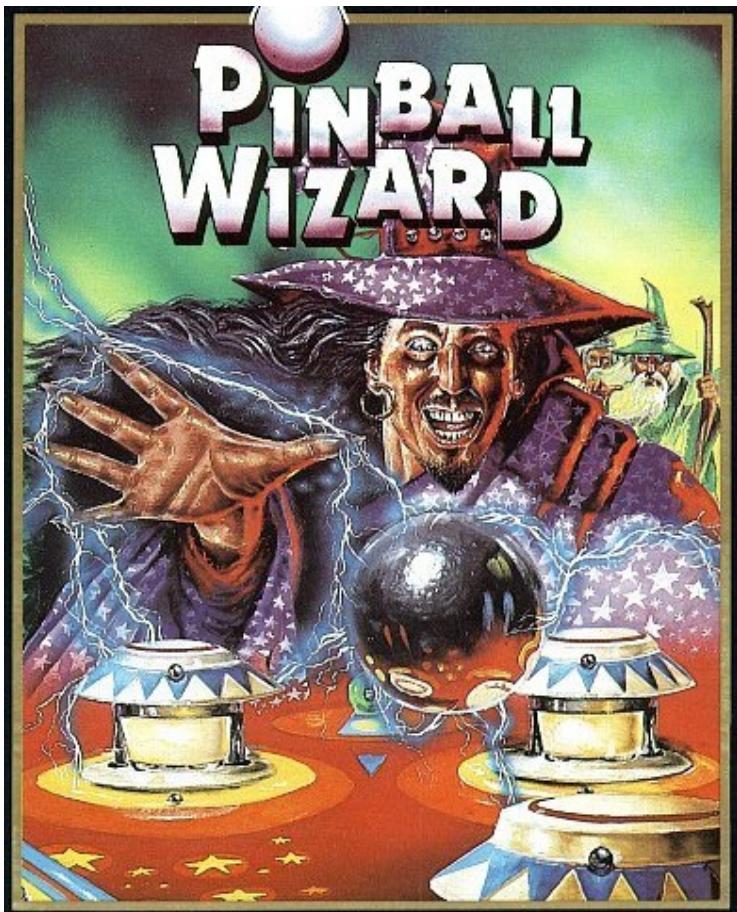


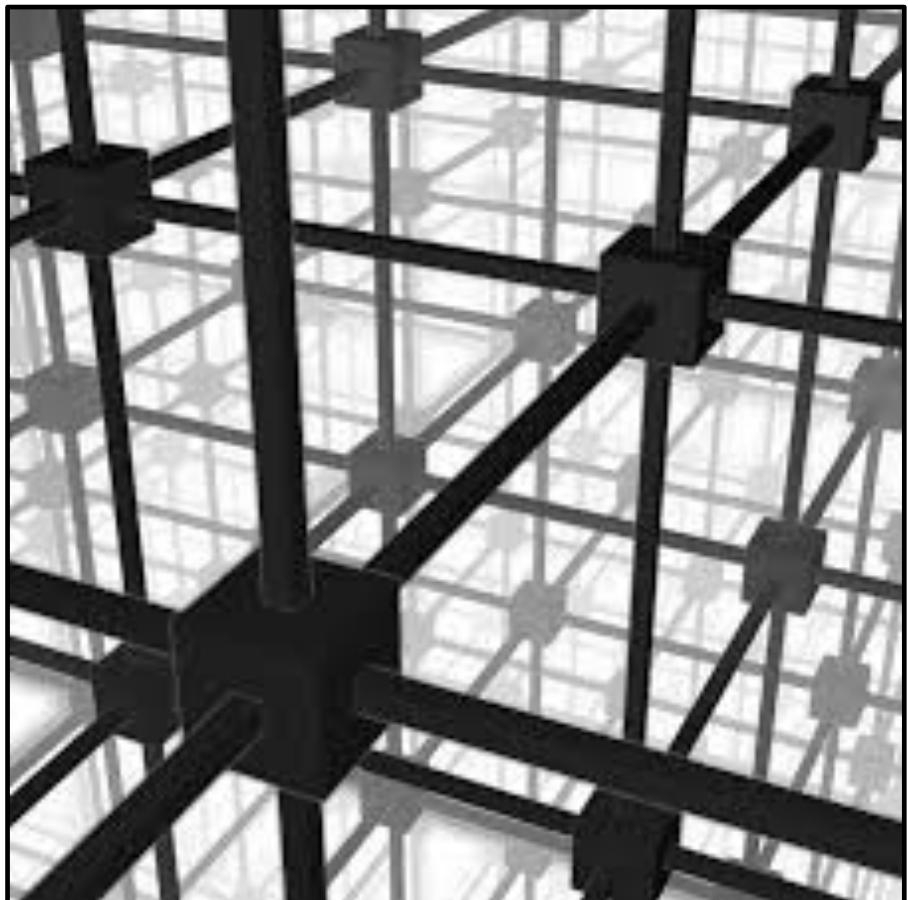
Computational Astrophysics

Solving for Gravity

Alexander Knebe, Universidad Autonoma de Madrid



VS



Solving for Gravity

■ full set of equations

- collisionless matter (e.g. dark matter)

$$\frac{d\vec{x}_{DM}}{dt} = \vec{v}_{DM}$$

$$\frac{d\vec{v}_{DM}}{dt} = -\nabla\phi$$

- Poisson's equation

$$\Delta\phi = 4\pi G \rho_{tot}$$

- collisional matter (e.g. gas)

$$\frac{\partial\rho}{\partial t} + \nabla \cdot (\rho\vec{v}) = 0$$

$$\frac{\partial(\rho\vec{v})}{\partial t} + \nabla \cdot \left(\rho\vec{v} \otimes \vec{v} + \left(p + \frac{1}{2\mu} B^2 \right) \vec{I} - \frac{1}{\mu} \vec{B} \otimes \vec{B} \right) = \rho (-\nabla\phi)$$

$$\frac{\partial(\rho E)}{\partial t} + \nabla \cdot \left(\left[\rho E + p + \frac{1}{2\mu} B^2 \right] \vec{v} - \frac{1}{\mu} [\vec{v} \cdot \vec{B}] \vec{B} \right) = \rho\vec{v} \cdot (-\nabla\phi) + (\Gamma - L)$$

- ideal gas equations

$$p = (\gamma - 1)\rho\varepsilon$$

$$\rho\varepsilon = \rho E - \frac{1}{2}\rho v^2$$

- Maxwell's equation

$$\frac{\partial\vec{B}}{\partial t} = -\nabla \times (\vec{v} \times \vec{B})$$

Solving for Gravity

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$$\begin{aligned}\frac{d\vec{x}_{DM}}{dt} &= \vec{v}_{DM} \\ \frac{d\vec{v}_{DM}}{dt} &= -\nabla\phi\end{aligned}$$

leap-frog integration

- Poisson's equation

$$\Delta\phi = 4\pi G\rho_{tot}$$

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later...

$$\begin{aligned} \frac{\partial\rho}{\partial t} + \nabla \cdot (\rho\vec{v}) &= 0 \\ \frac{\partial(\rho\vec{v})}{\partial t} + \nabla \cdot \left(\rho\vec{v} \otimes \vec{v} + \left(p + \frac{1}{2\mu} B^2 \right) \vec{I} - \frac{1}{\mu} \vec{B} \otimes \vec{B} \right) &= \rho (-\nabla\phi) \\ \frac{\partial(\rho E)}{\partial t} + \nabla \cdot \left(\left[\rho E + p + \frac{1}{2\mu} B^2 \right] \vec{v} - \frac{1}{\mu} [\vec{v} \cdot \vec{B}] \vec{B} \right) &= \rho\vec{v} \cdot (-\nabla\phi) + (\Gamma - L) \end{aligned}$$

hyperbolic partial differential equations:
solutions are wave-like, i.e. perturbations need time to travel...

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electrodynamics lecture...

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physics lecture...

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now!

- Poisson's equation

$$\Delta\phi = 4\pi G \rho_{tot}$$

**elliptical partial differential equation:
solution obtainable via FFT (for constant coefficients)**

- ideal gas equations

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Computational Astrophysics

Solving for Gravity

- Poisson's equation vs. Poisson's integral
- tree codes
- PM codes

- **Poisson's equation vs. Poisson's integral**
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- PM codes

Solving for Gravity

- Poisson's equation

$$\Delta\Phi(\vec{x}) = 4\pi G\rho(\vec{x})$$

Solving for Gravity

■ Poisson's equation

$$\Delta\Phi(\vec{x}) = 4\pi G \rho(\vec{x})$$

■ Poisson's integral

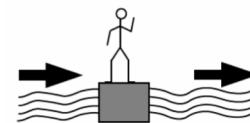
$$\Phi(\vec{x}) = \iiint G(\vec{x} - \vec{x}') \rho(\vec{x}') d^3x' ; \quad G(\vec{x}) = \frac{1}{4\pi x}$$

Solving for Gravity

- Poisson's equation

$$\Delta\Phi(\vec{x}) = 4\pi G \rho(\vec{x})$$

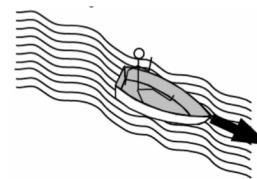
grid approach



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particle approach



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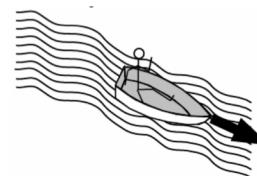
grid approach

$$\Delta\Phi(\vec{r}_{i,j,k}) = 4\pi G \rho(\vec{r}_{i,j,k})$$

($\vec{r}_{i,j,k}$ = position of centre of grid cell (i,j,k))

■ Poisson's integral

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particle approach

Solving for Gravity

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$$\rho(\vec{r}) = \sum_{i=1}^N m_i \delta_{\text{Dirac}}(\vec{r} - \vec{r}_i)$$

particle approach

$$\vec{F}(\vec{r}_i) = - \sum_{i \neq j} \frac{Gm_i m_j}{(r_i - r_j)^3} (\vec{r}_i - \vec{r}_j)$$

Solving for Gravity

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grid approach

$$\Delta\Phi(\vec{r}_{i,j,k}) = 4\pi G \rho(\vec{r}_{i,j,k})$$

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weapon of choice: AMR codes

■ Poisson's integral

$$\Phi(\vec{x}) = \iiint G(\vec{x} - \vec{x}') \rho(\vec{x}') d^3x' ; \quad G(\vec{x}) = \frac{1}{4\pi x}$$

weapon of choice: tree codes

particle approach

$$\vec{F}(\vec{r}_i) = - \sum_{i \neq j} \frac{Gm_i m_j}{(r_i - r_j)^3} (\vec{r}_i - \vec{r}_j)$$

Computational Astrophysics

Solving for Gravity

- Poisson's equation vs. Poisson's integral
- **tree codes**
- PM codes

Solving for Gravity

- direct particle-particle summation (PP)

$$\vec{F}_i(\vec{r}_i) = - \sum_{i \neq j} \frac{Gm_i m_j}{(r_i - r_j)^3} (\vec{r}_i - \vec{r}_j) \quad \forall i \in N$$

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✓ advantage: easy to code

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$$N \times N = N^2$$

- | | |
|--------------|--|
| ✓ advantage: | easy to code |
| ✗ drawback: | extremely time consuming (N^2 operations) |

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overcoming the “ N^2 ” issue?!

- | | |
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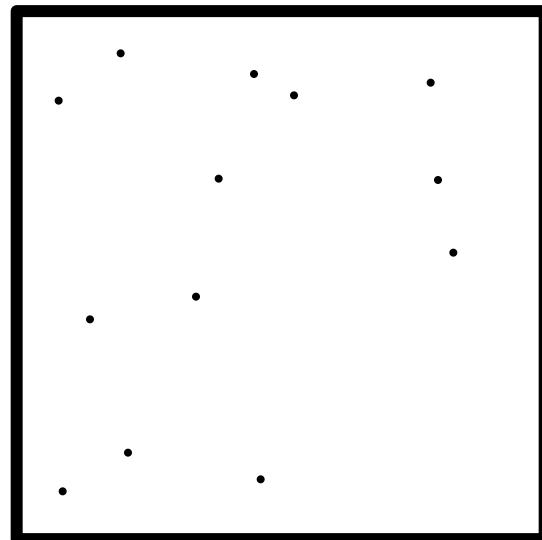
organizing particles into a “tree structure” will give $N \log(N)$ operations

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- generating the tree:

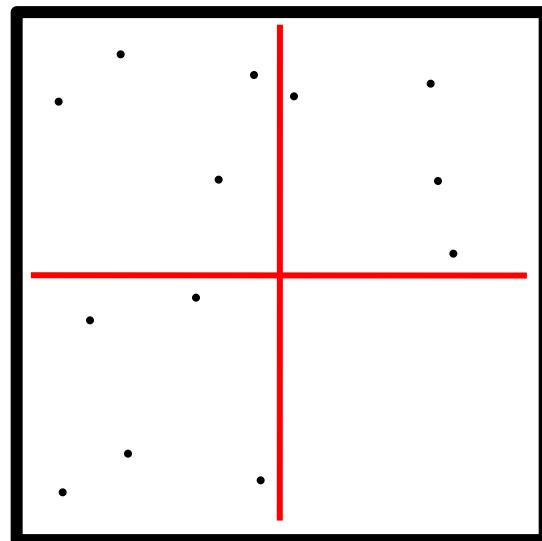


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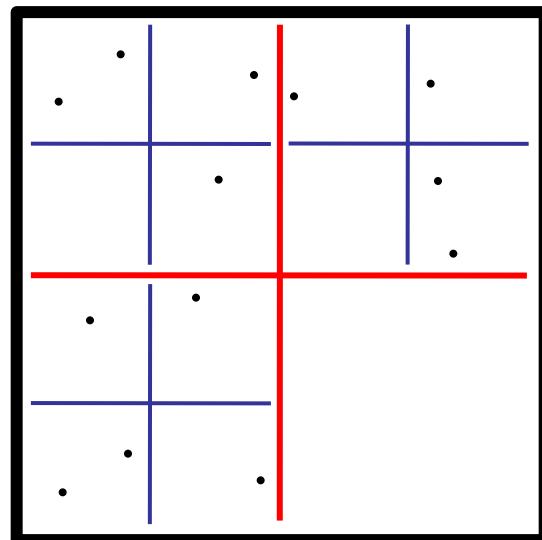


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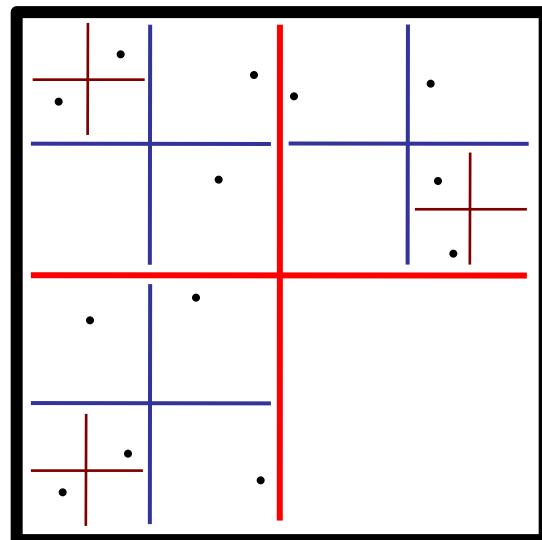


