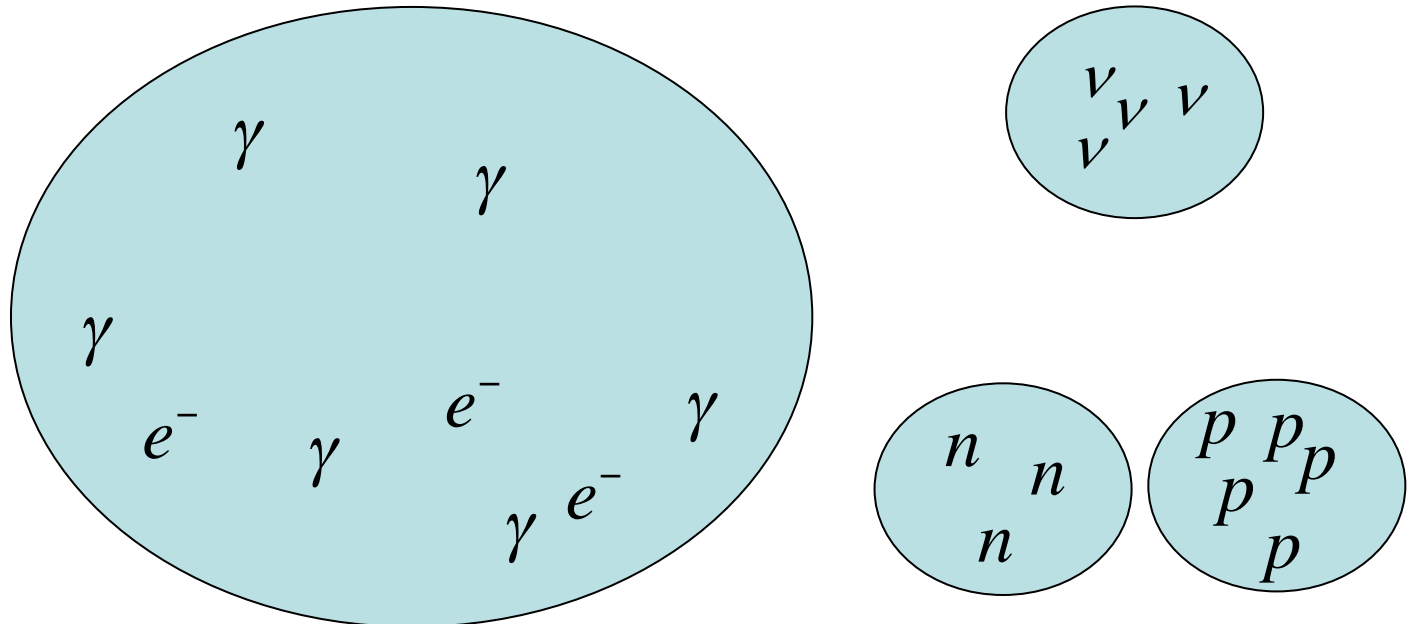
$$\nu_e + n \rightleftharpoons p + e^-$$

$$e^+ + n \rightleftharpoons p + \bar{\nu}_e$$

$$n \rightleftharpoons p + e^- + \bar{\nu}_e$$

} weak interaction

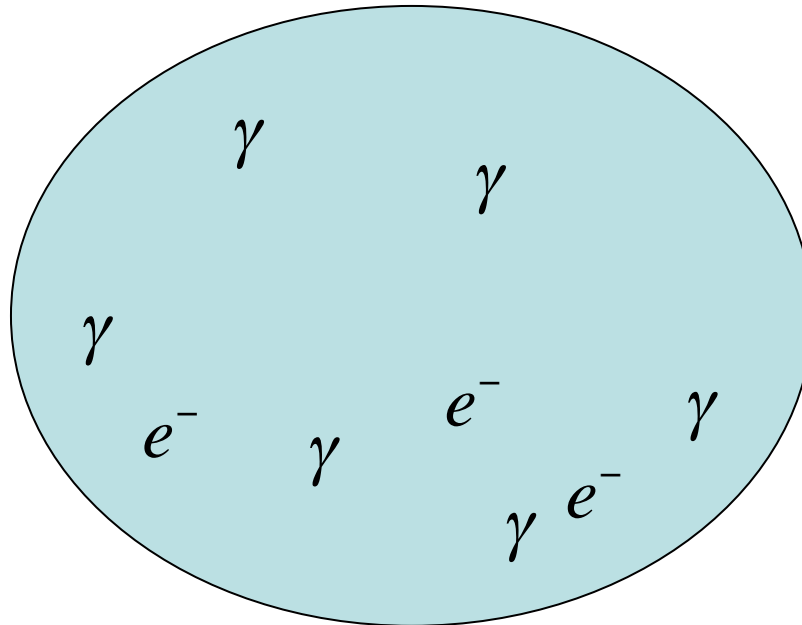
- weak interaction freezes out at $T \approx 0.8$ MeV



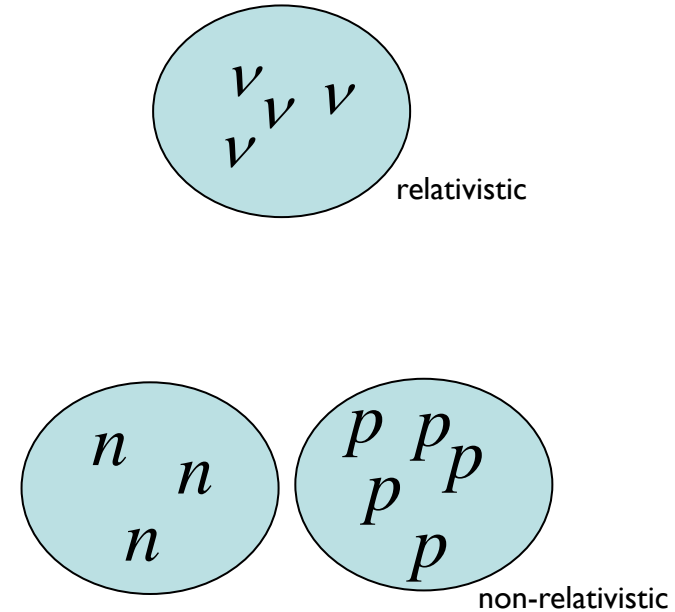
▪ inventory ($T < 0.8\text{MeV}$):

- relativistic particles in equilibrium: e^-, e^+
- decoupled relativistic particles: ν
- decoupled non-relativistic particles: n, p

thermal bath



decoupled species



▪ inventory ($T < 0.8\text{MeV}$):

- relativistic particles in equilibrium:
- decoupled relativistic particles:
- decoupled non-relativistic particles:

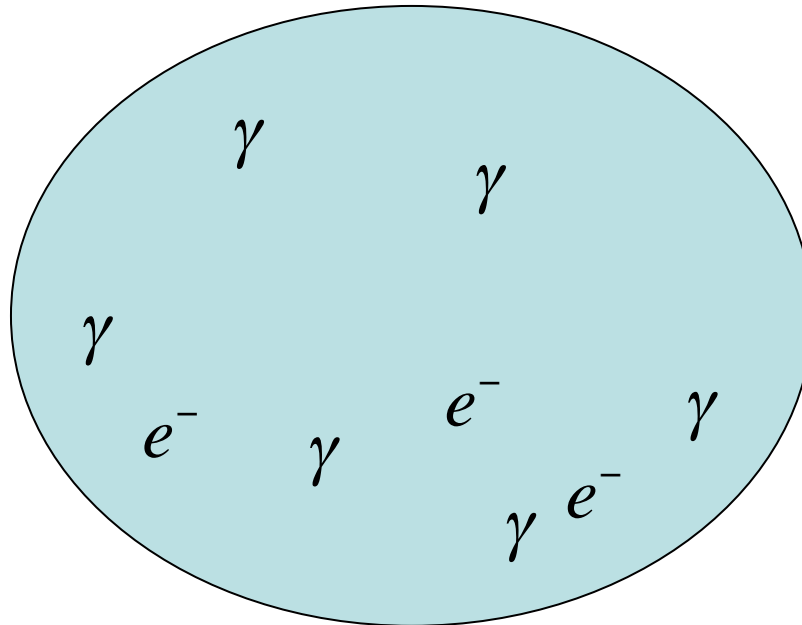
e^-, e^+

ν

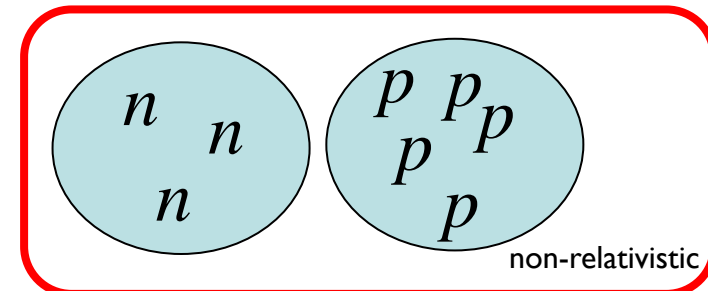
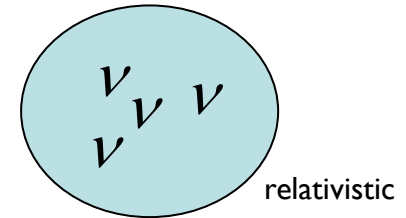
n, p

building blocks for nuclei!

thermal bath



decoupled species



▪ inventory ($T < 0.8\text{MeV}$):

- relativistic particles in equilibrium:
- decoupled relativistic particles:
- decoupled non-relativistic particles:

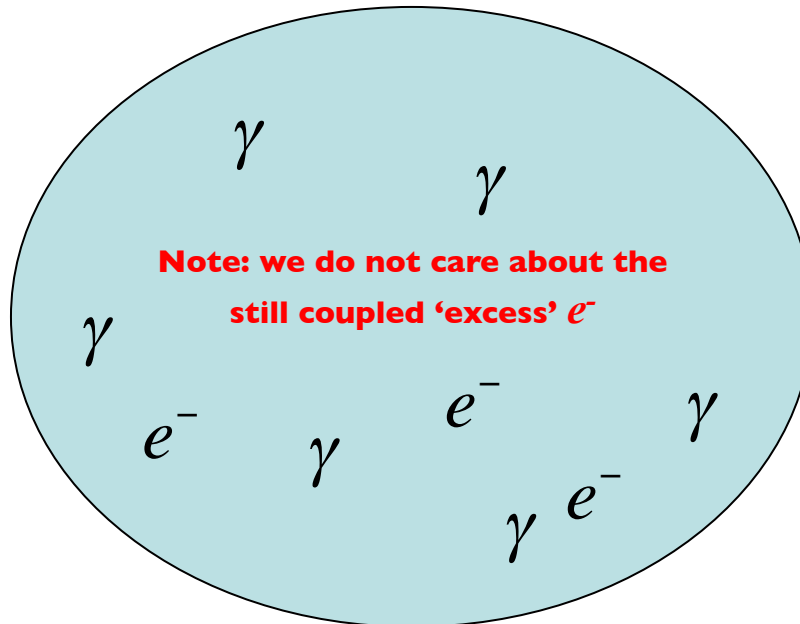
e^-, e^+

ν

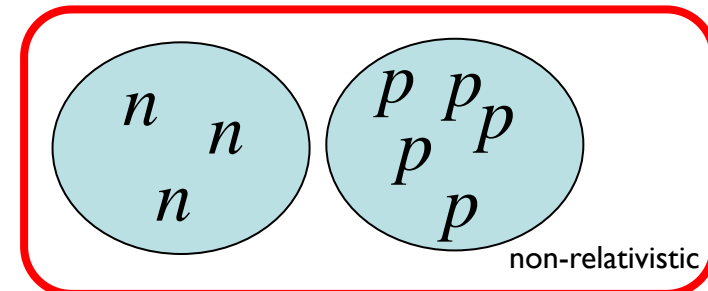
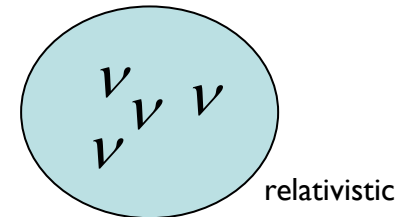
n, p

building blocks for nuclei!

thermal bath



decoupled species



▪ inventory ($T < 0.8\text{MeV}$):

- relativistic particles in equilibrium:

e^-, e^+

- decoupled relativistic particles:

ν

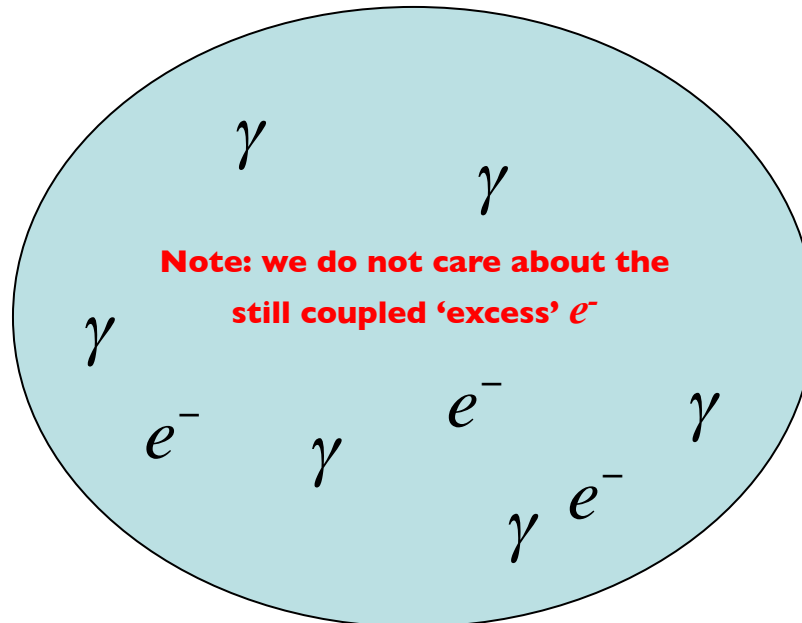
- decoupled non-relativistic particles:

n, p

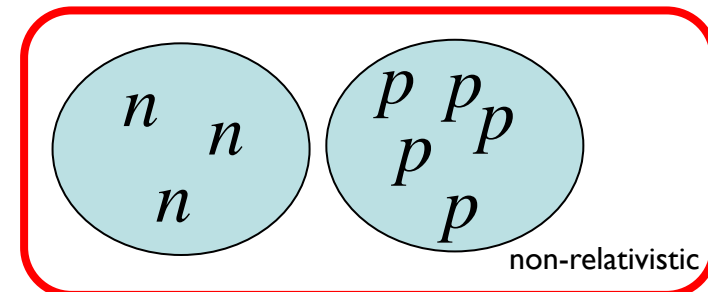
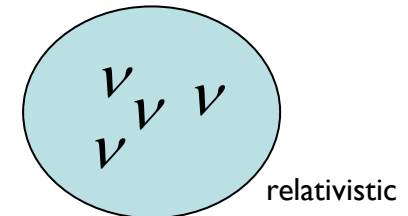
building blocks for nuclei!

let's calculate the number densities of n and p ...

thermal bath



decoupled species

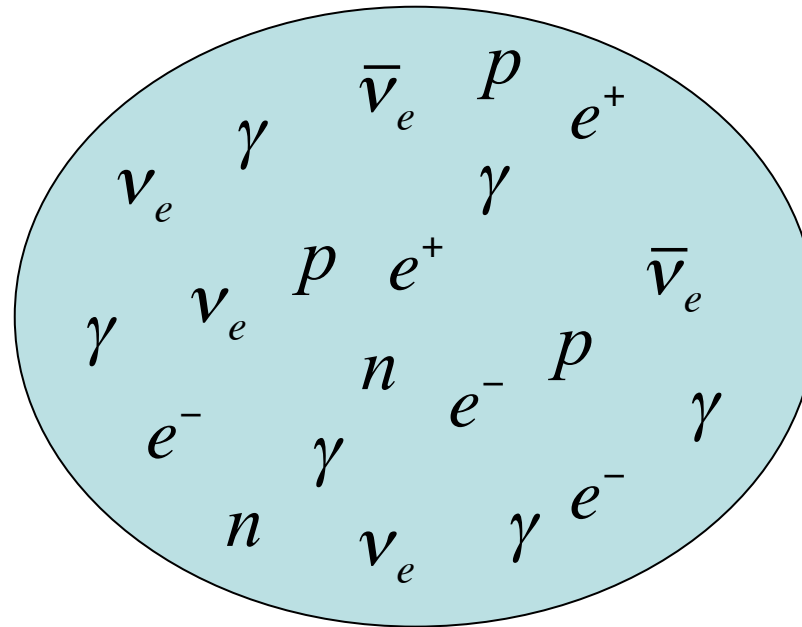


- neutron-to-proton ratio:*

the abundance of neutrons determines
how many nuclei beyond ($A=1$) can be formed...

*we will measure neutron abundance relative to protons...

- neutron-to-proton ratio:



thermal bath

$$\nu_e + n \Leftrightarrow p + e^-$$

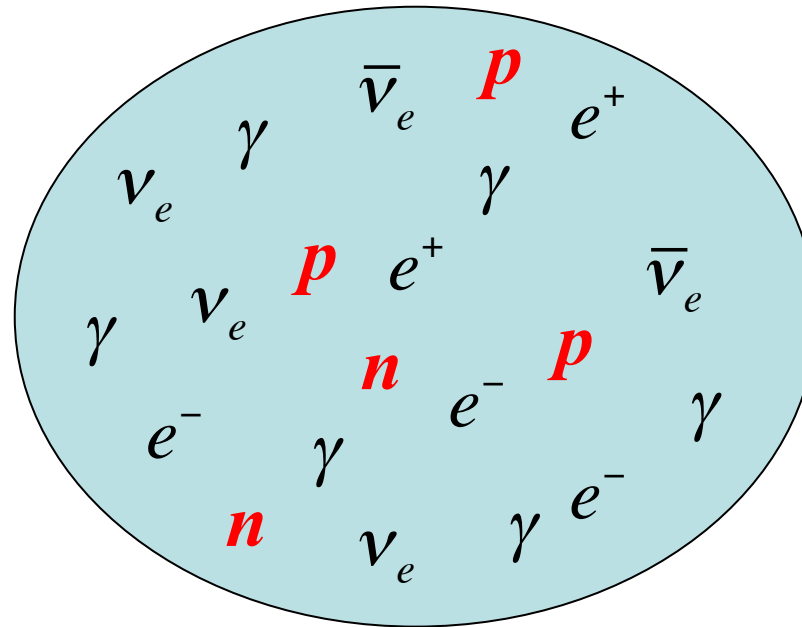
$$e^+ + n \Leftrightarrow p + \bar{\nu}_e$$

$$n \Leftrightarrow p + e^- + \bar{\nu}_e$$

$$e^- + \gamma \Leftrightarrow e^- + \gamma$$

- neutron-to-proton ratio:

$$n_A = ?$$



thermal bath

$$\nu_e + n \Leftrightarrow p + e^-$$

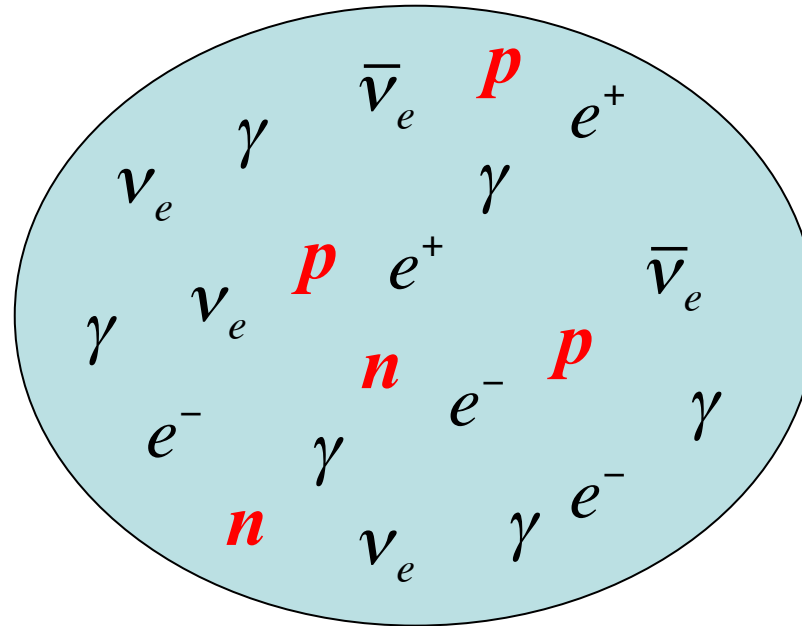
$$e^+ + n \Leftrightarrow p + \bar{\nu}_e$$

$$n \Leftrightarrow p + e^- + \bar{\nu}_e$$

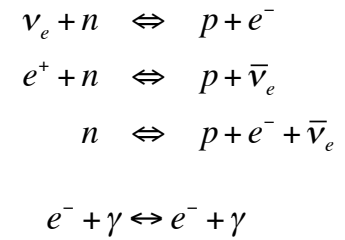
$$e^- + \gamma \Leftrightarrow e^- + \gamma$$

- neutron-to-proton ratio:

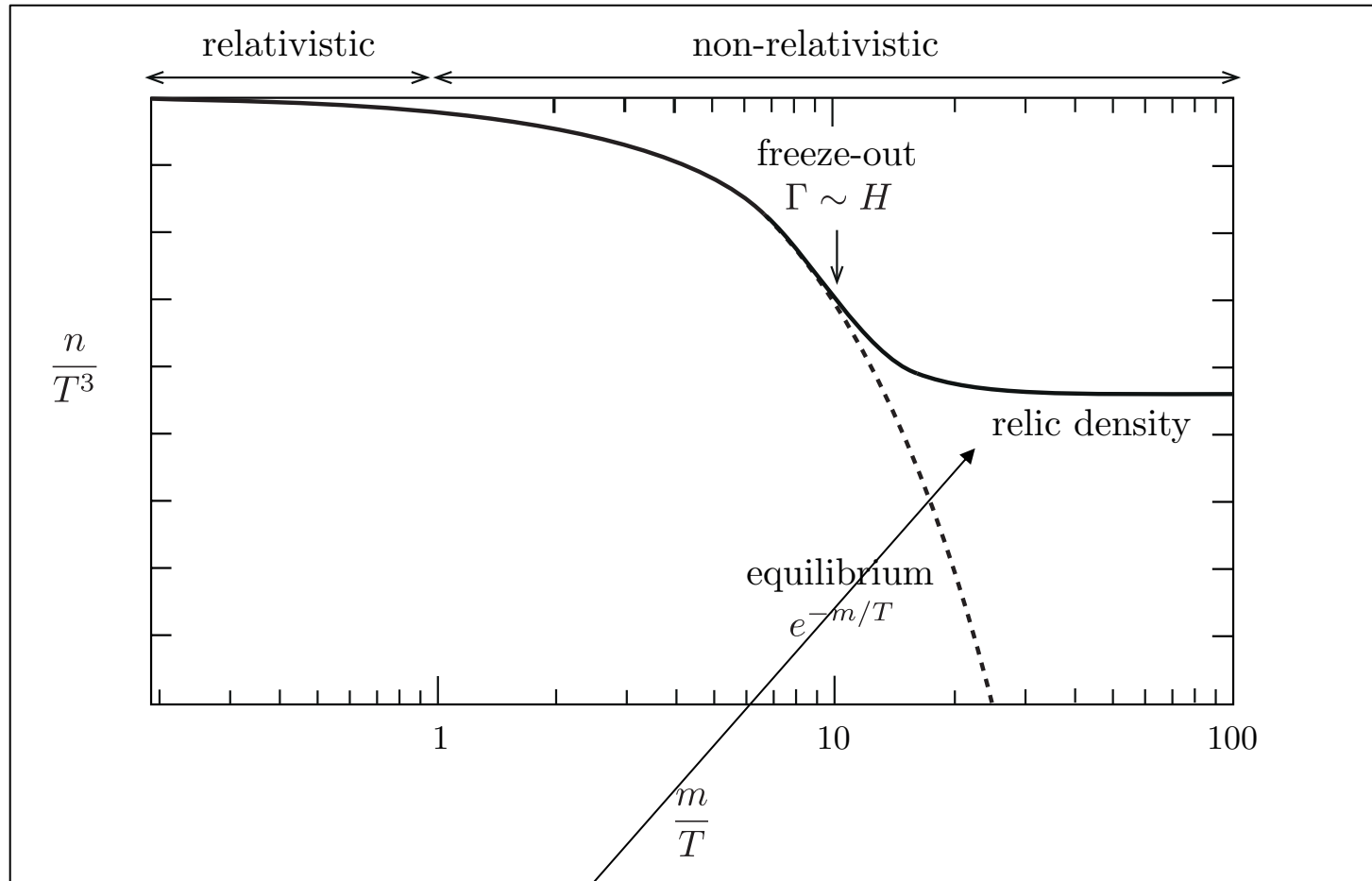
$$n_A = g_A \left(\frac{m_A kT}{2\pi\hbar^2} \right)^{3/2} e^{-(m_A - \mu_A)c^2/kT} \quad (\text{non-relativistic particles})$$



thermal bath



- neutron-to-proton ratio:



we start before the freeze-out of weak interaction,
as the equilibrium distributions will be 'frozen' in at freeze-out

- neutron-to-proton ratio:

$$n_A = g_A \left(\frac{m_A kT}{2\pi\hbar^2} \right)^{3/2} e^{-(m_A - \mu_A)c^2/kT} \quad (\text{non-relativistic particles})$$

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$$\frac{n_n}{n_p} = \left(\frac{m_n}{m_p} \right)^{3/2} e^{-Q/kT} e^{(\mu_n - \mu_p)/kT} \quad (Q=1.293 \text{ MeV})$$

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what about μ_n and μ_p ?

(for relativistic species we could neglect μ , but not so for non-relativistic)

■ neutron-to-proton ratio:

$$n_A = g_A \left(\frac{m_A kT}{2\pi\hbar^2} \right)^{3/2} e^{-(m_A - \mu_A)c^2/kT} \quad (\text{non-relativistic particles})$$

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■ chemical potential:

- energy absorbed or released during chemical reaction

- chemical equilibrium*: $A + B \leftrightarrow C + D \xrightarrow{\text{chemical equilibrium}} \mu_A + \mu_B = \mu_C + \mu_D$

*nuclear reactions are faster than cosmic expansion...

- neutron-to-proton ratio:

$$n_A = g_A \left(\frac{m_A kT}{2\pi\hbar^2} \right)^{3/2} e^{-(m_A - \mu_A)c^2/kT} \quad (\text{non-relativistic particles})$$

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- chemical potential:

- energy absorbed or released during chemical reaction

- chemical equilibrium: $A + B \leftrightarrow C + D \xrightarrow{\text{chemical equilibrium}} \mu_A + \mu_B = \mu_C + \mu_D$

what reaction is going on?

- neutron-to-proton ratio:

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what about μ_n and μ_p ?