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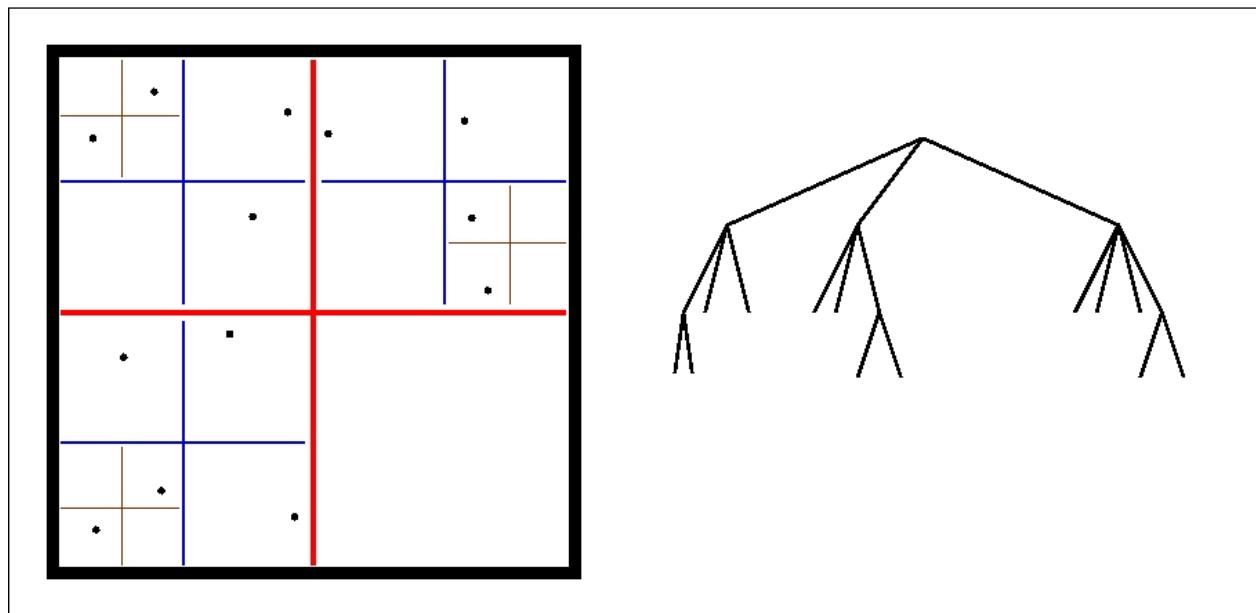


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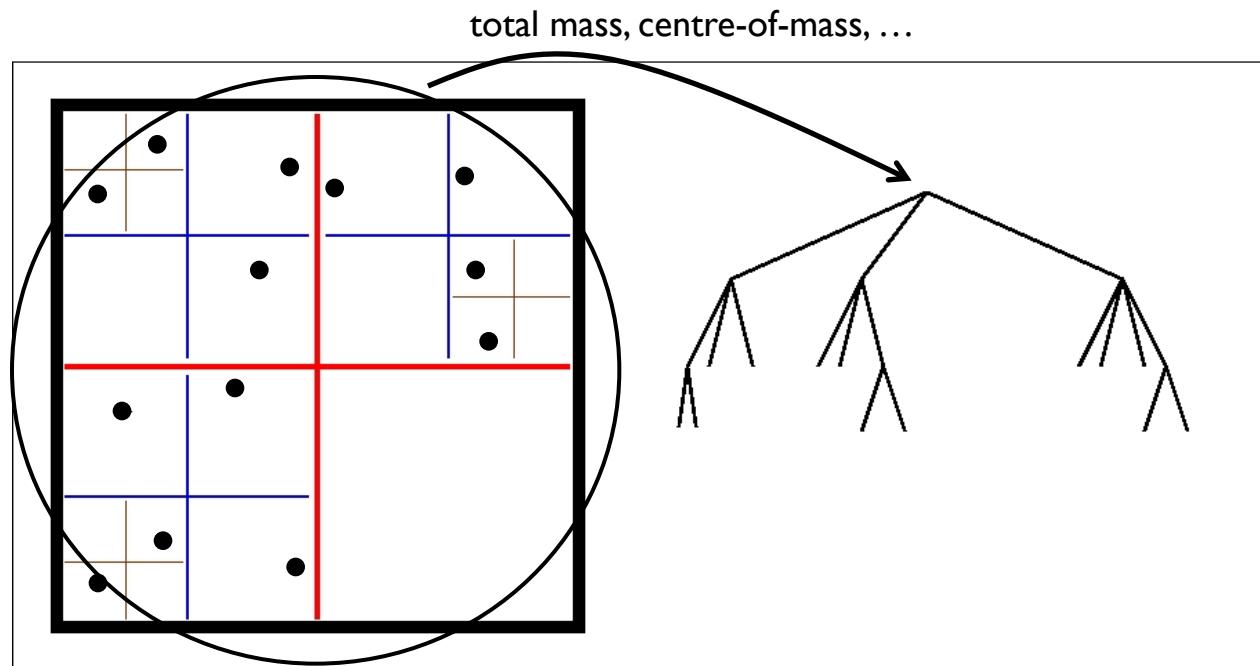


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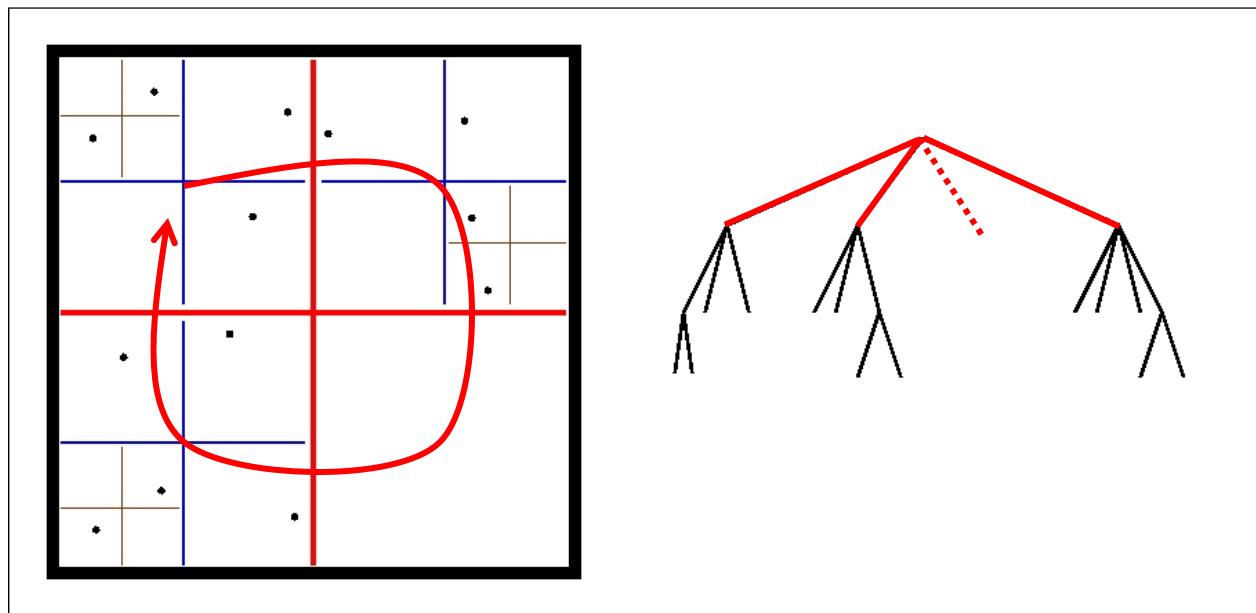


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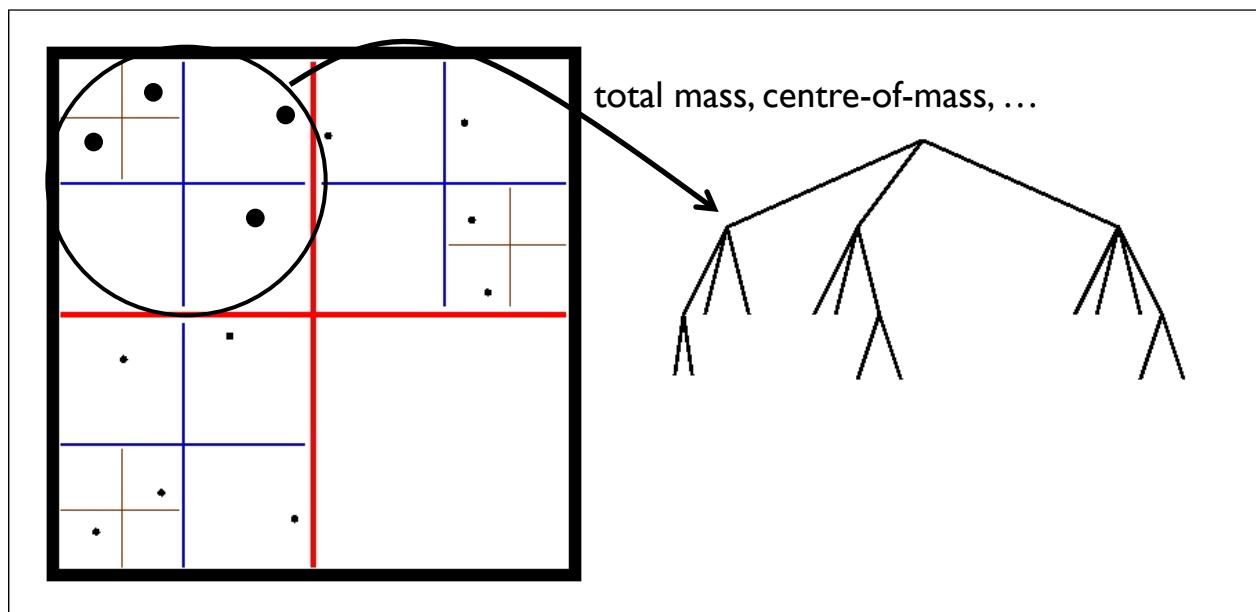


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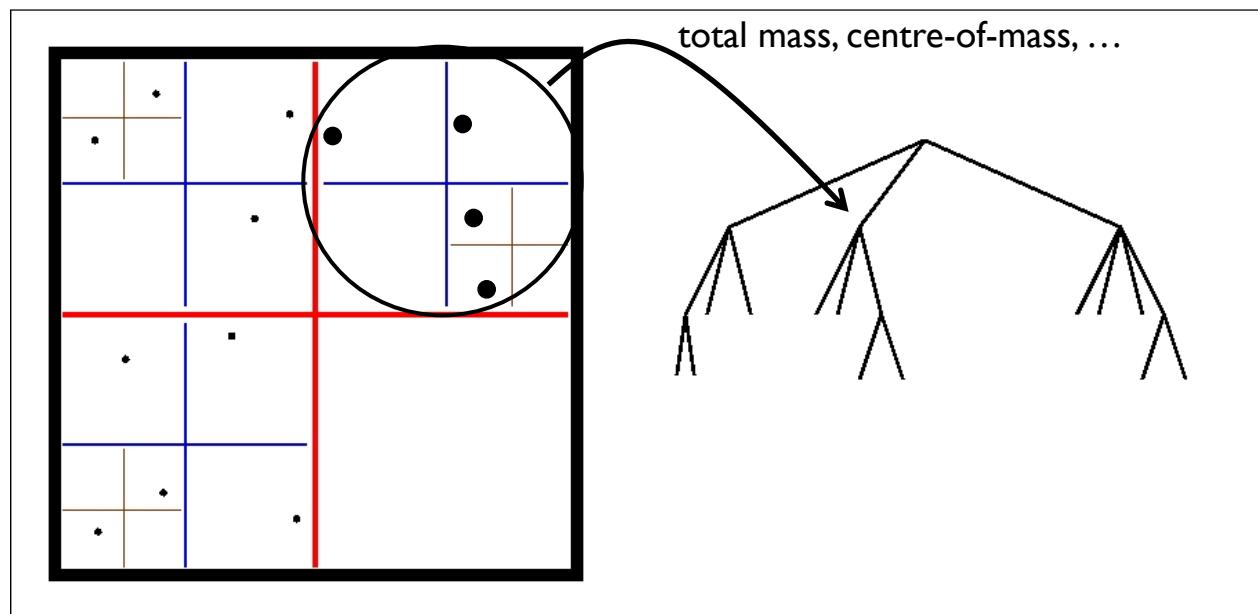


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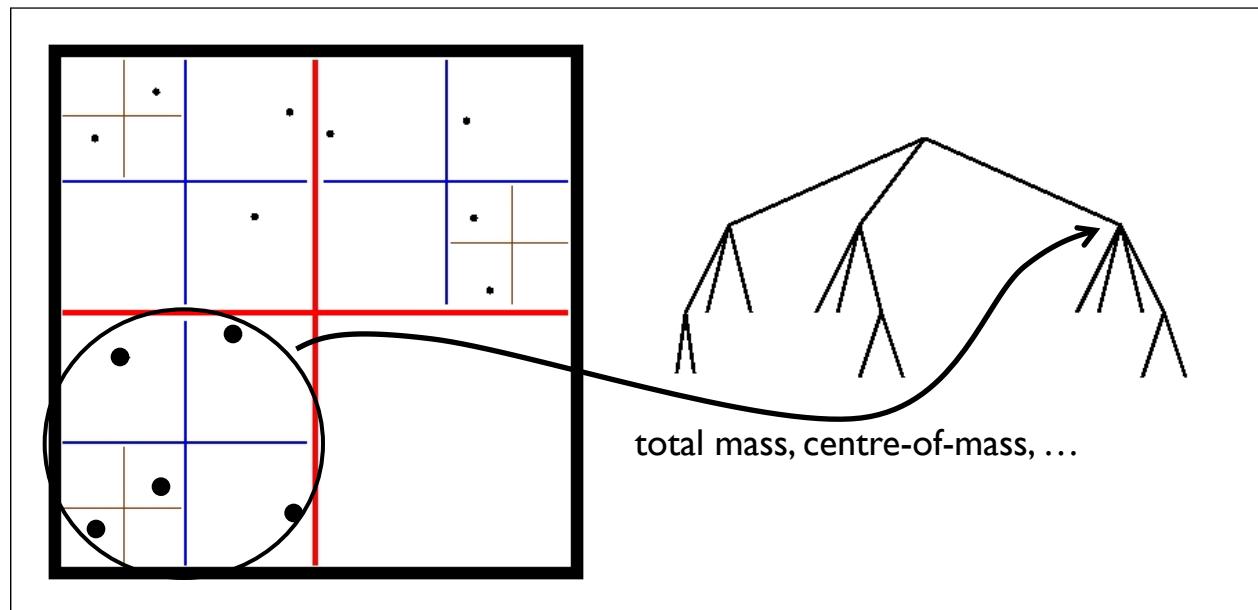