- β -decay:

$$\nu_e + n \leftrightarrow p + e^- \quad \rightarrow \mu_{\nu_e} + \mu_n = \mu_p + \mu_e$$

- neutron-to-proton ratio:

$$n_A = g_A \left(\frac{m_A kT}{2\pi\hbar^2} \right)^{3/2} e^{-(m_A - \mu_A)c^2/kT} \quad (\text{non-relativistic particles})$$

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the Universe is neutral $\Rightarrow n_{e^-} - n_{e^+} = n_p$

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what about μ_{ν_e} and μ_e ?

the Universe is neutral $\Rightarrow n_{e^-} - n_{e^+} = n_p$

$$n_{e^-} - n_{e^+} = \frac{2T^3}{6\pi^2} \left(\pi^2 \left(\frac{\mu_e}{T} \right) + \left(\frac{\mu_e}{T} \right)^3 \right)$$

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expression using
 η = baryon-to-photon ratio

$$n_p \approx \eta n_\gamma = \eta \frac{2\zeta(3)}{\pi^2} T^3$$

- neutron-to-proton ratio:

$$n_A = g_A \left(\frac{m_A kT}{2\pi\hbar^2} \right)^{3/2} e^{-(m_A - \mu_A)c^2/kT} \quad (\text{non-relativistic particles})$$

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what about μ_{ν_e} and μ_e ?

the Universe is neutral $\Rightarrow n_{e^-} - n_{e^+} = n_p$

$$\frac{1}{6} \left(\pi^2 \left(\frac{\mu_e}{T} \right) + \left(\frac{\mu_e}{T} \right)^3 \right) = \eta \zeta(3)$$

- neutron-to-proton ratio:

$$n_A = g_A \left(\frac{m_A kT}{2\pi\hbar^2} \right)^{3/2} e^{-(m_A - \mu_A)c^2/kT} \quad (\text{non-relativistic particles})$$

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the Universe is neutral $\Rightarrow n_{e^-} - n_{e^+} = n_p$

$$\left(\frac{\mu_e}{T} \right) + \left(\frac{\mu_e}{T} \right)^3 = \frac{6\zeta(3)}{\pi^2} \eta$$

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what about μ_{ν_e} and μ_e ?

the Universe is neutral $\Rightarrow n_{e^-} - n_{e^+} = n_p$

$$\Rightarrow \frac{\mu_e}{T} \approx \frac{6}{\pi^2} \zeta(3) \eta \xrightarrow{\eta \approx 10^{-9}} \mu_e \ll T$$

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$$n_A = g_A \left(\frac{m_A kT}{2\pi\hbar^2} \right)^{3/2} e^{-(m_A - \mu_A)c^2/kT} \quad (\text{non-relativistic particles})$$

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what about μ_{ν_e} ?

CMB measurements by WMAP/Planck $\Rightarrow \frac{\mu_{\nu_e}}{T} < 0.2$

- neutron-to-proton ratio:

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$$\frac{n_n}{n_p} = \left(\frac{m_n}{m_p} \right)^{3/2} e^{-Q/kT} e^{-\mu_{\nu_e}/kT} \quad (Q=1.293 \text{ MeV})$$

what about μ_{ν_e} ?

$$\text{CMB measurements by WMAP/Planck} \Rightarrow \frac{\mu_{\nu_e}}{T} < 0.2 \xrightarrow{\text{BBN}} \frac{\mu_{\nu_e}}{T} < 10^{-10}$$

- neutron-to-proton ratio:

$$n_A = g_A \left(\frac{m_A kT}{2\pi\hbar^2} \right)^{3/2} e^{-(m_A - \mu_A)c^2/kT} \quad (\text{non-relativistic particles})$$

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$$kT > 1.3 \text{ MeV} \quad \Rightarrow \quad n_n \approx n_p$$

$$kT < 1.3 \text{ MeV} \quad \Rightarrow \quad n_n < n_p$$

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- weak-interaction freezes out at $kT \leq 0.8 \text{ MeV} < Q$

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- weak-interaction freezes out at $kT \leq 0.8 \text{ MeV} < Q$

\Rightarrow we therefore end up with less neutrons than protons!

- neutron-to-proton ratio:

$$n_A = g_A \left(\frac{m_A kT}{2\pi\hbar^2} \right)^{3/2} e^{-(m_A - \mu_A)c^2/kT} \quad (\text{non-relativistic particles})$$

$$\frac{n_n}{n_p} = \left(\frac{m_n}{m_p} \right)^{3/2} e^{-Q/kT} \quad (Q=1.293 \text{ MeV})$$

$$kT \approx 0.72 \text{ MeV}^*$$

$$\Rightarrow$$

$$\frac{n_n}{n_p} \approx \frac{1}{6} \quad (\text{ratio of free neutrons \& protons!?!})$$

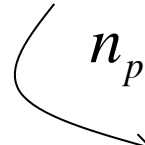
*Note, the weak interaction will not instantaneously freeze-out at 0.8MeV

- neutron-to-proton ratio:

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 $p + n \rightarrow D + \gamma$

BBN could now start ... or not?

- neutron-to-proton ratio:

$$\frac{n_n}{n_p} \approx \frac{1}{6} \quad (\text{ratio of free neutrons \& protons!?!})$$

\curvearrowright

$$p + n \rightarrow D + \gamma$$

Deuterium bottleneck:

D easily photo-dissociated by γ until $kT_D \cong 0.086\text{MeV}$

- neutron-to-proton ratio:

$$\frac{n_n}{n_p} \approx \frac{1}{6} \quad (\text{ratio of free neutrons \& protons!?!})$$

\curvearrowright

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Deuterium bottleneck:

D easily photo-dissociated by γ until $kT_D \cong 0.086\text{MeV}$,
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D easily photo-dissociated by γ until $kT_D \cong 0.086\text{MeV}$,
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something else happens in those 140s!

- neutron-to-proton ratio:

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\curvearrowright

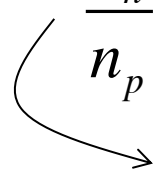
$$p + n \rightarrow D + \gamma$$

Deuterium bottleneck:

D easily photo-dissociated by γ until $kT_D \cong 0.086\text{MeV}$,
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→ neutrons have a limited lifetime $\tau_n \approx 887\text{ s}$

- neutron-to-proton ratio:

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→ neutrons have a limited lifetime $\tau_n \approx 887\text{ s}$

- ⇒ from $kT \approx 0.72\text{ MeV}$ until $kT_D \approx 0.086\text{ MeV}$ we have a race between...
- remaining free neutrons decaying away
 - remaining free neutrons being incorporated into nuclei (e.g. D)

- neutron-to-proton ratio:

$$\frac{n_n}{n_p} \approx \frac{1}{6} \quad (\text{ratio of free neutrons \& protons!?!})$$

$$\curvearrowright \quad \text{p} + \text{n} \rightarrow \text{D} + \gamma$$

Deuterium bottleneck:

D easily photo-dissociated by γ until $kT_D \approx 0.086 \text{ MeV}$,
and going from 0.72 MeV to 0.086 MeV takes about $t \approx 140 \text{ s}$

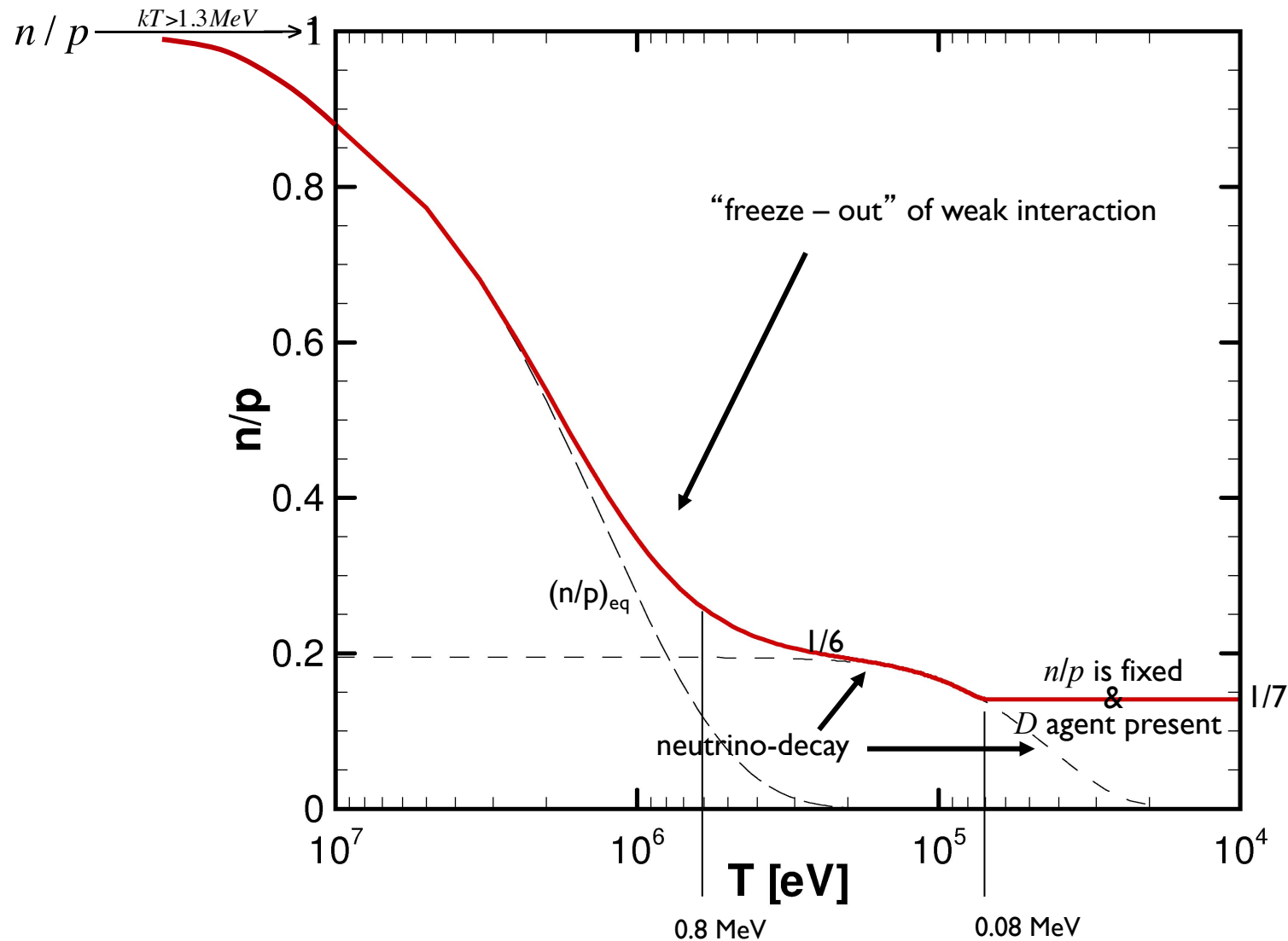
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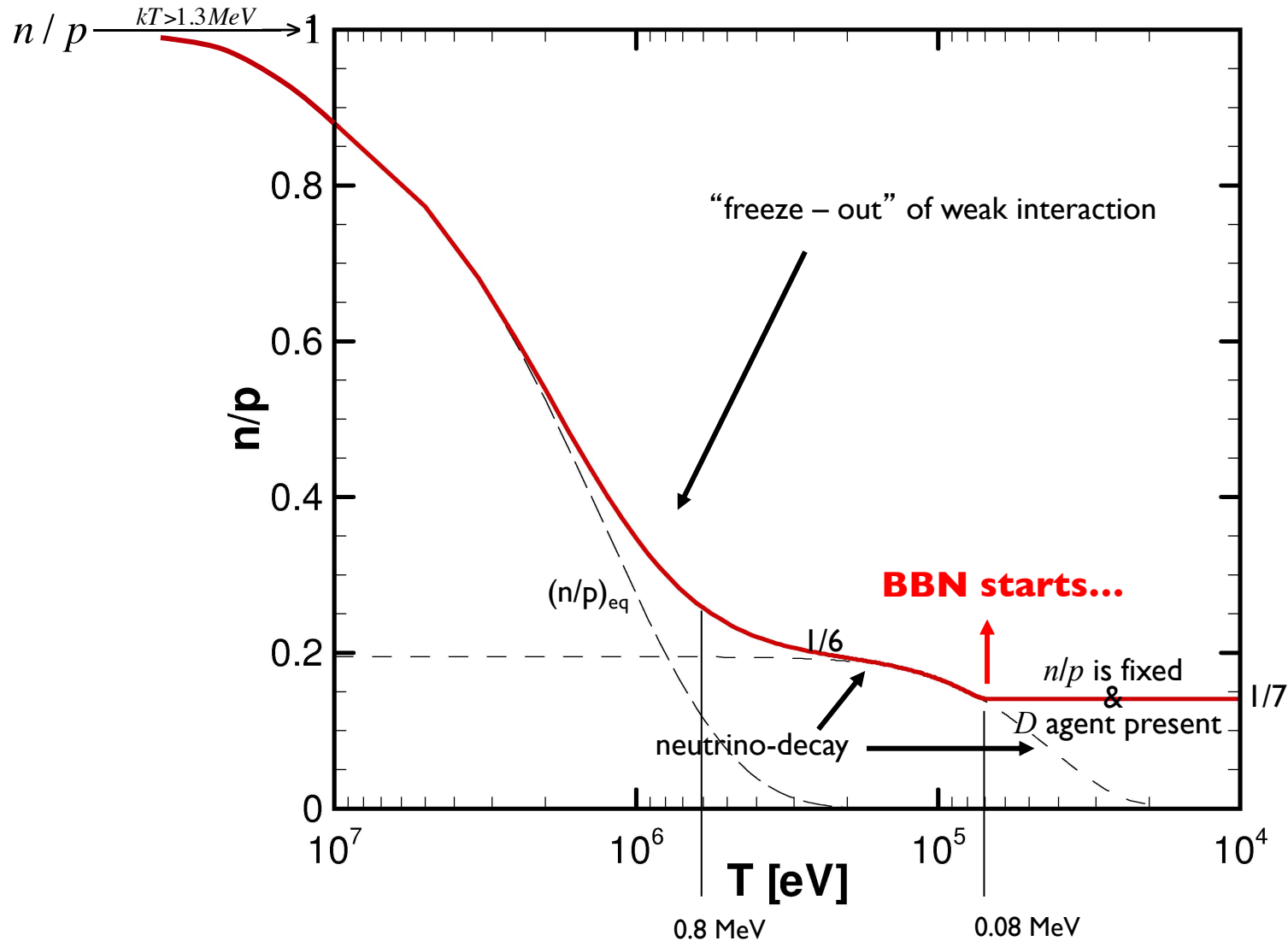
- remaining free neutrons decaying away
- remaining free neutrons being incorporated into nuclei (e.g. D)

$$\Rightarrow \frac{n_n}{n_p} = \frac{1}{6} e^{-t/\tau_n} \approx \frac{1}{7}$$

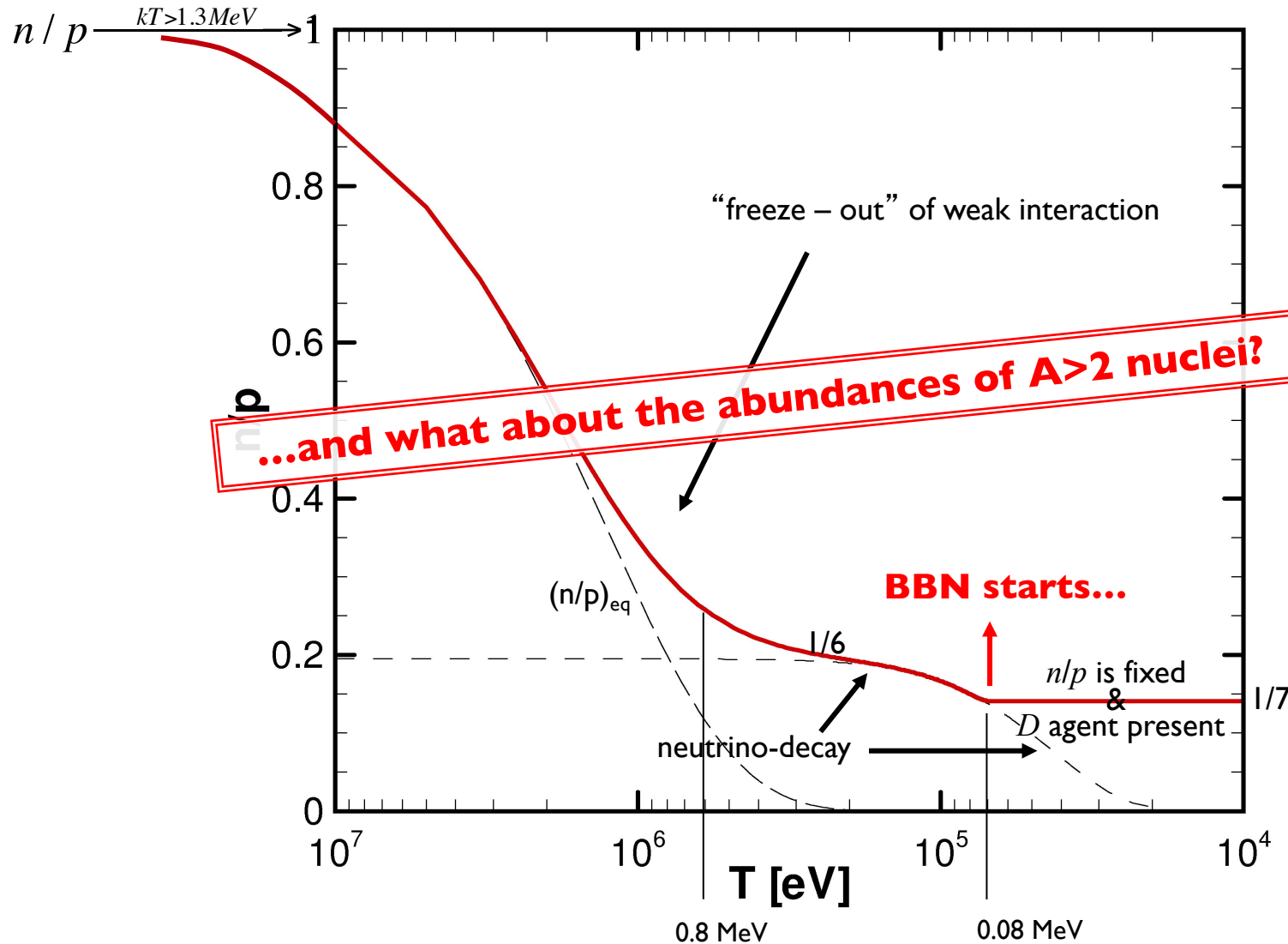
■ neutron-to-proton ratio:



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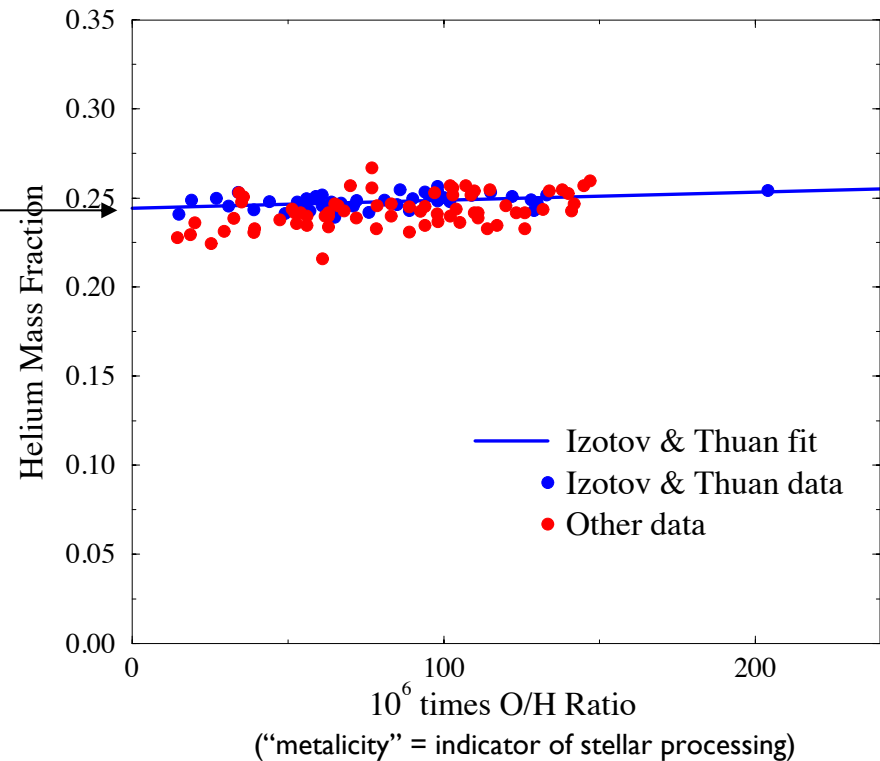
- neutron-to-proton ratio: primordial Helium abundance:

$$\Rightarrow X_{\text{He}} = Y = \frac{m_{\text{He}}}{m_{\text{He}} + m_{\text{H}}} = ?$$

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?



(Burles, Nollett & Turner, arxiv:9903300, Fig.4)

- neutron-to-proton ratio: primordial Helium abundance:

$$\frac{n_n}{n_p} = \frac{1}{7}$$

$$n_{He} = \frac{n_n}{2} \quad (\text{all neutrons end up in } {}^4\text{He}=2n+2p)$$

$$m_p \approx m_n$$

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$$m_p \approx m_n$$

$$\Rightarrow X_{He} = Y = \frac{4n_{He}}{4n_{He} + n_H} = \frac{2n_n}{2n_n + \underbrace{(n_p - n_n)}_{\text{protons not in } {}^4\text{He}}} = \frac{2n_n}{n_n + n_p} = \frac{2(n_n/n_p)}{1 + (n_n/n_p)} = \frac{2 \times 1/7}{1 + 1/7} = 0.25$$

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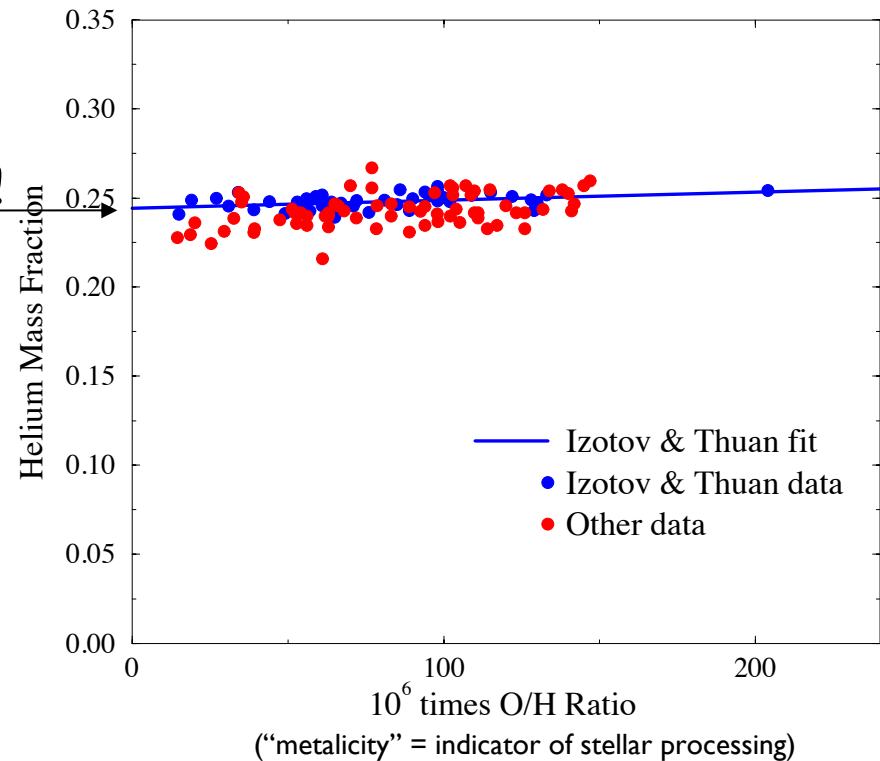
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$$\Rightarrow X_{He} = Y = \frac{m_{He}}{m_{He} + m_H} = 0.25 \quad \text{!}$$



(Burles, Nollett & Turner, 1999, Fig.4)

- neutron-to-proton ratio: primordial Helium abundance:

proper calculation leads to...

$$X_{He} = Y \approx 0.2454 + 0.0198(N_\nu - 3)$$

- nuclear statistical equilibrium: (non-relativistic nucleus $A = N_n + Z = \text{\#neutrons} + \text{\#protons}$)

- number density:
$$n_A = g_A \left(\frac{m_A kT}{2\pi\hbar^2} \right)^{3/2} e^{-(m_A - \mu_A)c^2/kT}$$

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→ relate to...

- proton number density n_p
- neutron number density n_n
- baryon-to-photon ratio η

- nuclear statistical equilibrium: (non-relativistic nucleus $A = N_n + Z = \text{\#neutrons} + \text{\#protons}$)

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- chemical reaction:
$$A \rightleftharpoons Z p + N_n n$$

- nuclear statistical equilibrium: (non-relativistic nucleus $A = N_n + Z = \text{\#neutrons} + \text{\#protons}$)

- number density:
$$n_A = g_A \left(\frac{m_A kT}{2\pi\hbar^2} \right)^{3/2} e^{-(m_A - \mu_A)c^2/kT}$$

- chemical equilibrium*:
$$\mu_A = Z\mu_p + N_n\mu_n = Z\mu_p + (A - Z)\mu_n$$

*nuclear reactions are faster than cosmic expansion

- nuclear statistical equilibrium: (non-relativistic nucleus $A = N_n + Z = \text{\#neutrons} + \text{\#protons}$)

- number density:
$$n_A = g_A \left(\frac{m_A kT}{2\pi\hbar^2} \right)^{3/2} e^{-(m_A - (Z\mu_p + (A-Z)\mu_n))c^2/kT}$$

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↓ ↓

eliminate μ_p and μ_n by...

- nuclear statistical equilibrium: (non-relativistic nucleus $A = N_n + Z = \text{\#neutrons} + \text{\#protons}$)

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\downarrow \downarrow

...introducing n_p and n_n

$$n_p = 2 \left(\frac{m_p kT}{2\pi\hbar^2} \right)^{3/2} e^{-(m_p - \mu_p) c^2 / kT}$$

$$n_n = 2 \left(\frac{m_n kT}{2\pi\hbar^2} \right)^{3/2} e^{-(m_n - \mu_n) c^2 / kT}$$

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$$n_A = g_A \left(\frac{m_A kT}{2\pi\hbar^2} \right)^{3/2} e^{-(m_A - (Z\mu_p + (A-Z)\mu_n))c^2/kT}$$

\downarrow \downarrow

...introducing n_p and n_n

$$e^{-\mu_p c^2/kT} = \frac{2}{n_p} \left(\frac{m_p kT}{2\pi\hbar^2} \right)^{3/2} e^{-m_p c^2/kT}$$

$$e^{-\mu_n c^2/kT} = \frac{2}{n_n} \left(\frac{m_n kT}{2\pi\hbar^2} \right)^{3/2} e^{-m_n c^2/kT}$$

⇔

$$n_p = 2 \left(\frac{m_p kT}{2\pi\hbar^2} \right)^{3/2} e^{-(m_p - \mu_p)c^2/kT}$$

$$n_n = 2 \left(\frac{m_n kT}{2\pi\hbar^2} \right)^{3/2} e^{-(m_n - \mu_n)c^2/kT}$$

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$$n_A = g_A \left(\frac{m_A kT}{2\pi\hbar^2} \right)^{3/2} e^{-(m_A - (Z\mu_p + (A-Z)\mu_n))c^2/kT}$$

...introducing n_p and n_n

$$\boxed{e^{-\mu_p c^2/kT}} = \frac{2}{n_p} \left(\frac{m_p kT}{2\pi\hbar^2} \right)^{3/2} e^{-m_p c^2/kT}$$

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⇔

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$$B_A = (Zm_p + (A - Z)m_n - m_A)c^2 \quad (\text{binding energy of nucleus})$$

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nucleus	${}^2\text{H}$	${}^3\text{H}$	${}^3\text{He}$	${}^4\text{He}$
B_A	2.2 MeV	8.48 MeV	7.72 MeV	28.3 MeV

(all above temperature of Universe at these times...)