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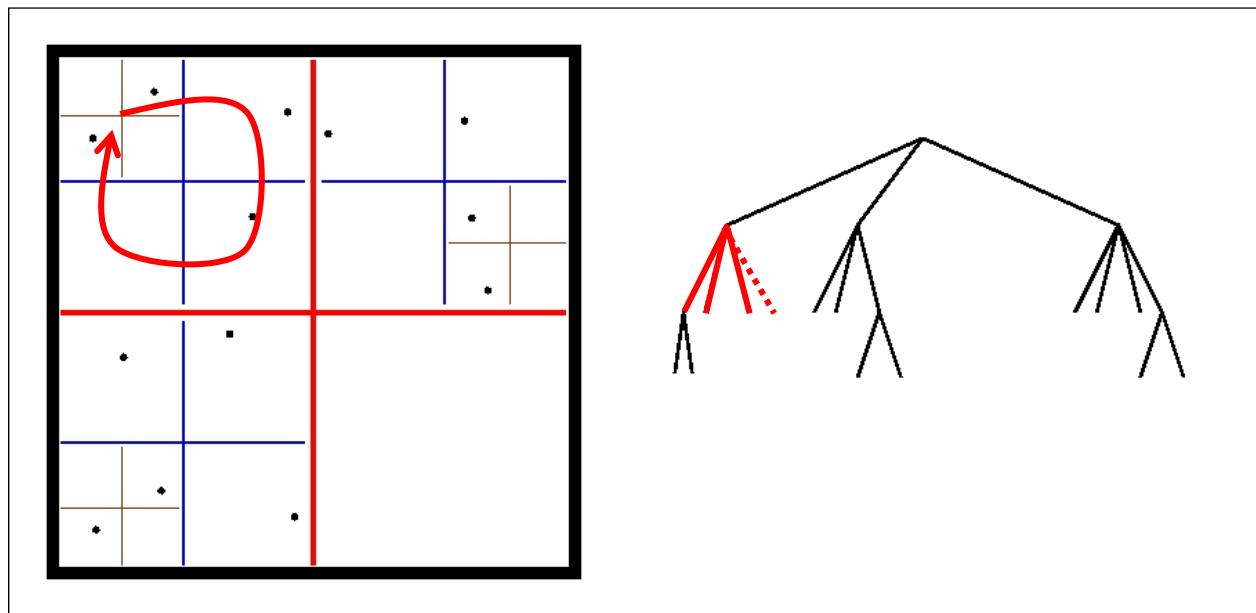


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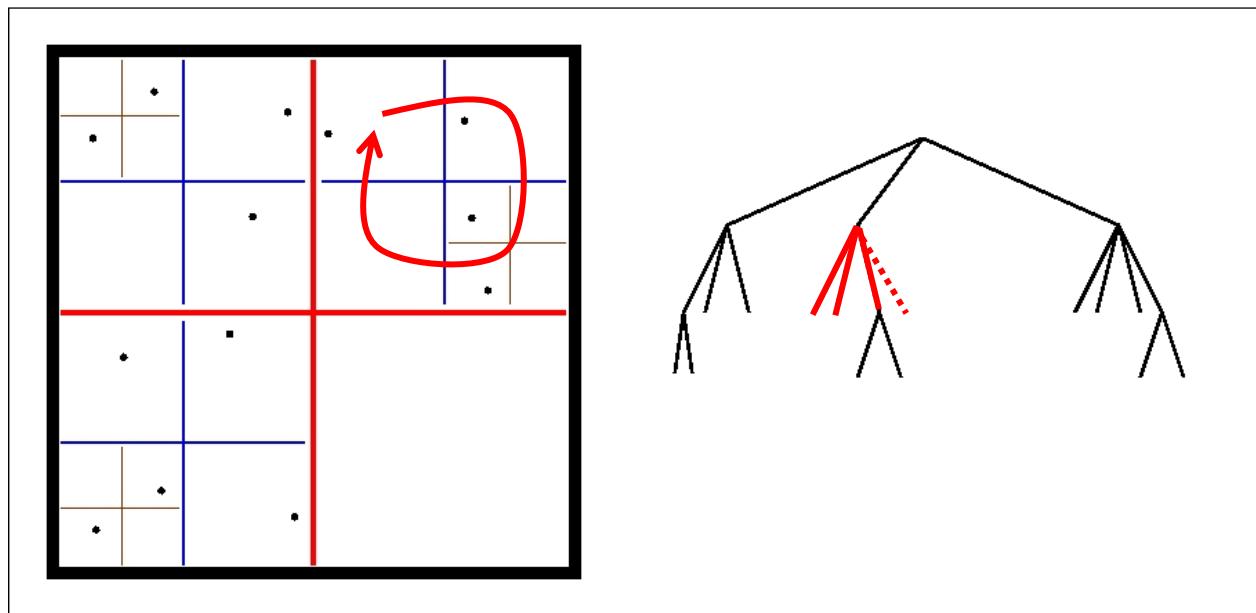


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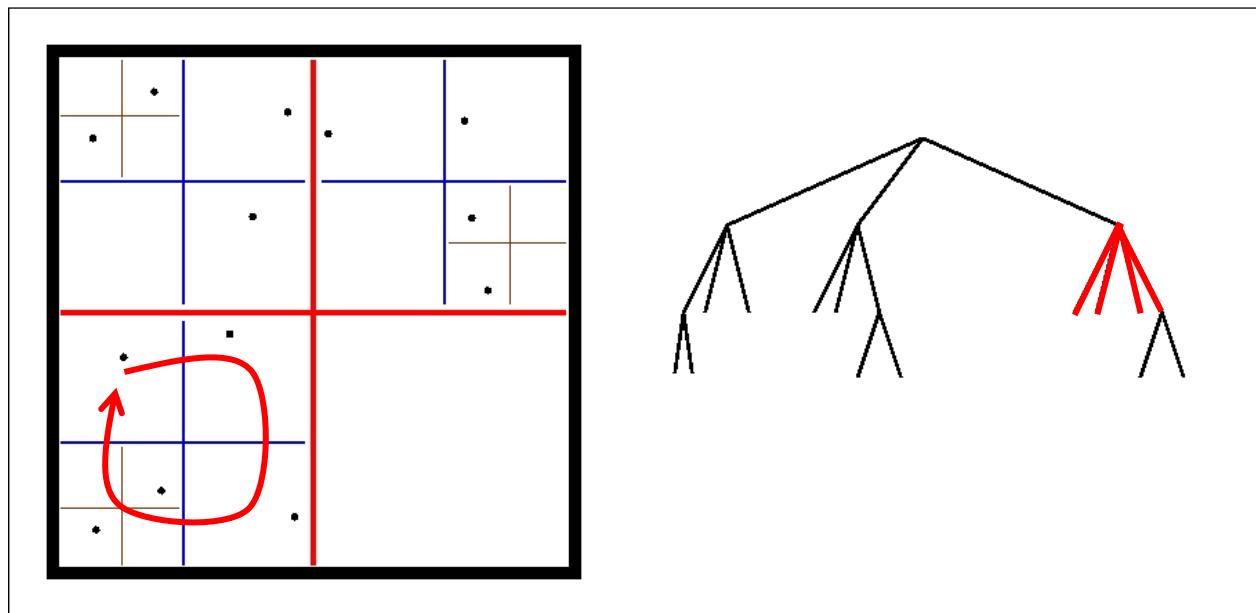


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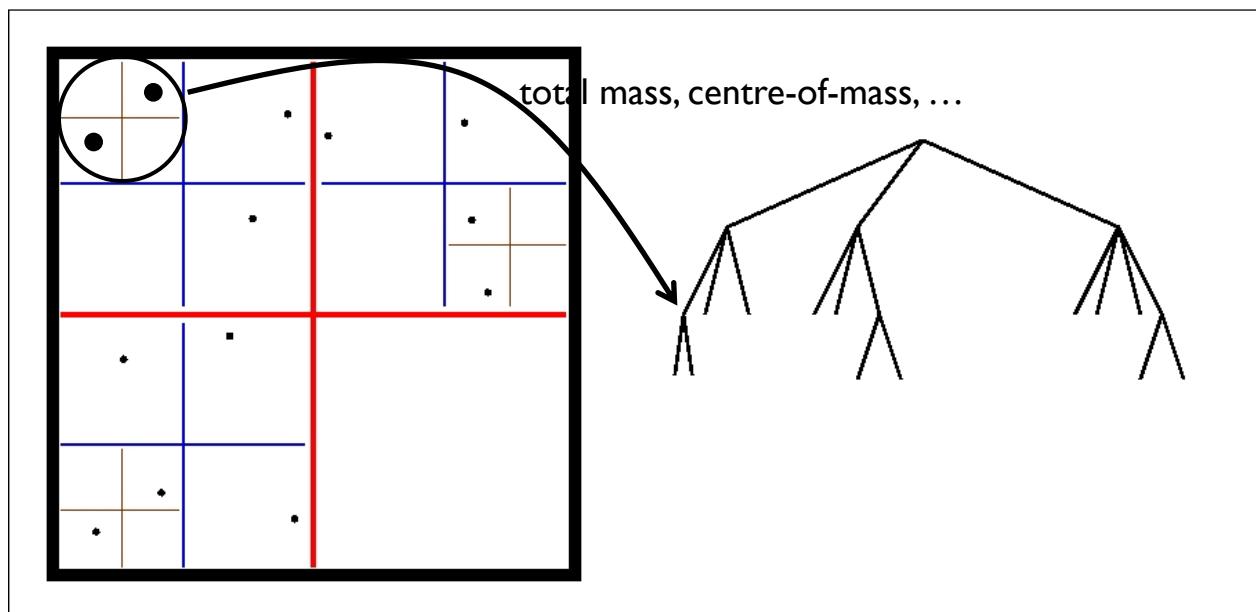


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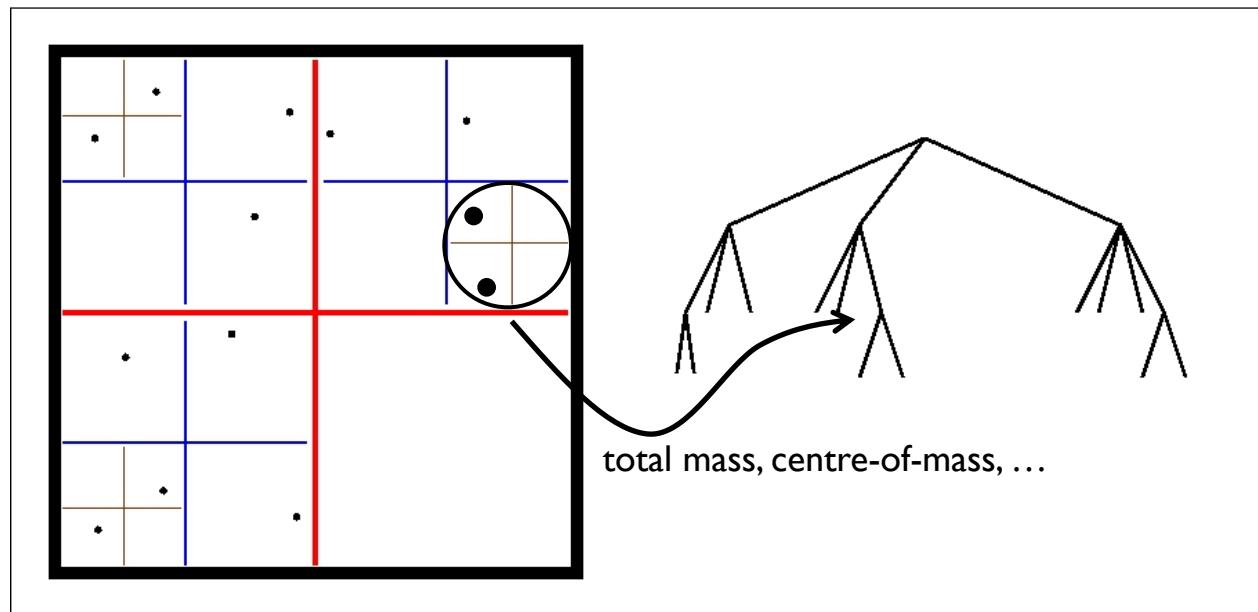


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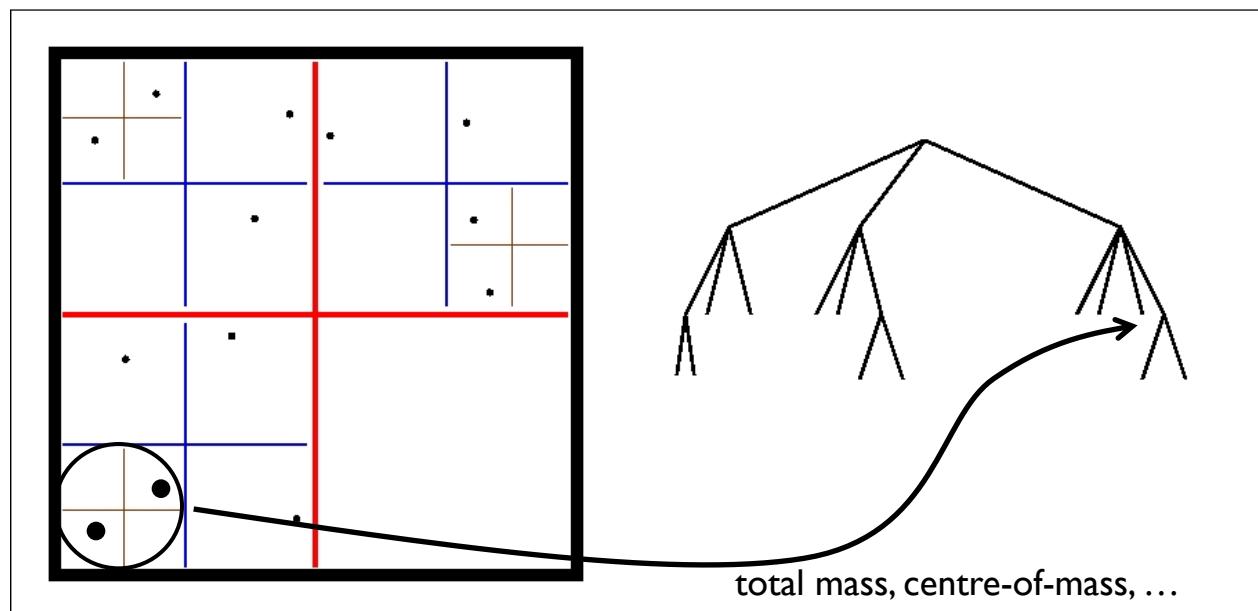