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replace with baryon-to-photon ratio*...

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$\eta = n_b/n_\gamma$: baryon-to-photon ratio

$$n_\gamma = \frac{2\xi(3)}{\pi^2} \left(\frac{k_B}{\hbar c} \right)^3 T^3$$

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mass fraction reduced to...

- proton mass fraction X_p
- neutron mass fraction X_n
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$$\frac{dX_A}{dt} = -\Gamma_A \left(X_A - (1 - X_A) e^{-B_A/kT} \right)$$

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mass fraction reduced to...

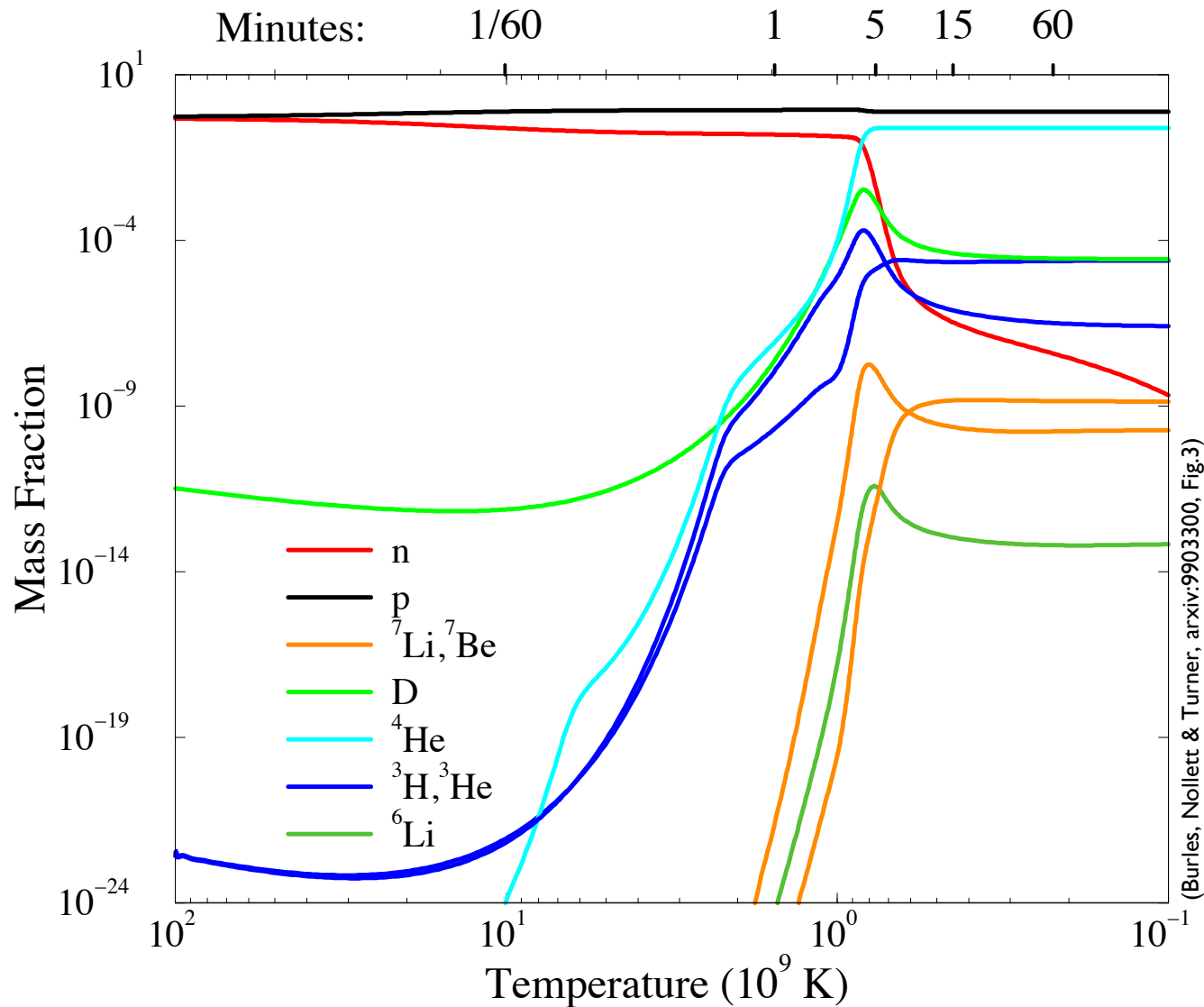
- proton mass fraction X_p
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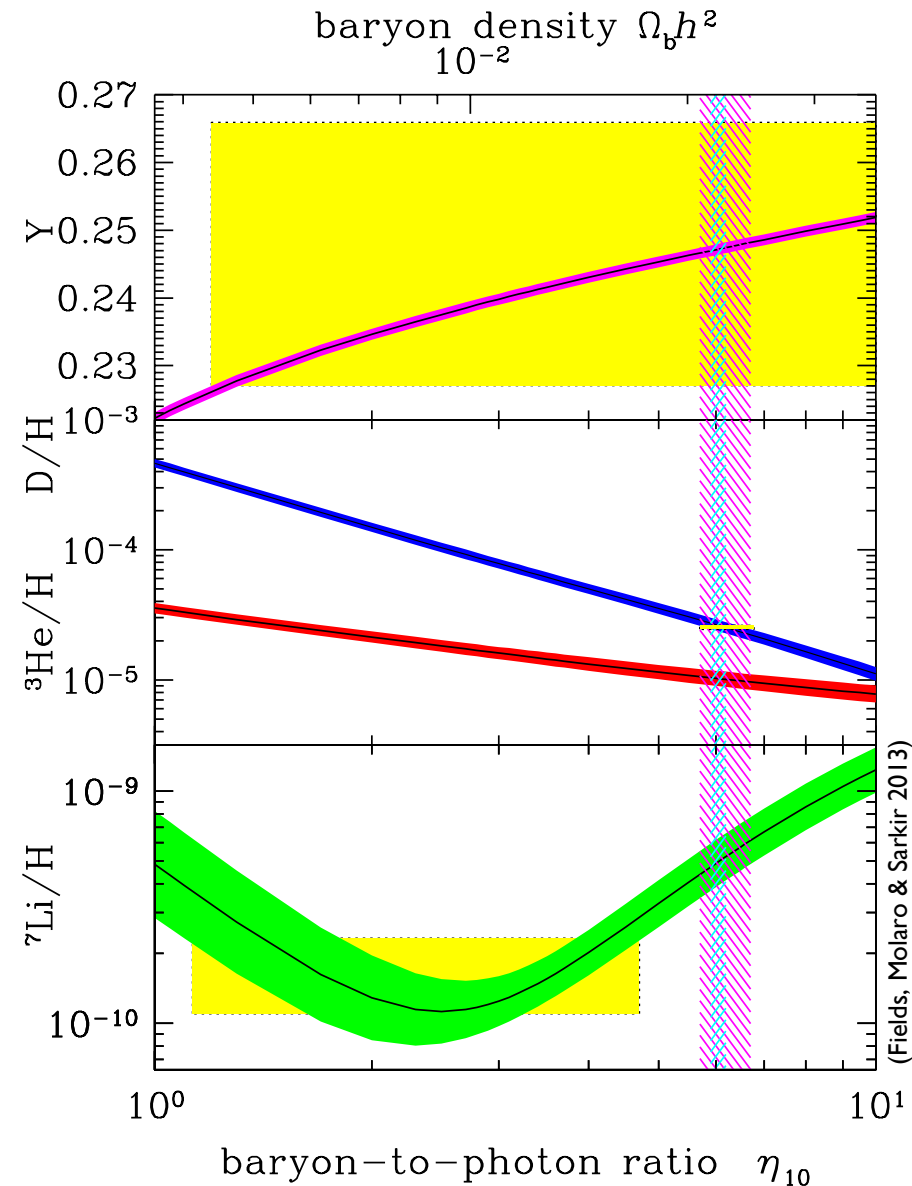
$$\frac{dX_A}{dt} = -\Gamma_A \left(X_A - (1 - X_A) e^{-B_A/kT} \right) \longrightarrow$$

- numerical calculations of cosmic evolution of mass fractions



dependence on baryon density:

- $\Omega_b \nearrow \rightarrow Y \nearrow$ (as BBN starts earlier)
- D and ${}^3\text{He}$ are primarily consumed
- $\Omega_b \nearrow \rightarrow {}^7\text{Be} \nearrow$
- $\Omega_b \nearrow \rightarrow {}^7\text{Li} \searrow \nearrow \nearrow$

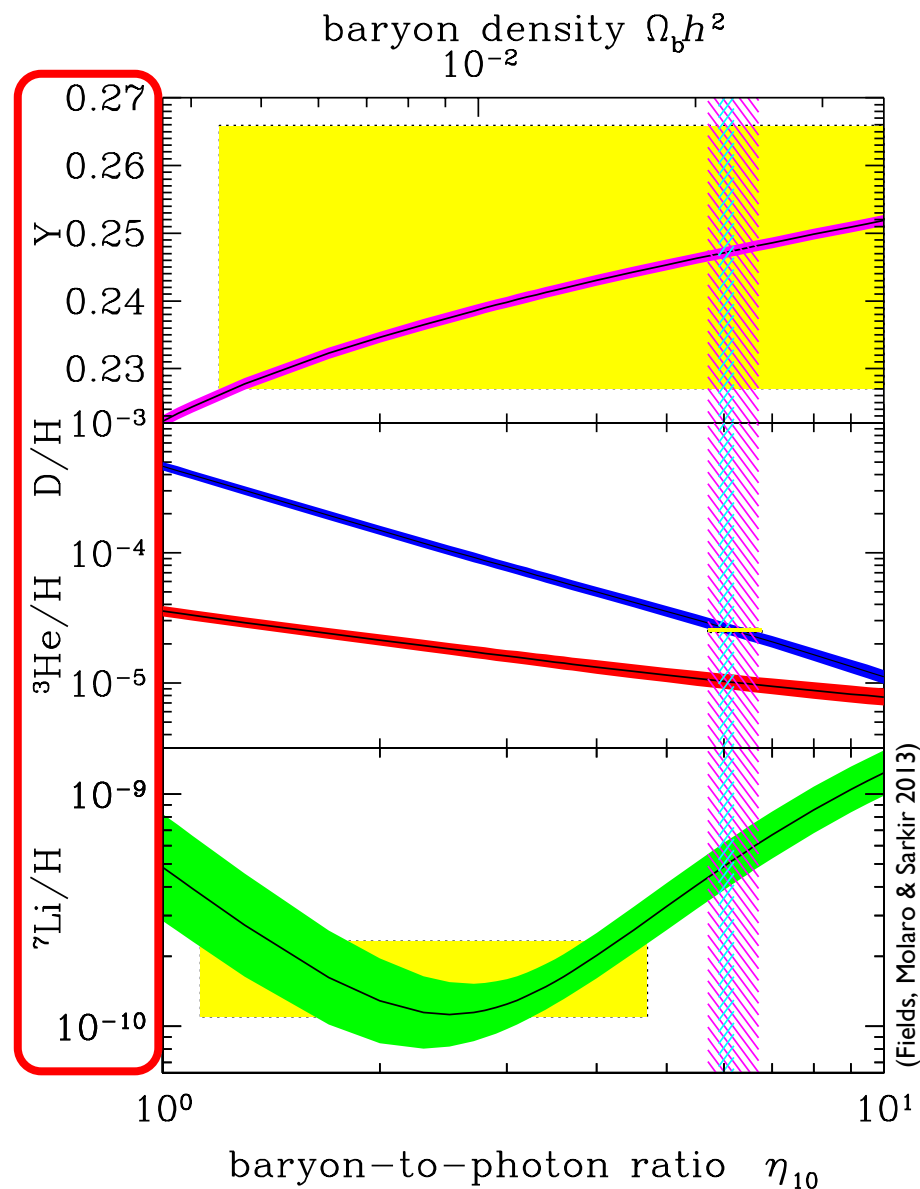


Remember: $Y = X_{4\text{He}}$

- dependence on baryon density:

9 orders of magnitude!