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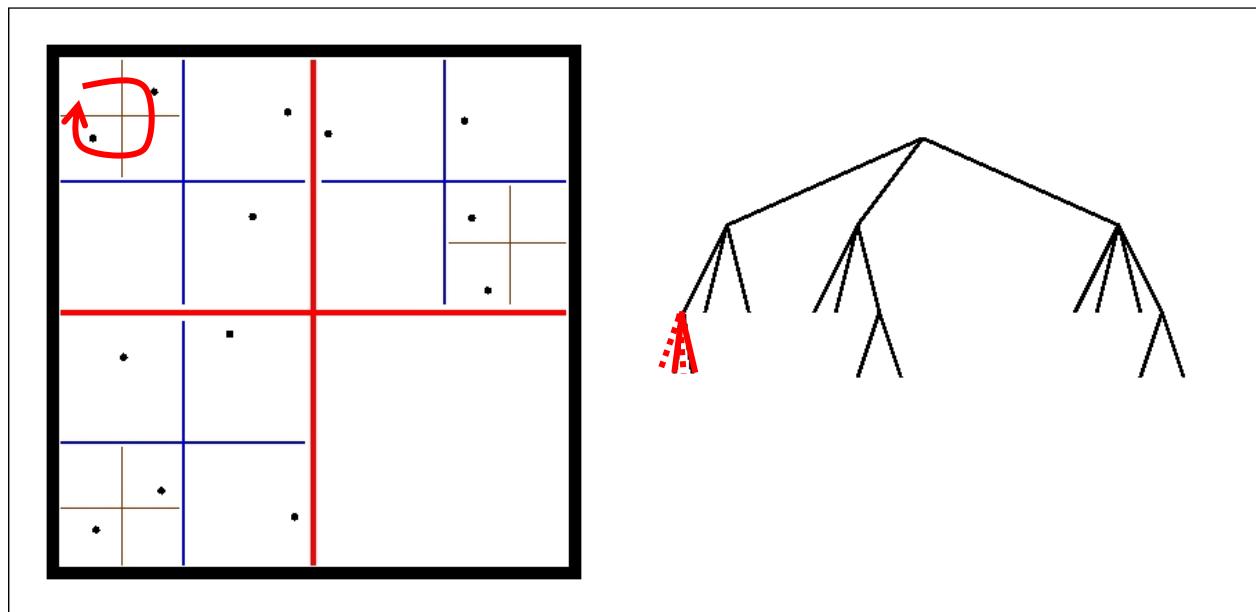


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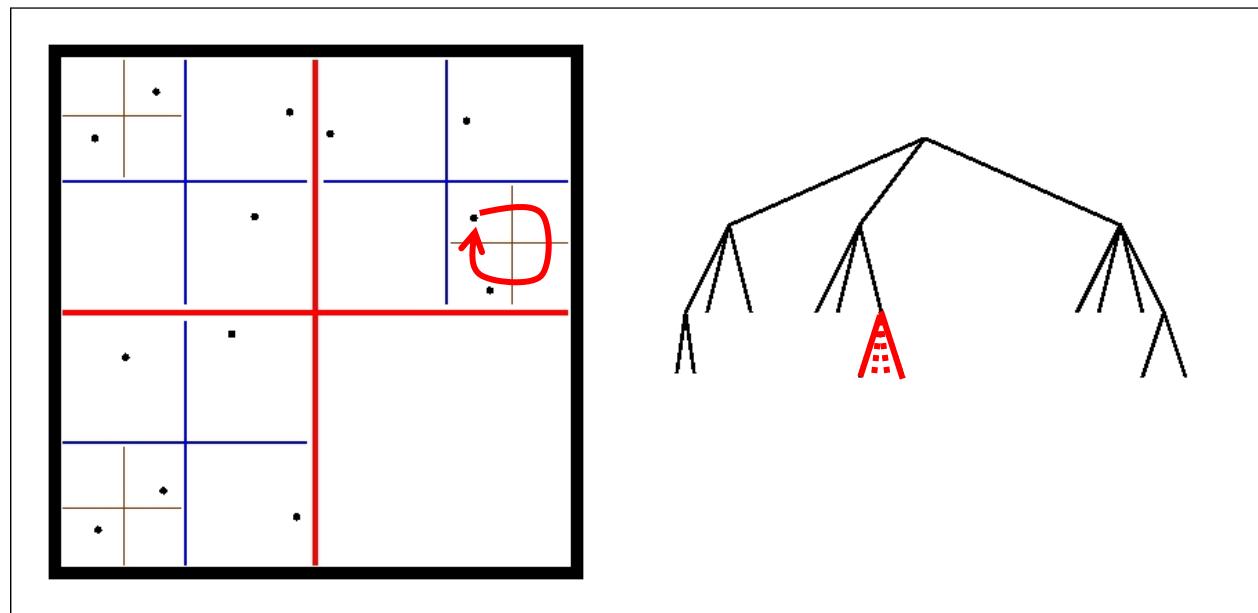


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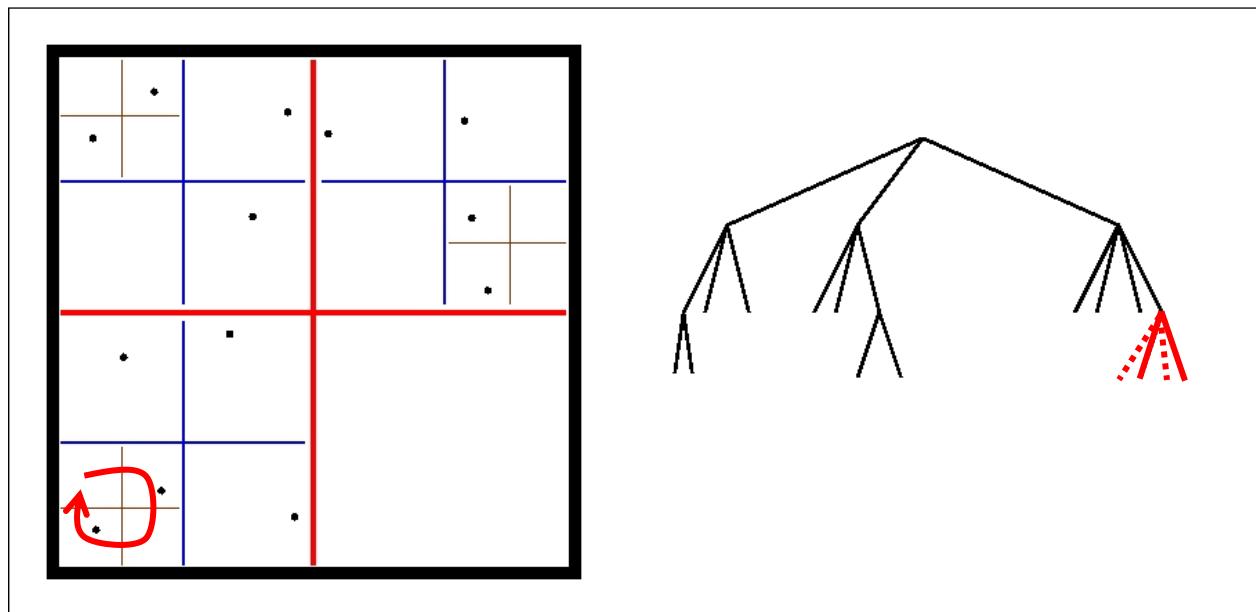


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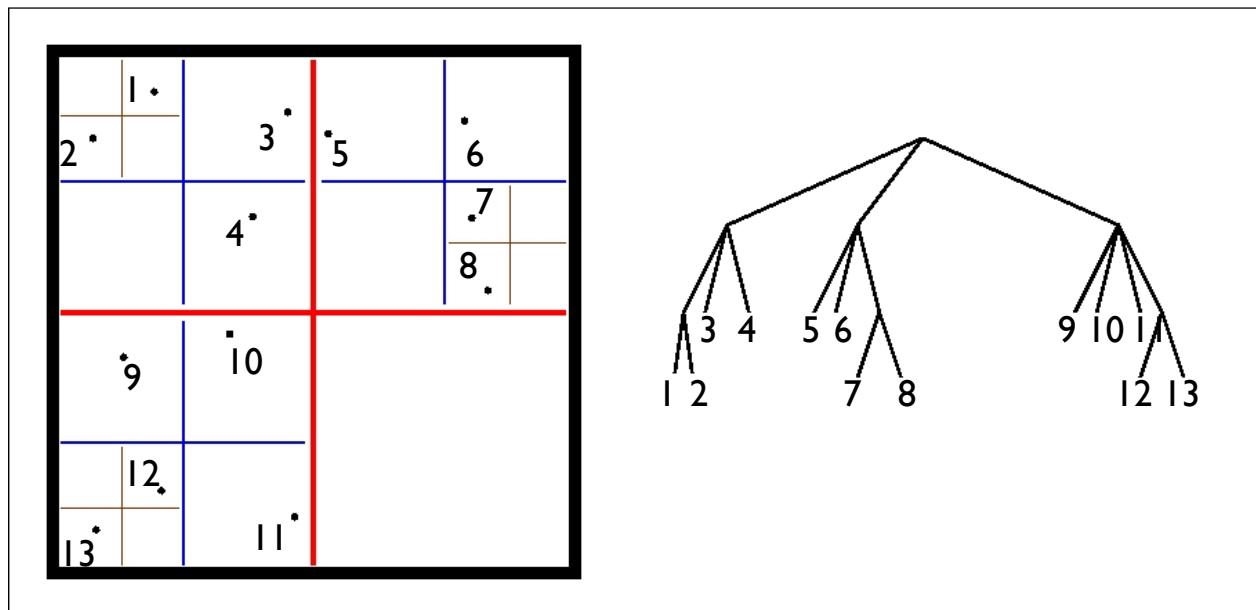
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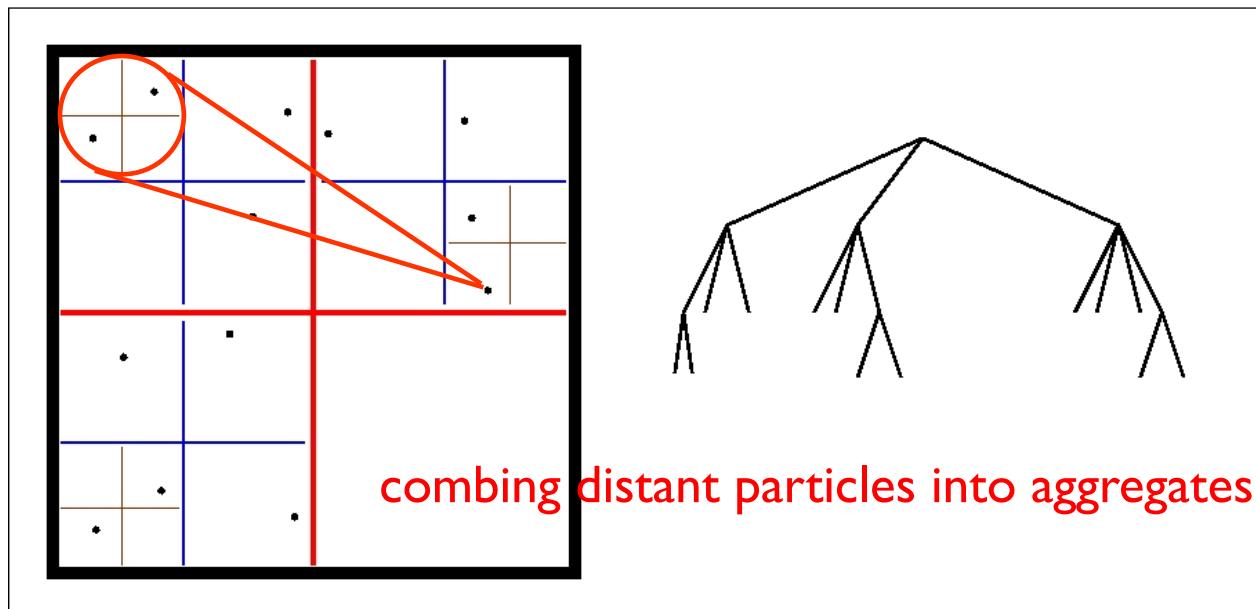


Solving for Gravity

- direct particle-particle summation (PP)

$$\vec{F}_i(\vec{r}_i) = - \sum_{i \neq j} \frac{Gm_i m_j}{(r_i - r_j)^3} (\vec{r}_i - \vec{r}_j) \quad \forall i \in N$$

- walking the tree ($\forall i \in N$):