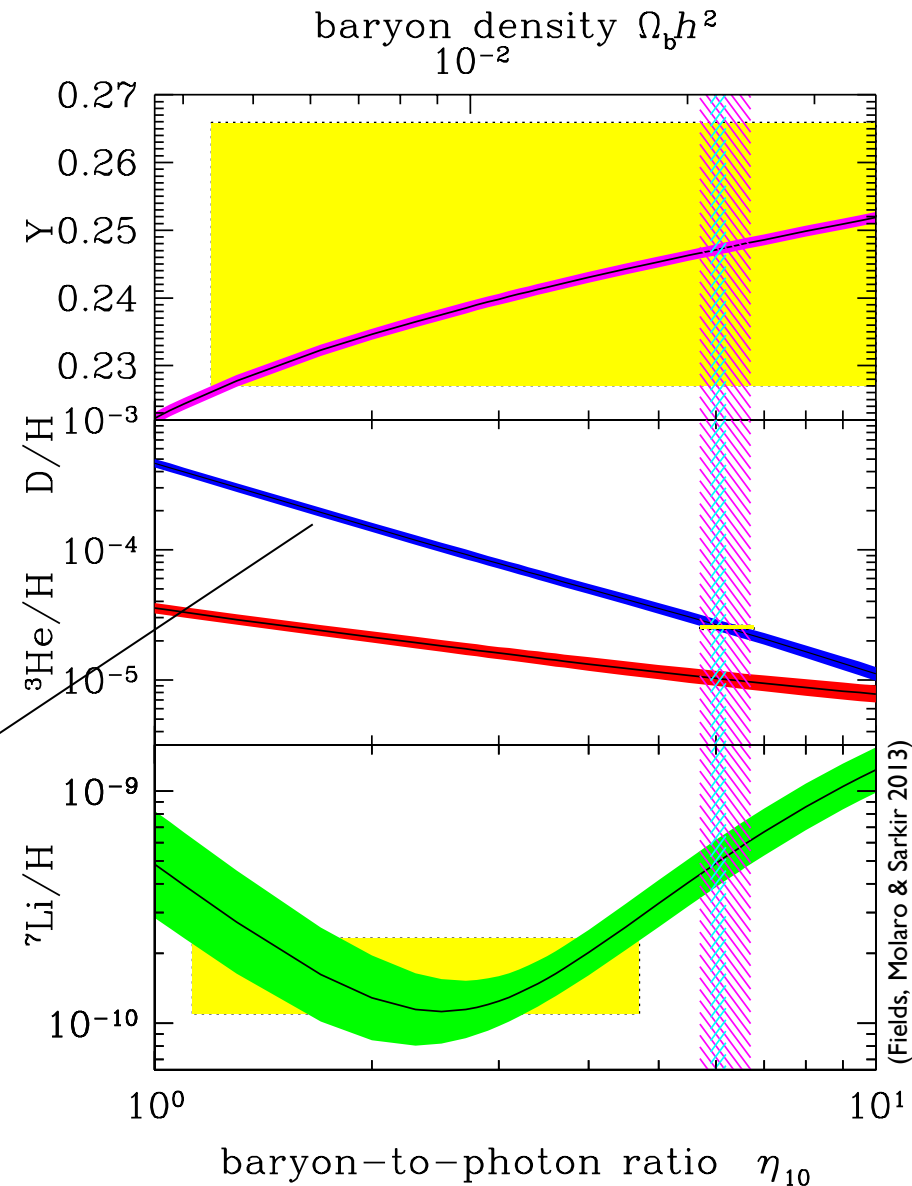
- $\Omega_b \uparrow \rightarrow Y \uparrow$ (as BBN starts earlier)
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dependence on baryon density:

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sensitive to baryon density!
("Baryometer")



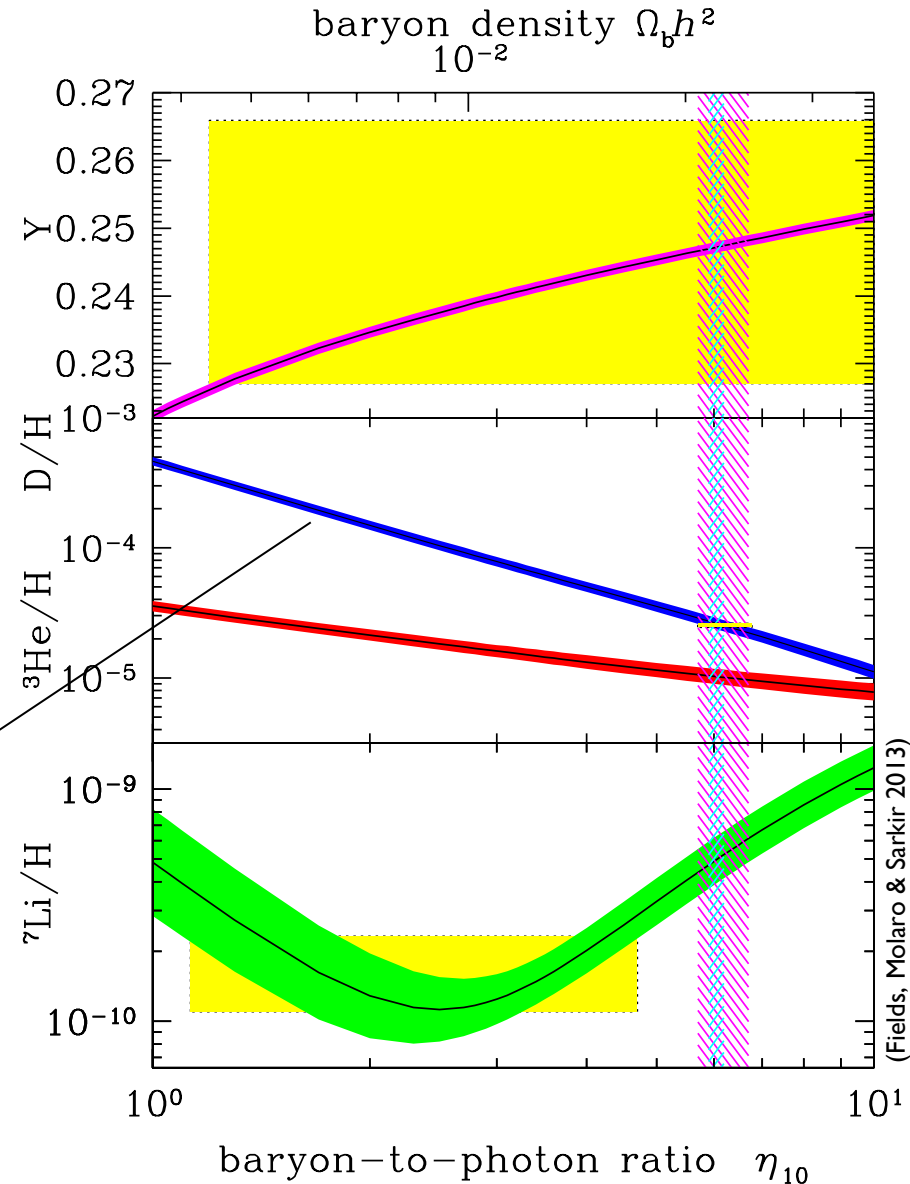
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sensitive to baryon density!
("Baryometer")

Note:

- D is only created during BBN \Rightarrow best Baryometer!
- D has highest sensitivity to Ω_b



dependence on baryon density:

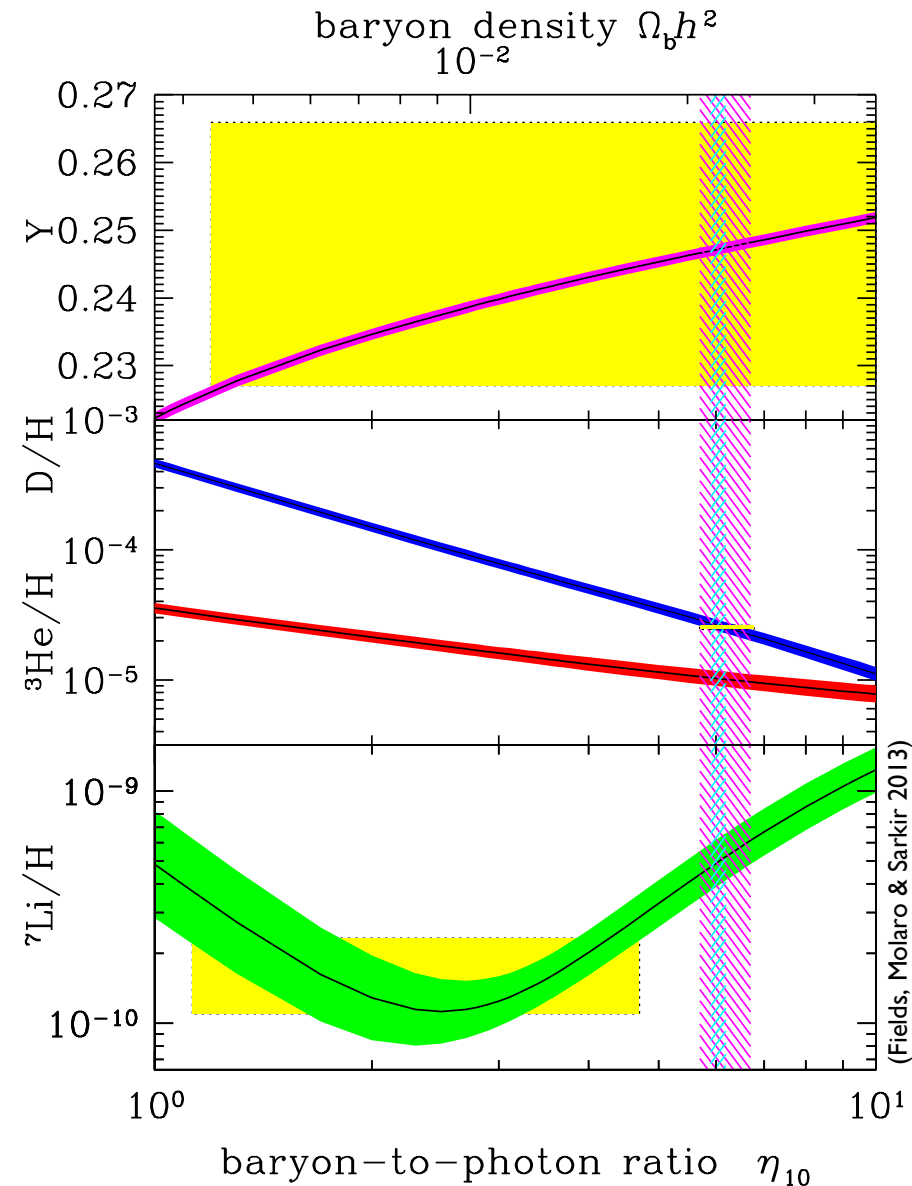
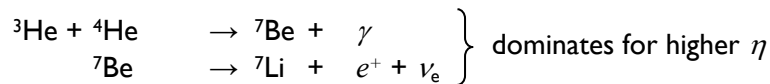
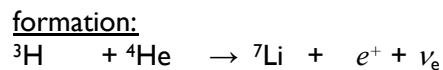
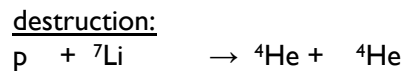
• $\Omega_b \nearrow \rightarrow Y \nearrow$ (as BBN starts earlier)

• D and ^3He are primarily consumed

• $\Omega_b \nearrow \rightarrow ^7\text{Be} \nearrow$

• $\Omega_b \nearrow \rightarrow ^7\text{Li} \searrow \nearrow$

^7Li shape:



▪ BBN depends on various ‘parameters’:

- g^* : $g^* \nearrow \rightarrow T_f \nearrow \rightarrow n/p \nearrow \rightarrow {}^4\text{He} \nearrow$
- τ_n : $\tau_n \nearrow \rightarrow T_f \nearrow \rightarrow n/p \nearrow \rightarrow {}^4\text{He} \nearrow$
- G_N : $G_N \nearrow \rightarrow T_f \nearrow \rightarrow n/p \nearrow \rightarrow {}^4\text{He} \nearrow$
- Q : $Q \nearrow \rightarrow n/p \searrow \rightarrow {}^4\text{He} \searrow$
- η : $\eta \nearrow \rightarrow X_A \nearrow \rightarrow {}^4\text{He} \nearrow \text{ \& D, T, } {}^3\text{He} \searrow$

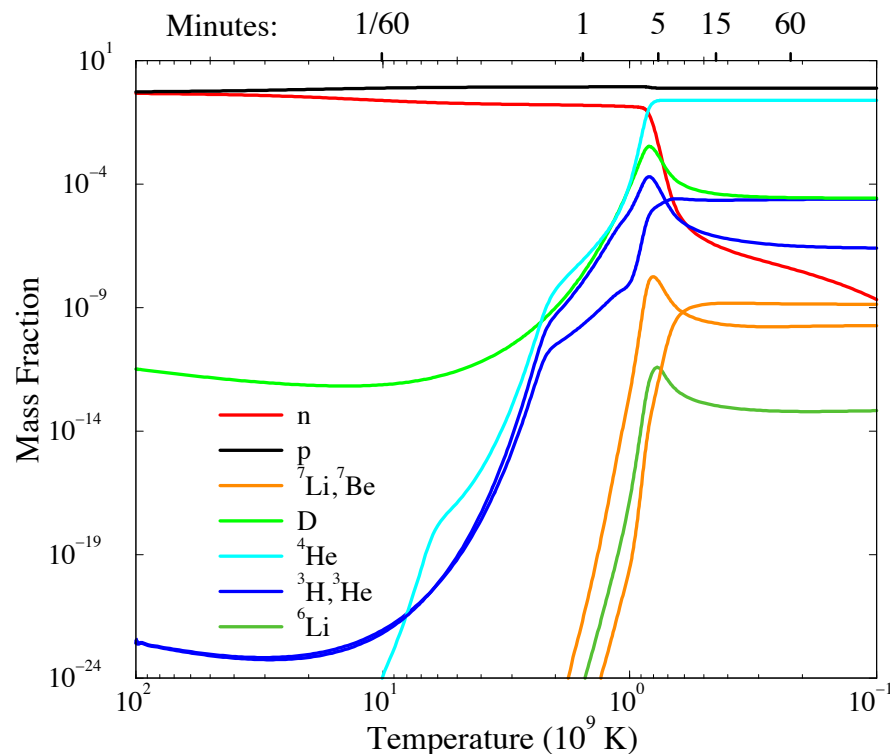
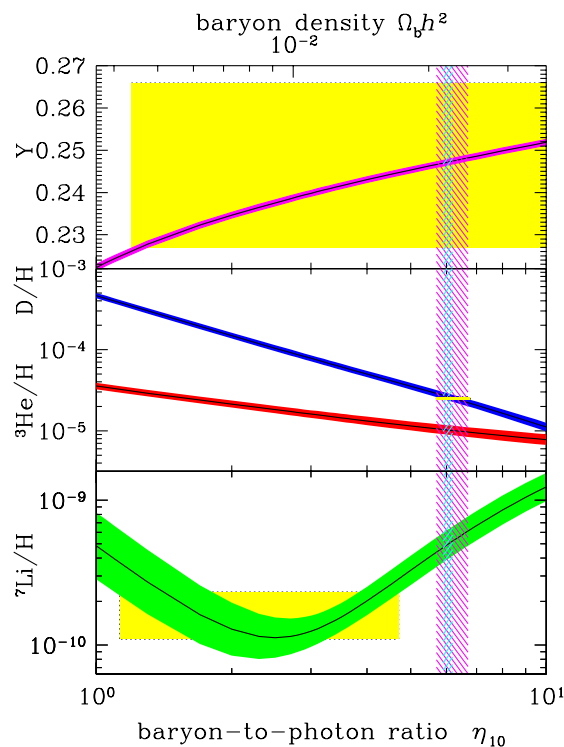
→ BBN is a probe of fundamental physics!