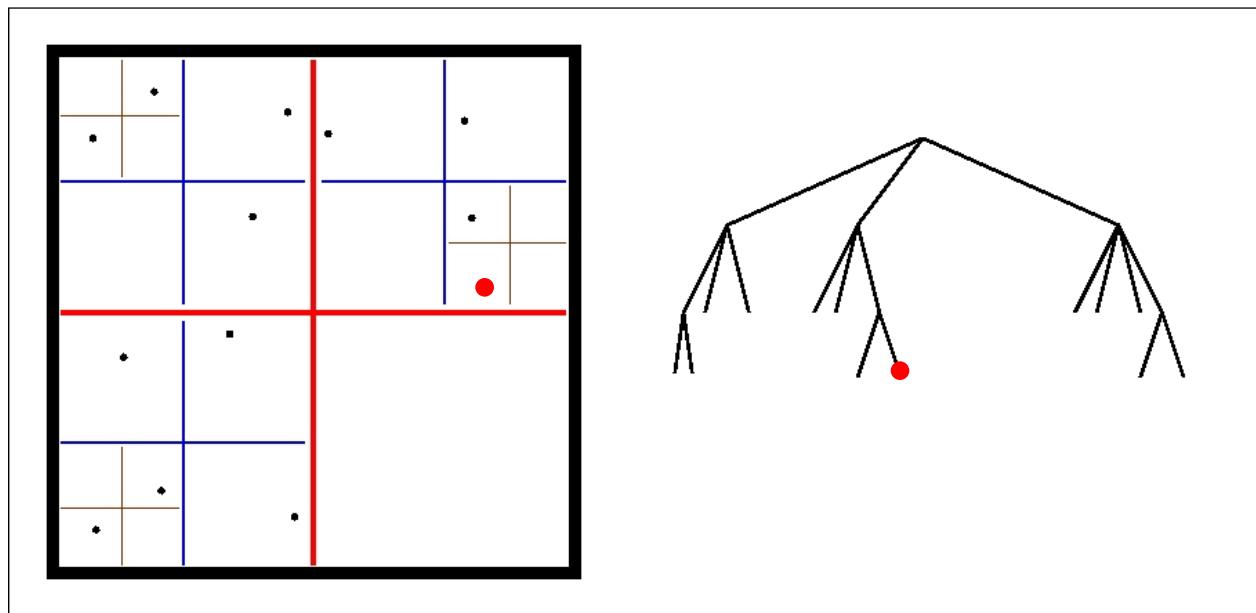


Solving for Gravity

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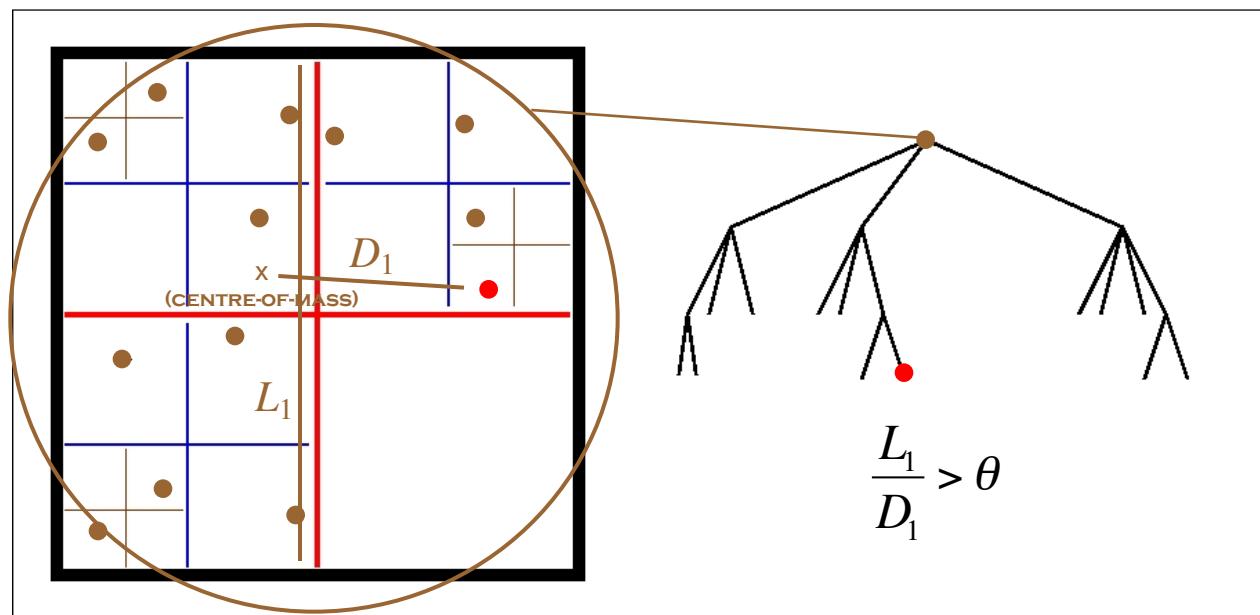


Solving for Gravity

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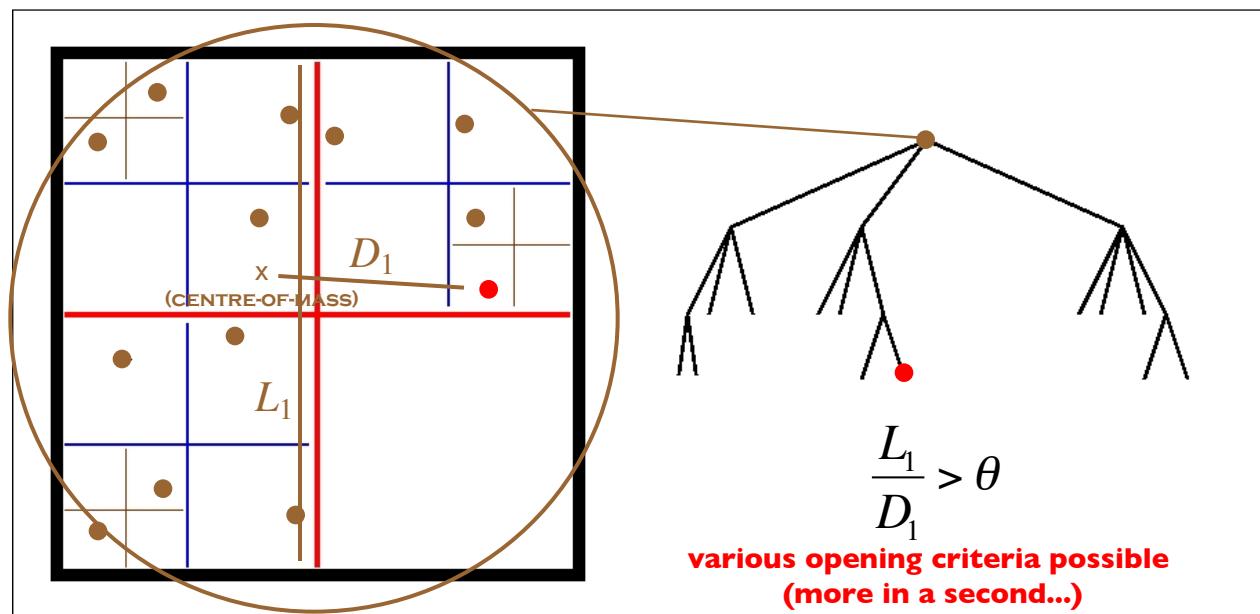


Solving for Gravity

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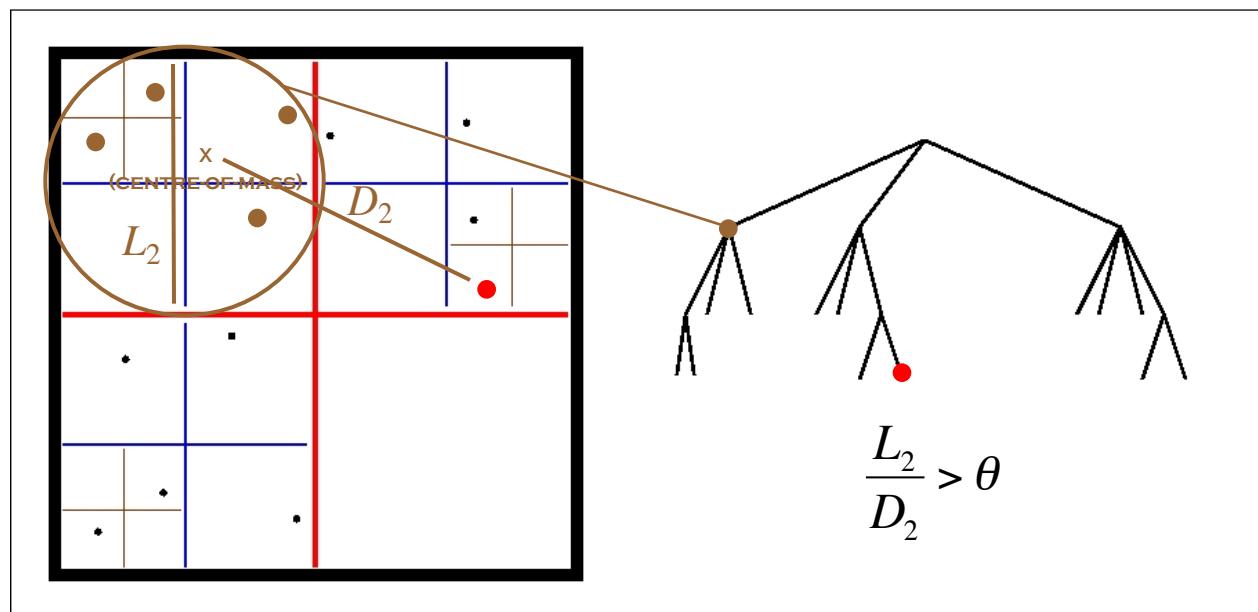


Solving for Gravity

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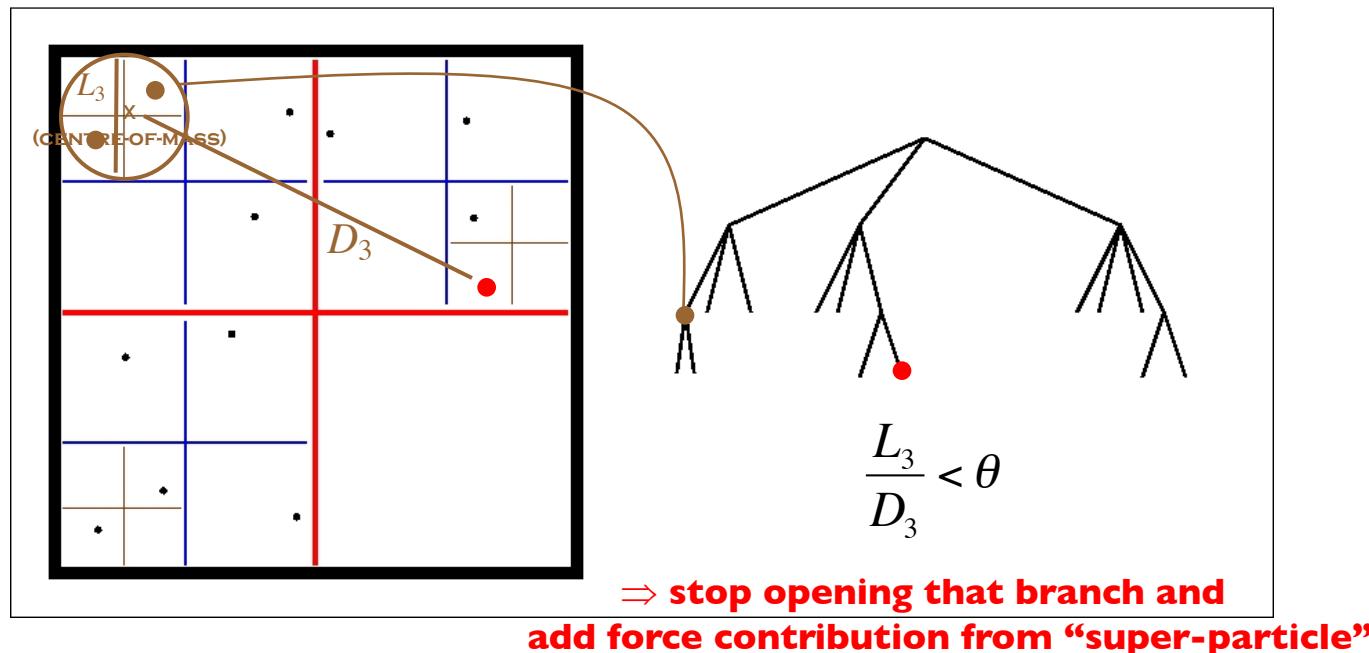


Solving for Gravity

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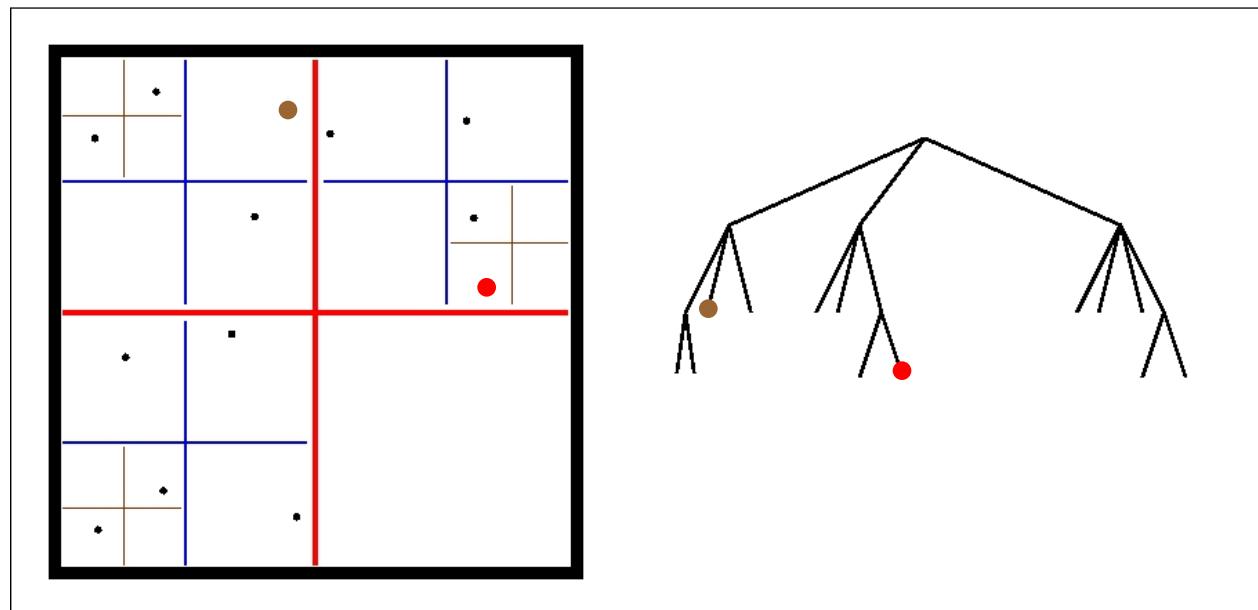


Solving for Gravity

- direct particle-particle summation (PP)

$$\vec{F}_i(\vec{r}_i) = - \sum_{i \neq j} \frac{Gm_i m_j}{(r_i - r_j)^3} (\vec{r}_i - \vec{r}_j) \quad \forall i \in N$$

- walking the tree ($\forall i \in N$):



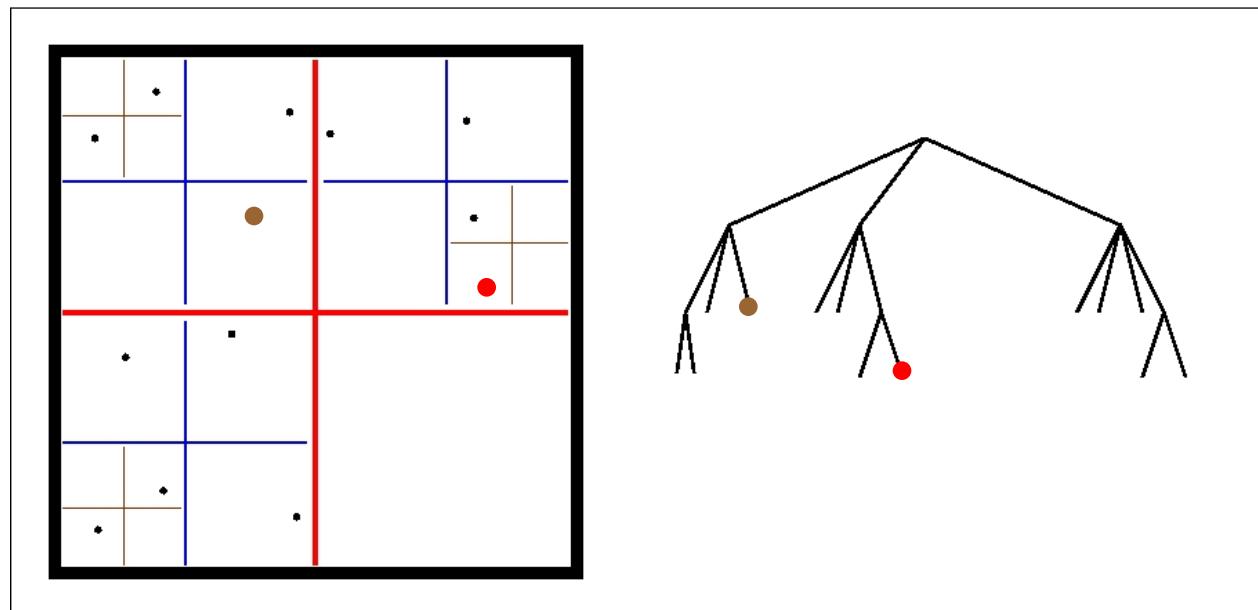
we still need to add the remaining contributions from that branch...

Solving for Gravity

- direct particle-particle summation (PP)

$$\vec{F}_i(\vec{r}_i) = - \sum_{i \neq j} \frac{Gm_i m_j}{(r_i - r_j)^3} (\vec{r}_i - \vec{r}_j) \quad \forall i \in N$$

- walking the tree ($\forall i \in N$):

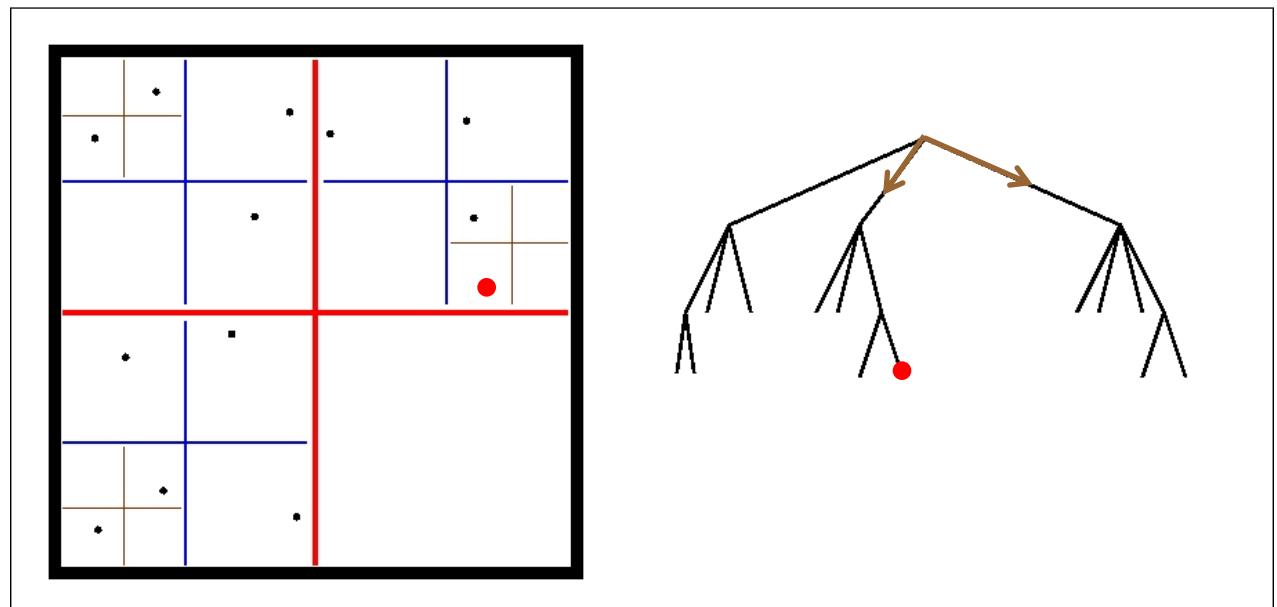


Solving for Gravity

- direct particle-particle summation (PP)

$$\vec{F}_i(\vec{r}_i) = - \sum_{i \neq j} \frac{Gm_i m_j}{(r_i - r_j)^3} (\vec{r}_i - \vec{r}_j) \quad \forall i \in N$$

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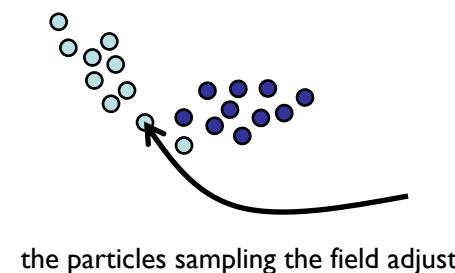


Solving for Gravity

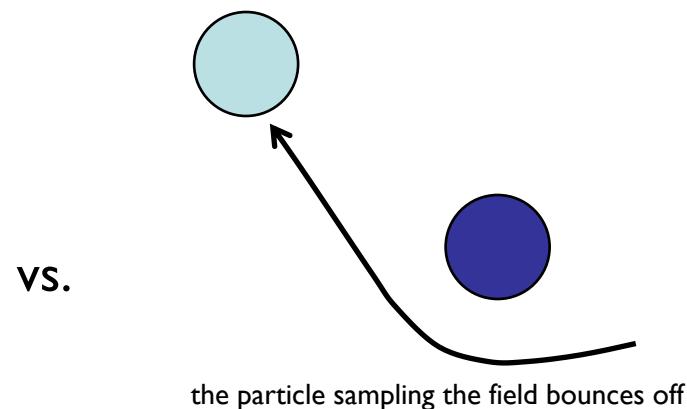
- direct particle-particle summation (PP)

$$\vec{F}_i(\vec{r}_i) = - \sum_{i \neq j} \frac{Gm_i m_j}{(r_i - r_j)^3} (\vec{r}_i - \vec{r}_j) \quad \forall i \in N$$

we use (collisionless) particles to sample $f(x, v, t)$



well sampled system



undersampled system

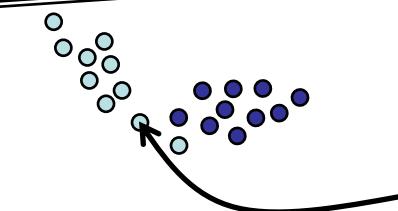
Solving for Gravity

- direct particle-particle summation (PP)

$$\vec{F}_i(\vec{r}_i) = - \sum_{i \neq j} \frac{Gm_i m_j}{(r_i - r_j)^3} (\vec{r}_i - \vec{r}_j) \quad \forall i \in N$$

we use (collisionless) particles to sample $f(x, v, t)$

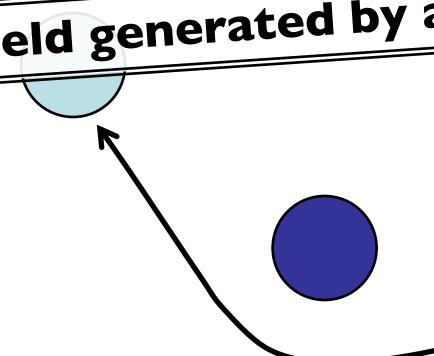
each particle should only feel the mean field generated by all particles!



the particles sampling the field adjust

well sampled system

vs.



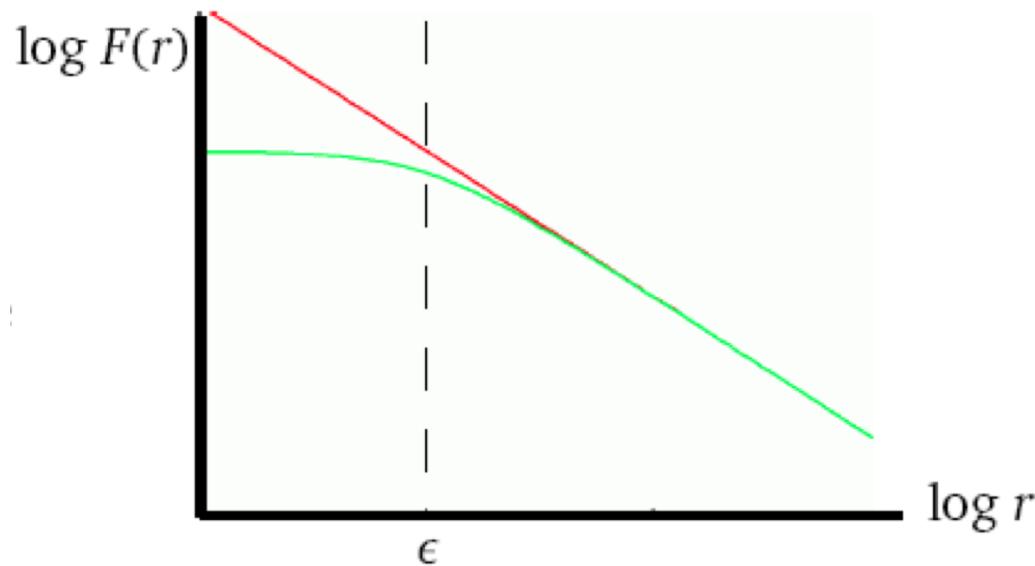
the particle sampling the field bounces off

undersampled system

Solving for Gravity

- direct particle-particle summation (PP)

$$\vec{F}_i(\vec{r}_i) = - \sum_{i \neq j} \frac{Gm_i m_j}{\left((r_i - r_j)^2 + \epsilon^2\right)^{3/2}} (\vec{r}_i - \vec{r}_j) \quad \forall i \in N$$



Solving for Gravity

- direct particle-particle summation (PP)

$$\vec{F}_i(\vec{r}_i) = - \sum_{i \neq j} \frac{Gm_i m_j}{\left((r_i - r_j)^2 + \varepsilon^2 \right)^{3/2}} (\vec{r}_i - \vec{r}_j) \quad \forall i \in N$$

“soften” the force to...

1. avoid the singularity for $r_i=r_j$
2. smooth mass density on small scales

Solving for Gravity

- direct particle-particle summation (PP)

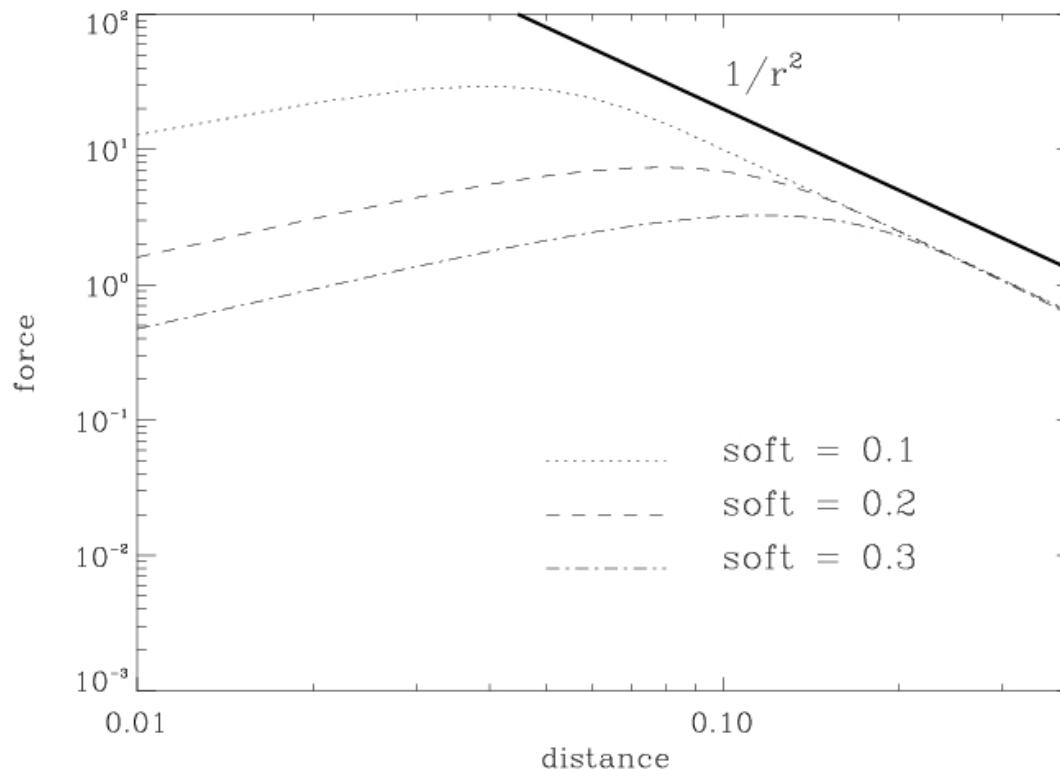
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ε determines the overall force resolution of the simulation

Solving for Gravity

- direct particle-particle summation (PP)

$$\vec{F}_i(\vec{r}_i) = - \sum_{i \neq j} \frac{Gm_i m_j}{\left((r_i - r_j)^2 + \epsilon^2\right)^{3/2}} (\vec{r}_i - \vec{r}_j) \quad \forall i \in N$$



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error budget?

Solving for Gravity

- direct particle-particle summation (PP)

$$\vec{F}_i(\vec{r}_i) = - \sum_{i \neq j} \frac{Gm_i m_j}{\left((r_i - r_j)^2 + \varepsilon^2\right)^{3/2}} (\vec{r}_i - \vec{r}_j) \quad \forall i \in N$$

- error estimate:

$$MISE = \left\langle \iiint \rho(\vec{x}) \left| \vec{F}(\vec{x}) - \vec{F}_{true}(\vec{x}) \right|^2 d^3x \right\rangle$$

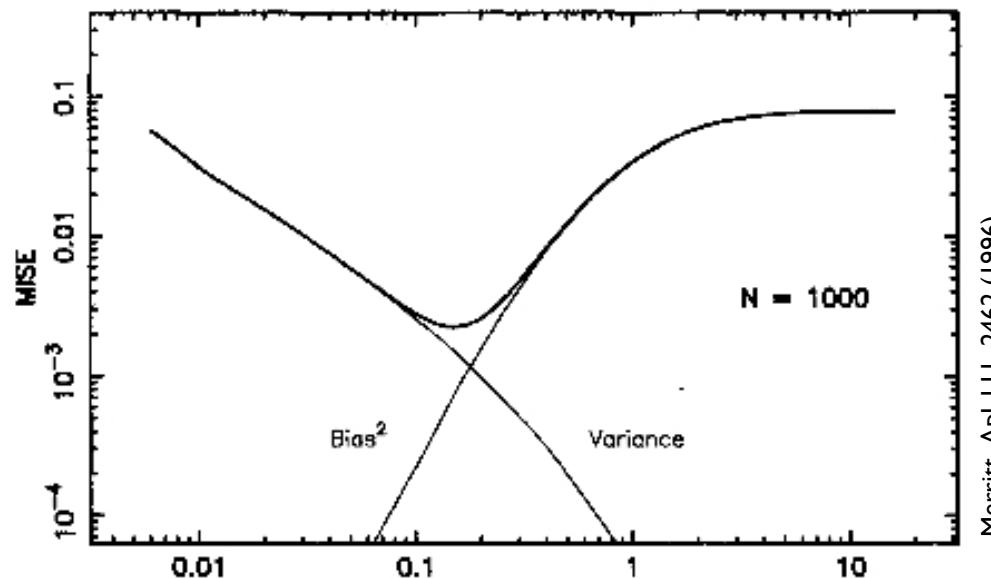
Solving for Gravity

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- error estimate:

$$MISE = \left\langle \iiint \rho(\vec{x}) |\vec{F}(\vec{x}) - \vec{F}_{true}(\vec{x})|^2 d^3x \right\rangle$$



interplay between N and ε : $N\varepsilon^3 = \text{const.}$

$\text{const.} = \left(\frac{B}{30}\right)^3$ for cosmological simulations (where B is the size of the cubical domain in 1D)