▪ BBN summary:

- production of $\text{H}, {}^2\text{H}, {}^3\text{H}, {}^3\text{He}, {}^4\text{He}, {}^7\text{Be}, {}^7\text{Li}$
- all other elements are in fact produced in stars
- mass fractions:

${}^2\text{H}$: deuterium D

${}^3\text{H}$: tritium T



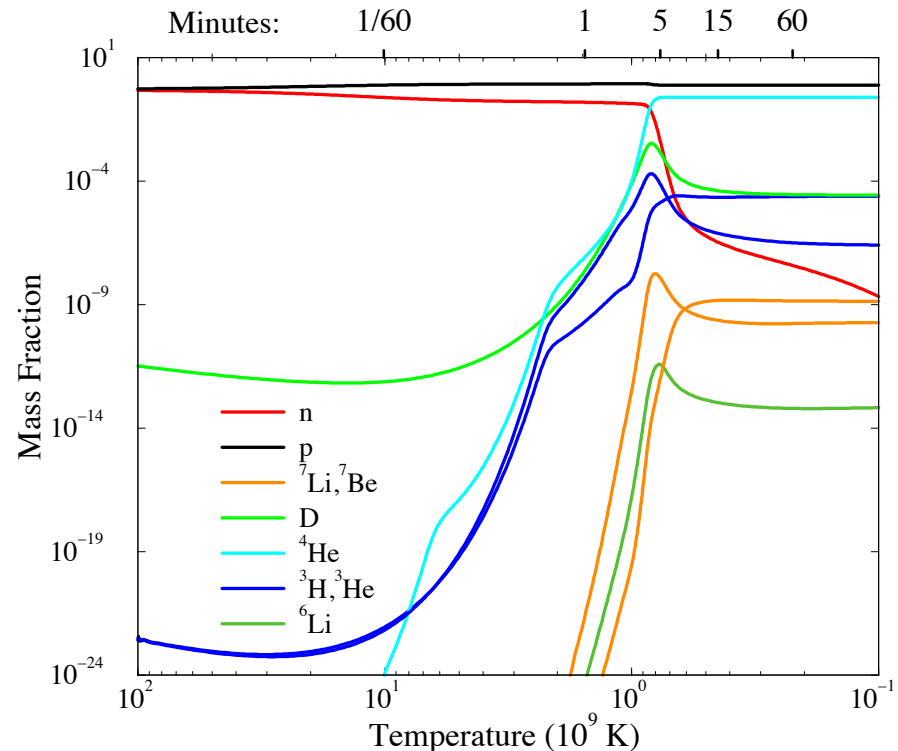
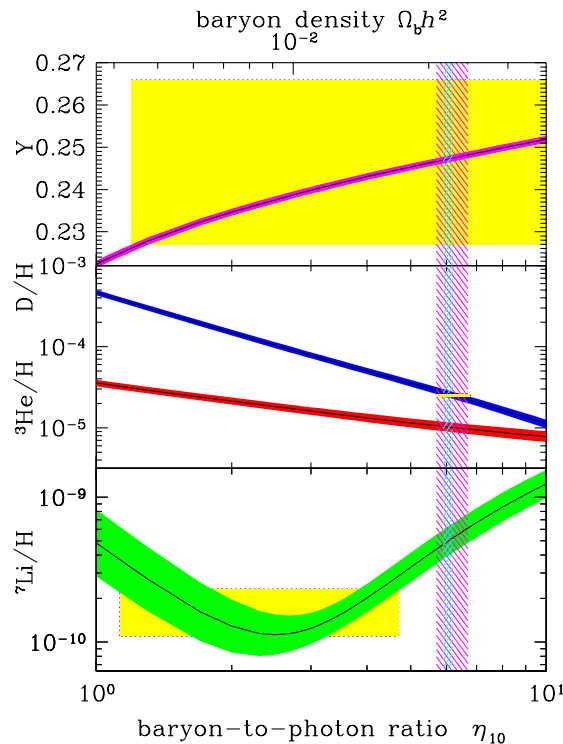
remember: we are forming *nuclei* and not atoms...

BBN summary:

- production of (H), ^2H , ^3H , ^3He , ^4He , ^7Be , ^7Li
- all other elements are in fact produced in stars
- mass fractions:

^2H : deuterium D

^3H : tritium T

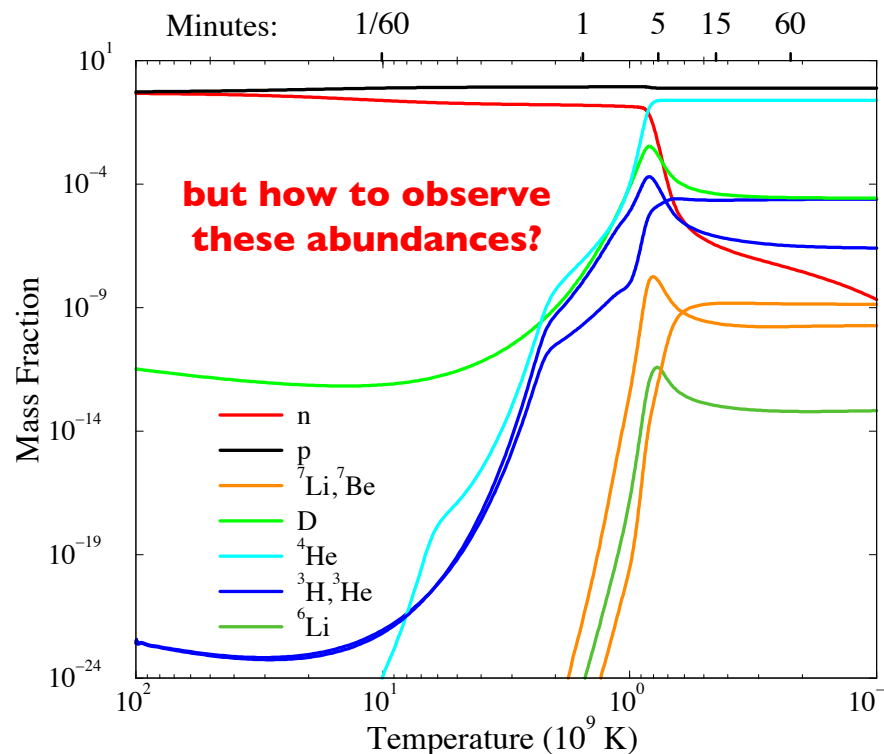
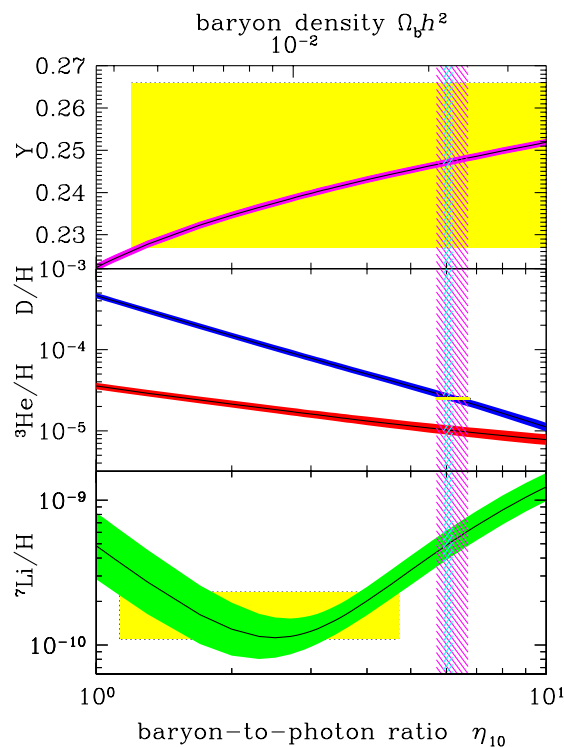


▪ **BBN summary:**

- production of (H), ²H, ³H, ³He, ⁴He, ⁷Be, ⁷Li
- all other elements are in fact produced in stars
- mass fractions:

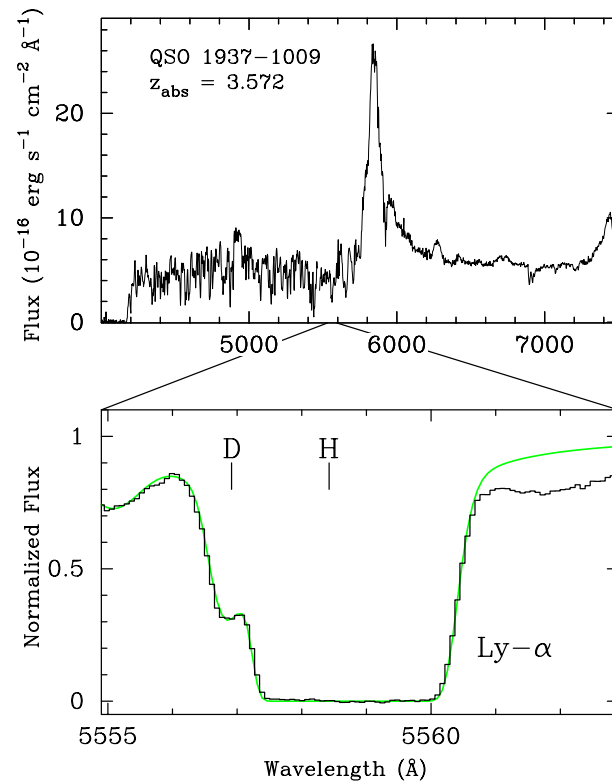
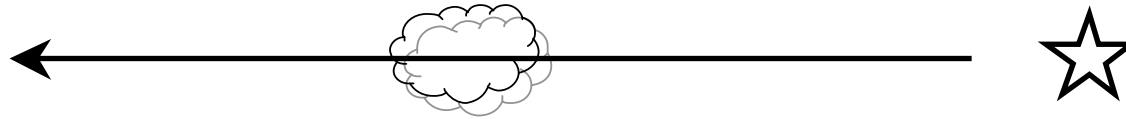
²H: deuterium D

³H: trillium T



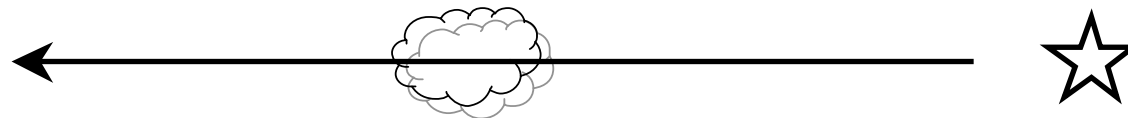
- introduction
- particle physics
- synthesis of elements
- big bang nucleosynthesis
- ***observations***

- Ly- α clouds (not polluted by stars):
 - line strength in QSO absorption spectra provide abundance measures

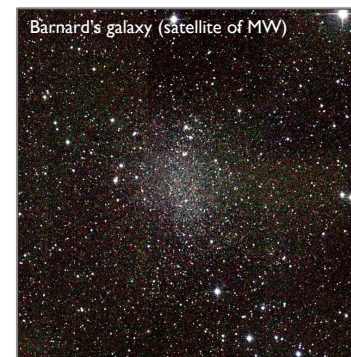


(Burles, Nollett & Turner, arxiv:9903300, Fig2)

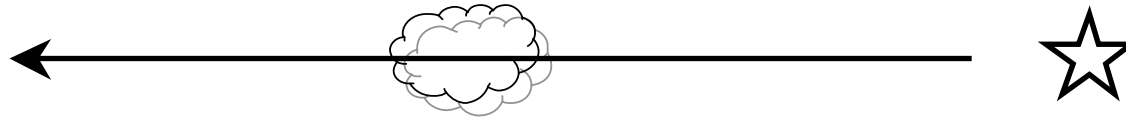
- Ly- α clouds (not polluted by stars):
 - line strength in QSO absorption spectra provide abundance measures



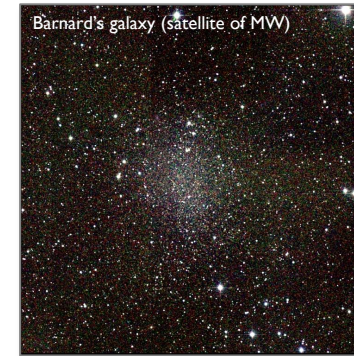
- nearby dwarf galaxies
 - high gas/star ratio and low metal/H in gas suggest that their interstellar medium is close to primordial