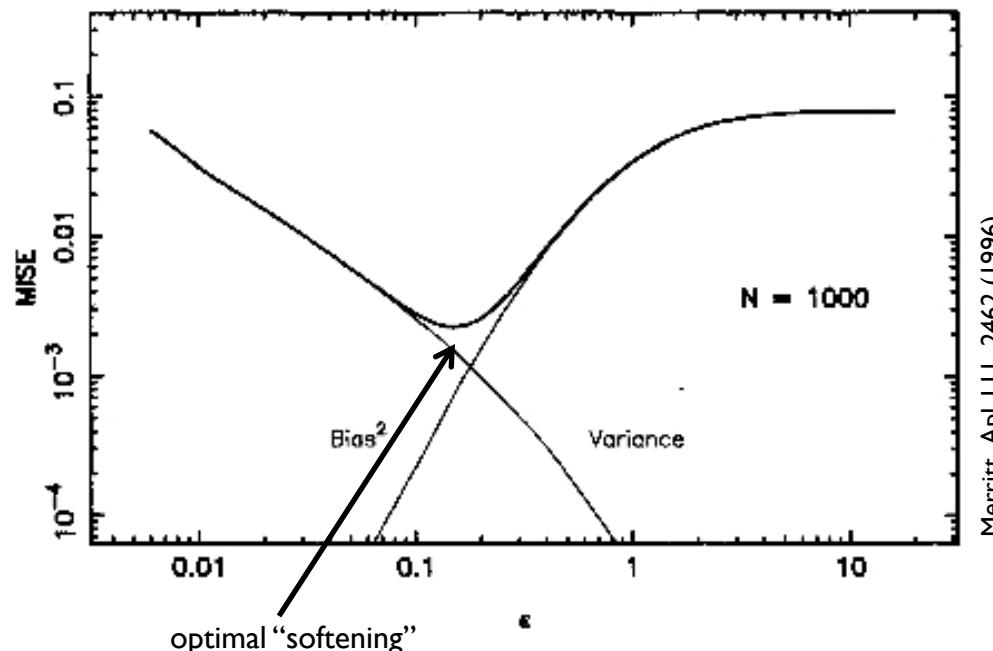
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Merritt, ApJ 111, 2462 (1996)

Computational Astrophysics

Solving for Gravity

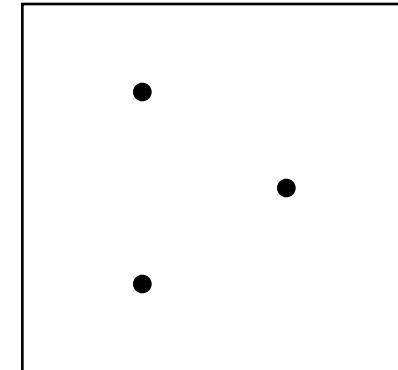
- Poisson's equation vs. Poisson's integral
- tree codes
- **PM codes**

Solving for Gravity

- numerically integrate Poisson's equation

$$\Delta \Phi(\vec{g}_{k,l,m}) = 4\pi G \rho(\vec{g}_{k,l,m})$$

$$\vec{F}(\vec{g}_{k,l,m}) = -m \nabla \Phi(\vec{g}_{k,l,m})$$

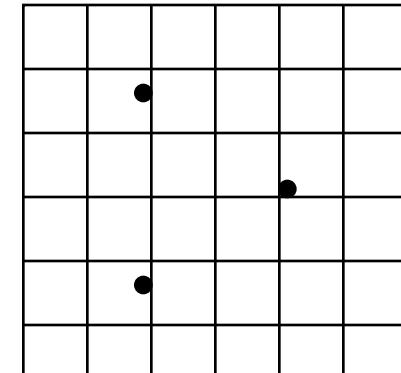


Solving for Gravity

- numerically integrate Poisson's equation

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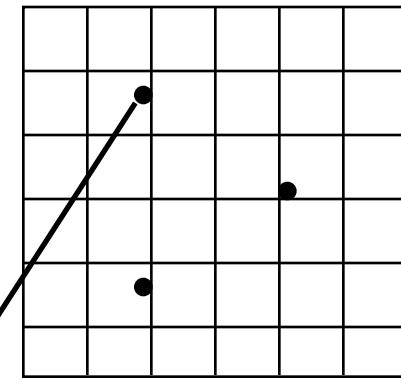
Solving for Gravity

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$$\vec{F}(\vec{g}_{k,l,m}) = -m \nabla \Phi(\vec{g}_{k,l,m})$$

I. calculate mass density on grid



$$\vec{x}_i \rightarrow \rho(\vec{g}_{k,l,m})$$

$$\forall i \in N$$

Solving for Gravity

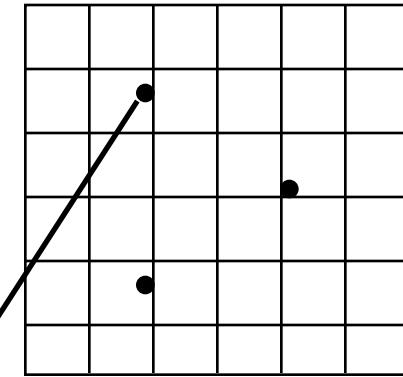
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1. calculate mass density on grid

2. solve Poisson's equation on grid



$$\vec{x}_i \rightarrow \rho(\vec{g}_{k,l,m})$$

$$\Phi(\vec{g}_{k,l,m})$$

$\forall i \in N$

Solving for Gravity

- numerically integrate Poisson's equation

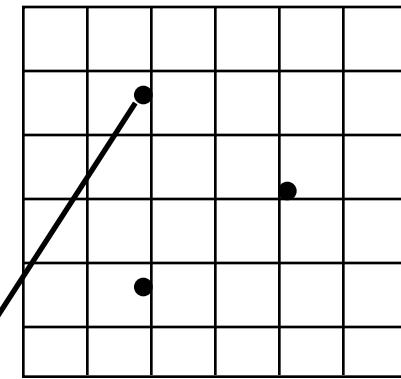
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$$\vec{F}(\vec{g}_{k,l,m}) = -m \nabla \Phi(\vec{g}_{k,l,m})$$

1. calculate mass density on grid

2. solve Poisson's equation on grid

3. differentiate potential to get forces



$$\vec{x}_i \rightarrow \rho(\vec{g}_{k,l,m})$$

$$\Phi(\vec{g}_{k,l,m})$$

$$\vec{F}(\vec{g}_{k,l,m})$$

$$\forall i \in N$$

Solving for Gravity

- numerically integrate Poisson's equation

$$\Delta\Phi(\vec{g}_{k,l,m}) = 4\pi G \rho(\vec{g}_{k,l,m})$$

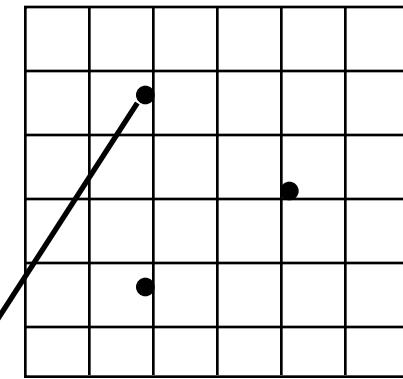
$$\vec{F}(\vec{g}_{k,l,m}) = -m \nabla \Phi(\vec{g}_{k,l,m})$$

1. calculate mass density on grid

2. solve Poisson's equation on grid

3. differentiate potential to get forces

4. interpolate forces back to particles



$$\vec{x}_i \rightarrow \rho(\vec{g}_{k,l,m}) \quad \forall i \in N$$

$$\Phi(\vec{g}_{k,l,m})$$

$$\vec{F}(\vec{g}_{k,l,m})$$

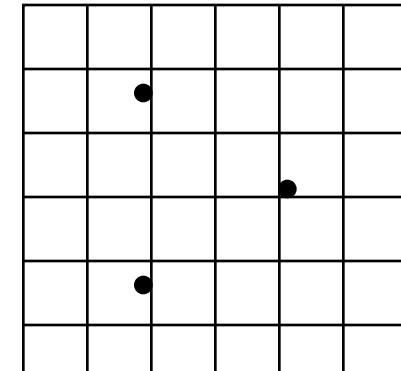
$$\vec{F}(\vec{g}_{k,l,m}) \rightarrow \vec{F}(\vec{x}_i) \quad \forall i \in N$$

Solving for Gravity

- numerically integrate Poisson's equation

$$\Delta\Phi(\vec{g}_{k,l,m}) = 4\pi G \rho(\vec{g}_{k,l,m})$$

$$\vec{F}(\vec{g}_{k,l,m}) = -m \nabla \Phi(\vec{g}_{k,l,m})$$



- calculate mass density on grid $\vec{x}_i \rightarrow \rho(\vec{g}_{k,l,m})$
- solve Poisson's equation on grid $\Phi(\vec{g}_{k,l,m})$
- differentiate potential to get forces $\vec{F}(\vec{g}_{k,l,m})$
- interpolate forces back to particles $\vec{F}(\vec{g}_{k,l,m}) \rightarrow \vec{F}(\vec{x}_i)$

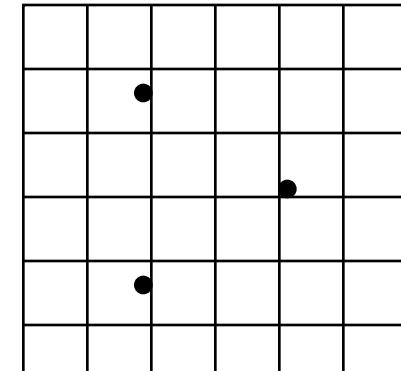
sounds like a waste of time and computer resources,
but **exceptionally fast** in practice

Solving for Gravity

- numerically integrate Poisson's equation

$$\Delta\Phi(\vec{g}_{k,l,m}) = 4\pi G \rho(\vec{g}_{k,l,m})$$

$$\vec{F}(\vec{g}_{k,l,m}) = -m \nabla \Phi(\vec{g}_{k,l,m})$$



1. calculate mass density on grid

$$\vec{x}_i \rightarrow \rho(\vec{g}_{k,l,m})$$

2. solve Poisson's equation on grid

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3. differentiate potential to get forces

$$\vec{F}(\vec{g}_{k,l,m})$$

4. interpolate forces back to particles

$$\vec{F}(\vec{g}_{k,l,m}) \rightarrow \vec{F}(\vec{x}_i)$$

Solving for Gravity

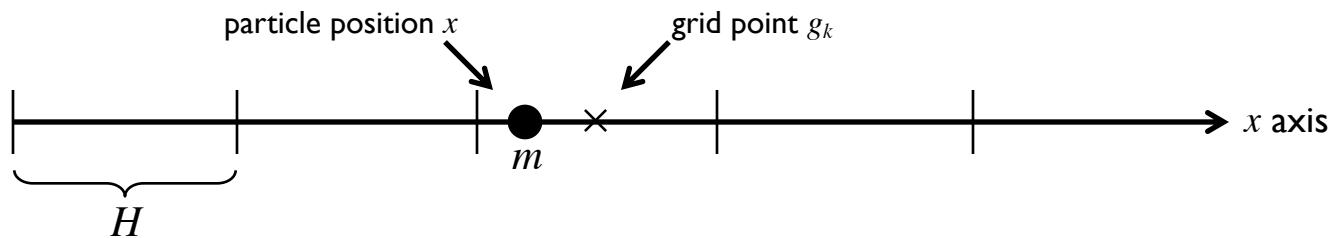
- density assignment schemes

$$\vec{x}_i \rightarrow \rho(\vec{g}_{k,l,m})$$

example: 1 particle on 1 dimensional grid

$$M(g_k) = mW(d) \quad d = |x - g_k|$$

$$\rho(g_k) = \frac{M(g_k)}{H}$$



Solving for Gravity

- density assignment schemes

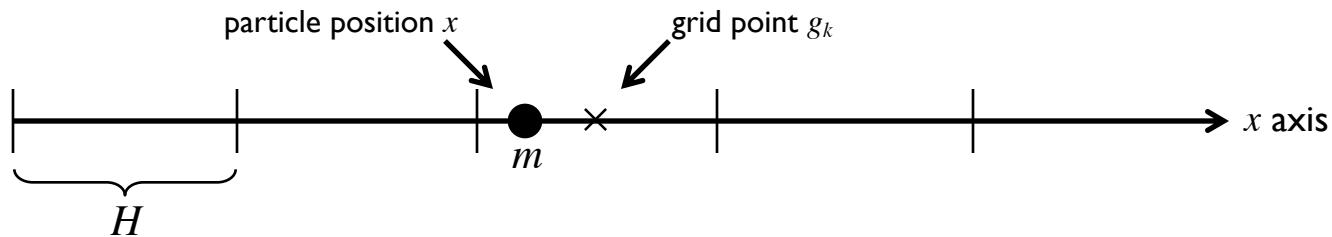
$$\vec{x}_i \rightarrow \rho(\vec{g}_{k,l,m})$$

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$$M(g_k) = m W(d) \quad d = |x - g_k|$$

mass assignment function

$$\rho(g_k) = \frac{M(g_k)}{H}$$



Solving for Gravity

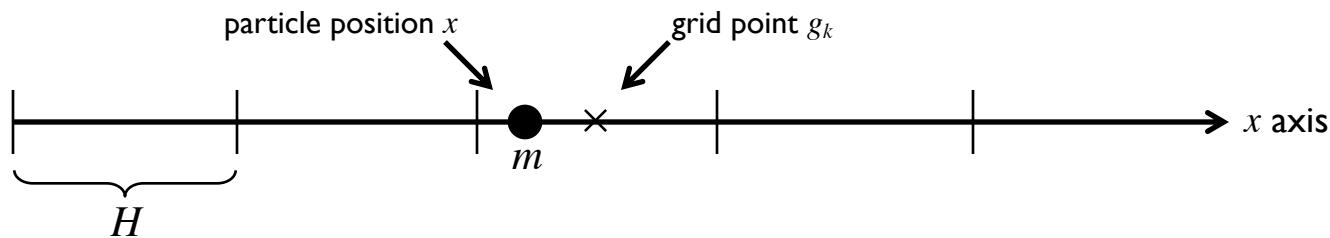
- density assignment schemes

$$\vec{x}_i \rightarrow \rho(\vec{g}_{k,l,m})$$

example: 1 particle on 1 dimensional grid

- hierarchy of mass assignment schemes:

- Nearest-Grid-Point NGP
- Could-In-Cell CIC
- Triangular-Shaped Cloud TSC
- ...