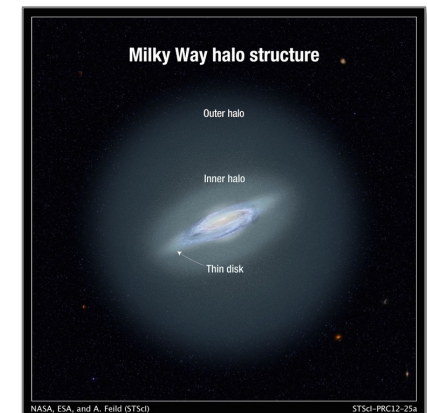
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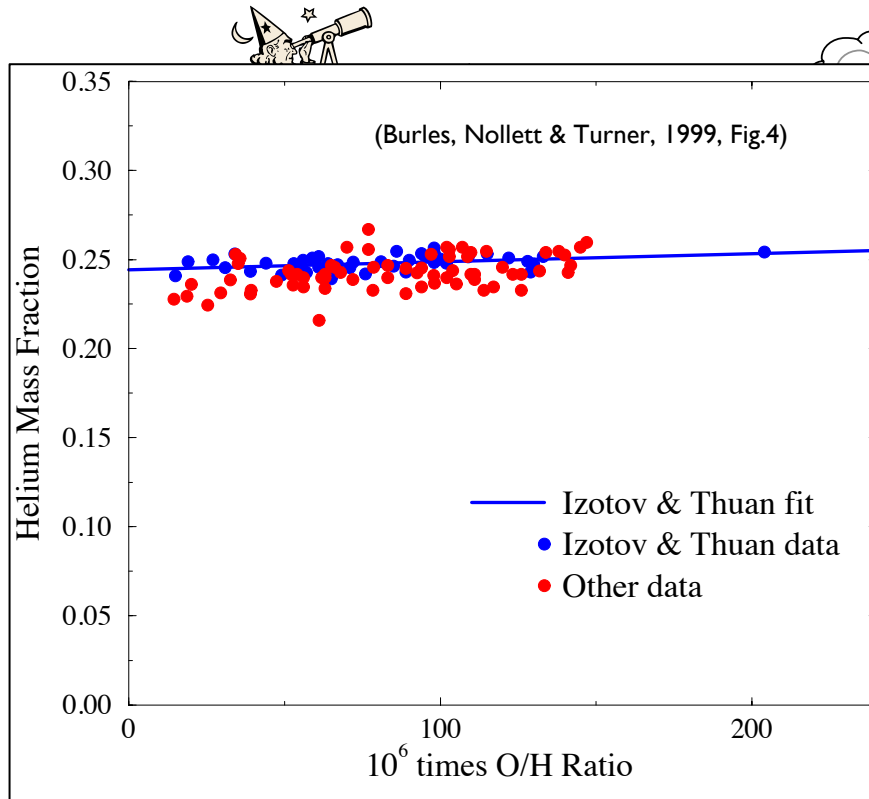
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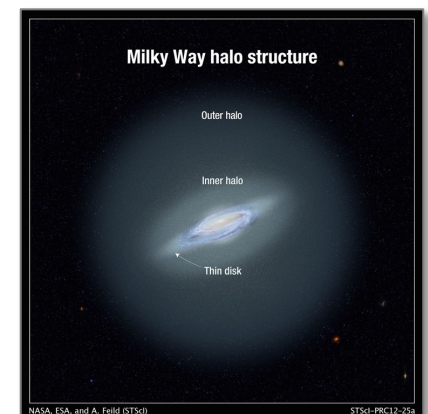
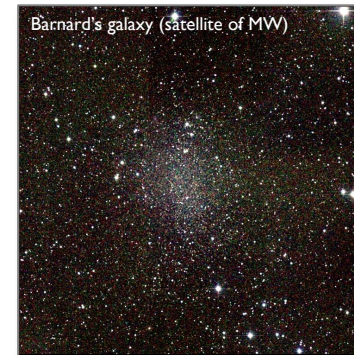
- Galactic halo
 - contains very old stellar population



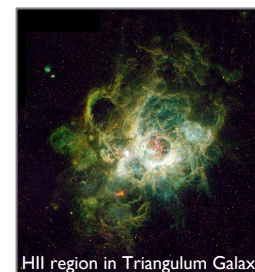
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in gas suggest
e to primordial



- HII regions
 - low-density cloud of partially ionized gas in which star formation took/takes place



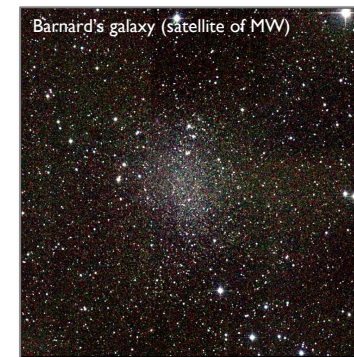
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• IGM observations: $\eta \approx 2.5 \times 10^{-5}$



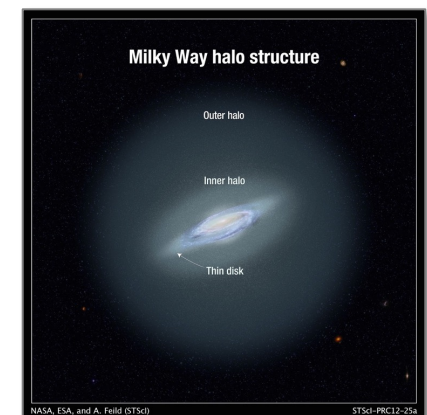
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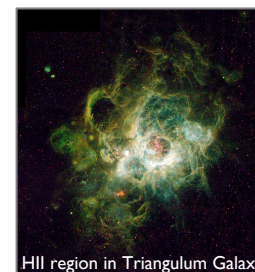
- Galactic halo

• ${}^7\text{Li}$ observed in spectra of cool low-mass stars in Galactic halo



- HII regions

• ${}^4\text{He}$ probed via emission from optical recombination lines in HII regions



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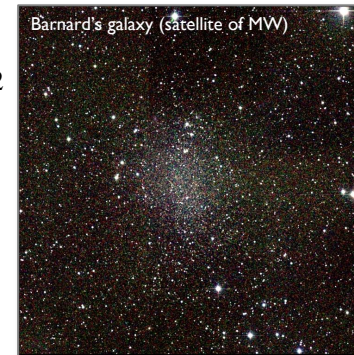
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$$\eta = \frac{n_b}{n_\gamma} = 10^{-10} \eta_{10} = 10^{-10} \cdot 274 \Omega_b h^2$$

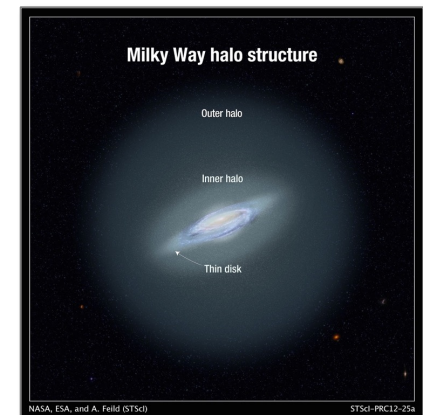
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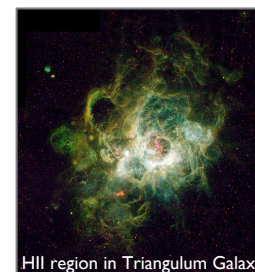
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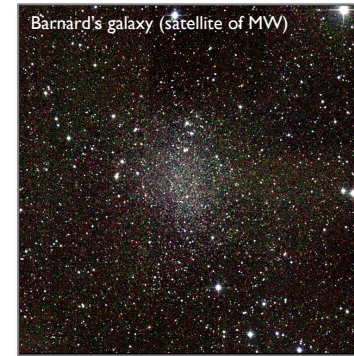


$$\Omega_b h^2 \approx 0.0214$$

- nearby dwarf galaxies



• ISM observations: $\eta \approx 1.6 \times 10^{-5}$

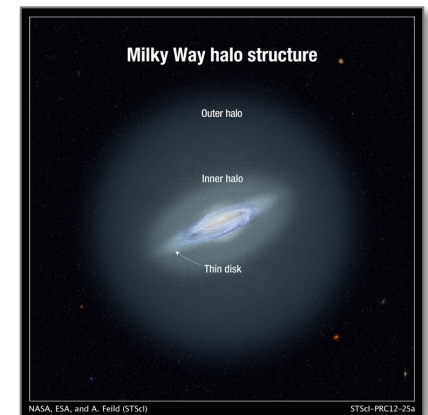


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$$\Omega_b h^2 \approx 0.0214$$

even BBN claims for the existence of non-baryonic matter!

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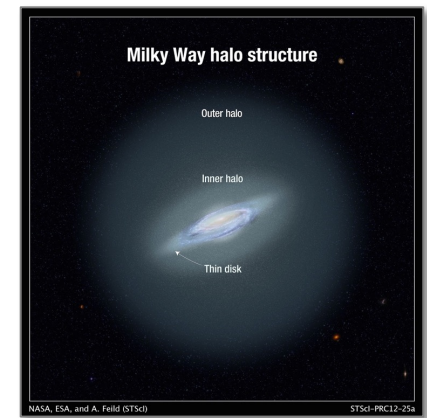


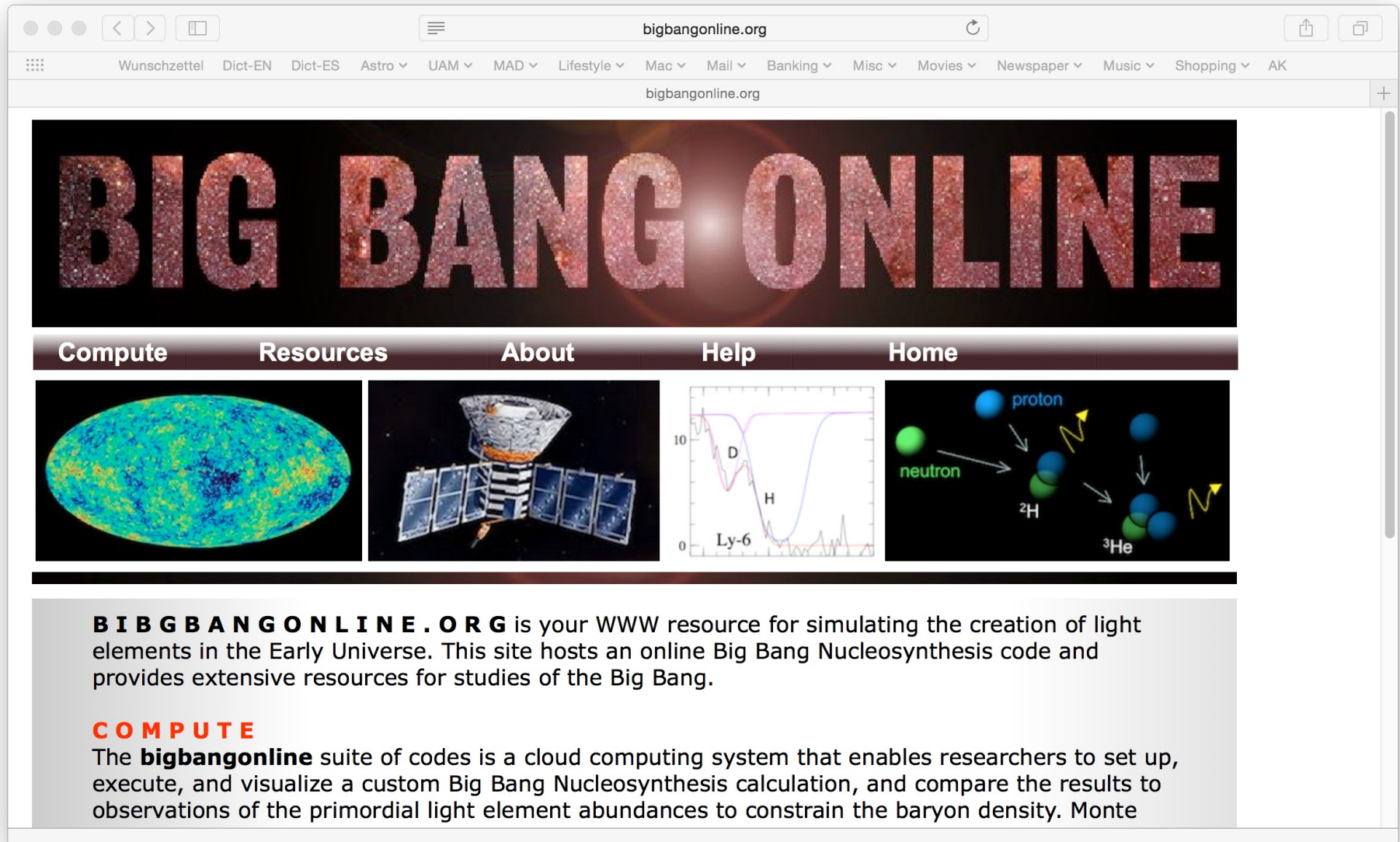
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The screenshot shows a web browser window with the URL `bigbangonline.org`. The browser's address bar and menu bar are visible. The website's main header features the text "BIG BANG ONLINE" in large, red, textured letters. Below this is a navigation bar with five links: "Compute", "Resources", "About", "Help", and "Home".

The main content area contains four images:

- A Cosmic Microwave Background (CMB) fluctuation map.
- An illustration of a satellite with solar panels.
- A plot showing the abundance of light elements (D, H, Ly-6) versus time.
- A diagram illustrating the fusion of a proton and a neutron to form ^2H and ^3He .

B I B G B A N G O N L I N E . O R G is your WWW resource for simulating the creation of light elements in the Early Universe. This site hosts an online Big Bang Nucleosynthesis code and provides extensive resources for studies of the Big Bang.

COMPUTE
The **bigbangonline** suite of codes is a cloud computing system that enables researchers to set up, execute, and visualize a custom Big Bang Nucleosynthesis calculation, and compare the results to observations of the primordial light element abundances to constrain the baryon density. Monte