Solving for Gravity

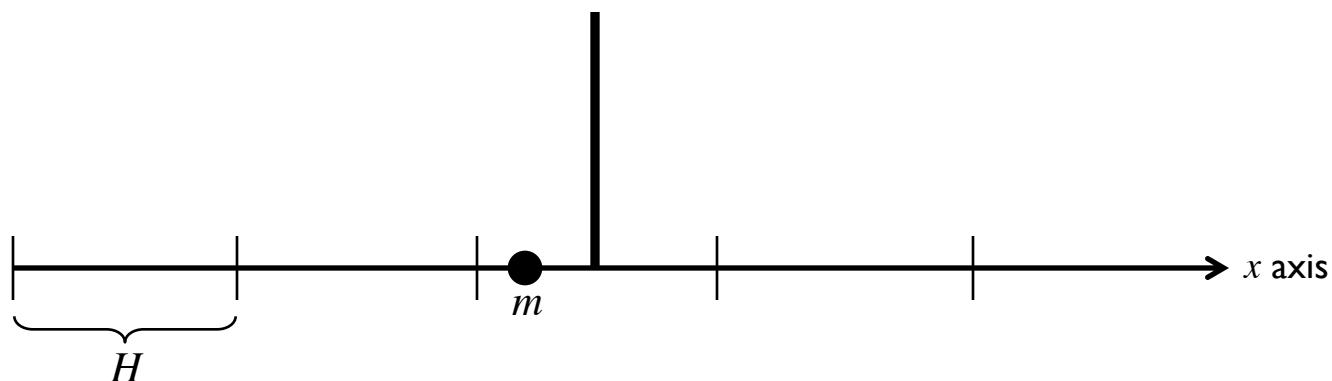
- density assignment schemes

$$\vec{x}_i \rightarrow \rho(\vec{g}_{k,l,m})$$

Nearest-Grid-Point (NGP):

mass assignment function:

$$W(d) = \begin{cases} 1 & d \leq H/2 \\ 0 & \text{otherwise} \end{cases}$$



Solving for Gravity

- density assignment schemes

$$\vec{x}_i \rightarrow \rho(\vec{g}_{k,l,m})$$

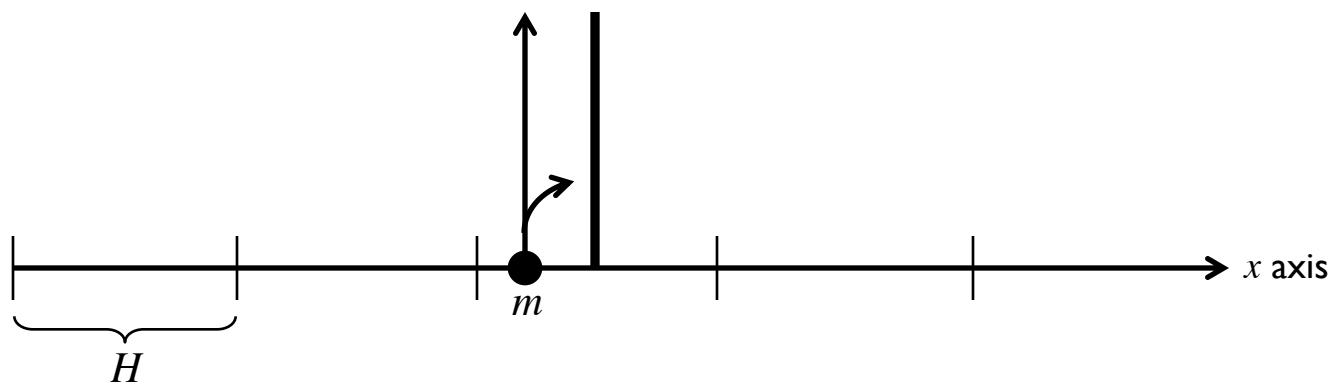
Nearest-Grid-Point (NGP):

particle shape:

$$S(x) = \delta(x)$$

mass assignment function:

$$W(d) = \begin{cases} 1 & d \leq H/2 \\ 0 & \text{otherwise} \end{cases}$$



Solving for Gravity

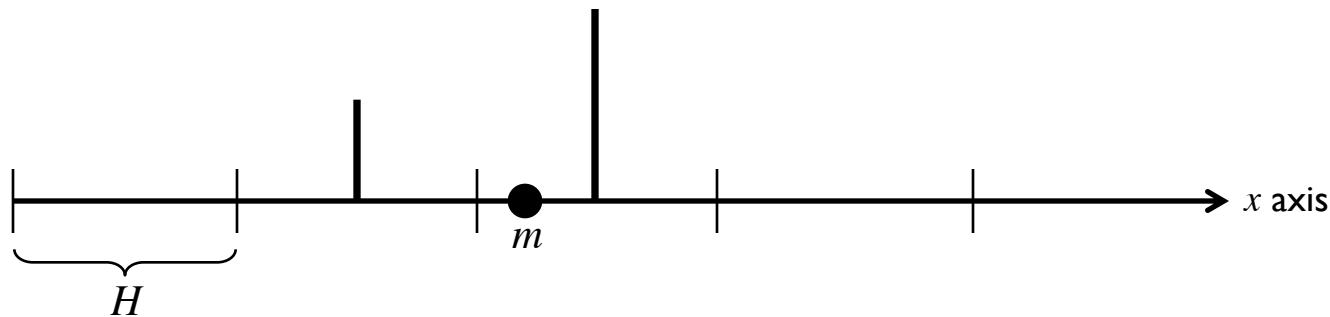
- density assignment schemes

$$\vec{x}_i \rightarrow \rho(\vec{g}_{k,l,m})$$

Cloud-In-Cell (CIC):

mass assignment function:

$$W(d) = \begin{cases} 1 - \frac{d}{H} & d \leq H \\ 0 & \text{otherwise} \end{cases}$$



Solving for Gravity

- density assignment schemes

$$\vec{x}_i \rightarrow \rho(\vec{g}_{k,l,m})$$

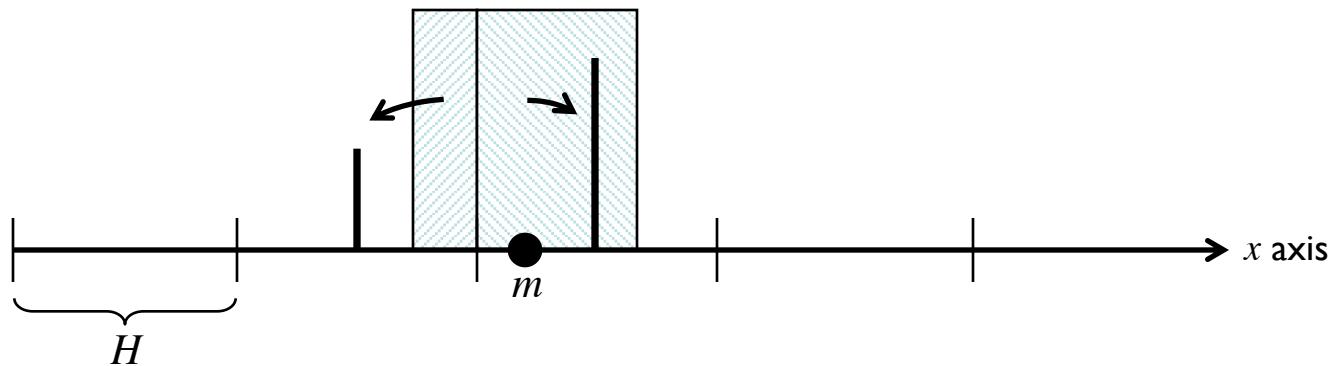
Cloud-In-Cell (CIC):

particle shape:

$$S(x) = \begin{cases} 1 & |x| \leq H/2 \\ 0 & \text{otherwise} \end{cases}$$

mass assignment function:

$$W(d) = \begin{cases} 1 - \frac{d}{H} & d \leq H \\ 0 & \text{otherwise} \end{cases}$$



Solving for Gravity

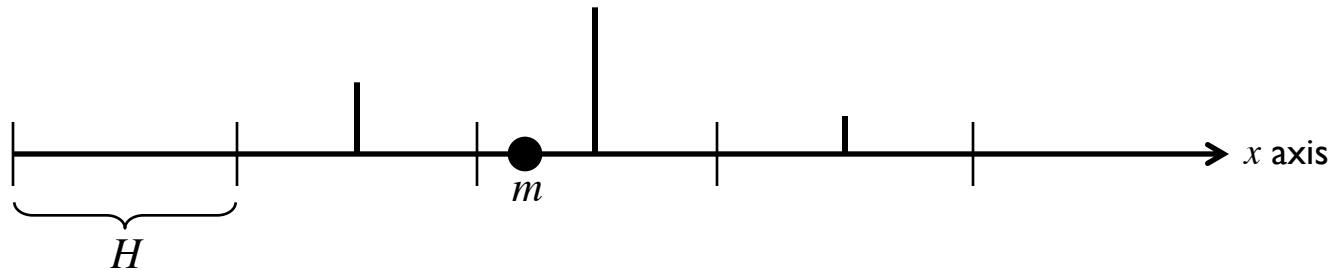
- density assignment schemes

$$\vec{x}_i \rightarrow \rho(\vec{g}_{k,l,m})$$

Triangular-Shaped-Cloud (TSC):

mass assignment function:

$$W(d) = \begin{cases} \frac{3}{4} - \left(\frac{d}{H}\right)^2 & d \leq \frac{H}{2} \\ \frac{1}{2} \left(\frac{3}{2} - \frac{d}{H} \right)^2 & \frac{H}{2} \leq d \leq \frac{3H}{2} \\ 0 & \text{otherwise} \end{cases}$$



Solving for Gravity

- density assignment schemes

$$\vec{x}_i \rightarrow \rho(\vec{g}_{k,l,m})$$

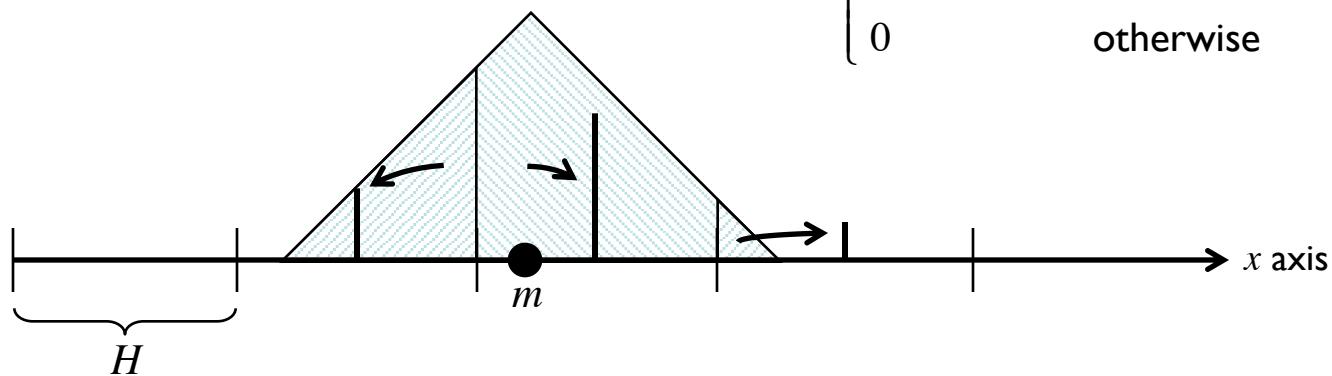
Triangular-Shaped-Cloud (TSC):

particle shape:

$$S(x) = \begin{cases} 1 - \frac{|x|}{H} & |x| \leq H \\ 0 & \text{otherwise} \end{cases}$$

mass assignment function:

$$W(d) = \begin{cases} \frac{3}{4} - \left(\frac{d}{H}\right)^2 & d \leq \frac{H}{2} \\ \frac{1}{2} \left(3 - \frac{d}{H}\right)^2 & \frac{H}{2} \leq d \leq \frac{3H}{2} \\ 0 & \text{otherwise} \end{cases}$$



Solving for Gravity

- density assignment schemes

$$\vec{x}_i \rightarrow \rho(\vec{g}_{k,l,m})$$

N particles on 3 dimensional grid

$$\vec{d} = \vec{x}_i - \vec{g}_{k,l,m}$$

$$M(\vec{g}_{k,l,m}) = \sum_{i=1}^N m_i W(|d_x|)W(|d_y|)W(|d_z|)$$

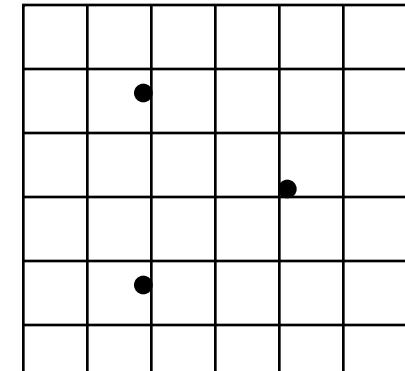
$$\rho(\vec{g}_{k,l,m}) = \frac{M(\vec{g}_{k,l,m})}{H^3}$$

Solving for Gravity

- numerically integrate Poisson's equation

$$\Delta\Phi(\vec{g}_{k,l,m}) = 4\pi G \rho(\vec{g}_{k,l,m})$$

$$\vec{F}(\vec{g}_{k,l,m}) = -m \nabla \Phi(\vec{g}_{k,l,m})$$



I. calculate mass density on grid

$$\vec{x}_i \rightarrow \rho(\vec{g}_{k,l,m})$$

2. solve Poisson's equation on grid

$$\Phi(\vec{g}_{k,l,m})$$

3. differentiate potential to get forces

$$\vec{F}(\vec{g}_{k,l,m})$$

4. interpolate forces back to particles

$$\vec{F}(\vec{g}_{k,l,m}) \rightarrow \vec{F}(\vec{x}_i)$$

Solving for Gravity

- numerically integrate Poisson's equation

$$\Delta\Phi_{k,l,m} = \rho_{k,l,m}$$

- relaxation technique: applicable and usable for **any** differential equation
- FTT technique: only applicable and usable for **linear** differential equation

Solving for Gravity

- numerically integrate Poisson's equation

$$\Delta\Phi_{k,l,m} = \rho_{k,l,m}$$

- relaxation technique: applicable and usable for **any** differential equation
- **FTT technique:** only applicable and usable for **linear** differential equation

Solving for Gravity

- numerically integrate Poisson's equation

$$\Delta \Phi_{k,l,m} = \rho_{k,l,m}$$

- Green's function method:
 - solve differential equation by Fourier transformation
 - applicable and usable for **linear** differential equations

Solving for Gravity

- numerically integrate Poisson's equation *fast fourier transform method*
 - Green's function method

$$\Delta \Phi = \rho$$
$$\Phi(\vec{x}) = \iiint G(\vec{x} - \vec{x}') \rho(\vec{x}') d^3x' ; \quad G(\vec{x}) = \frac{1}{4\pi x}$$

Solving for Gravity

- numerically integrate Poisson's equation *fast fourier transform method*
 - Green's function method

$$\Delta \Phi = \rho$$

$$\Phi(\vec{x}) = \iiint G(\vec{x} - \vec{x}') \rho(\vec{x}') d^3x' ; \quad G(\vec{x}) = \frac{1}{4\pi x}$$

$$\Phi = \rho \otimes G$$



FFT \rightarrow convolution becomes multiplication

$$\hat{\Phi} = \hat{\rho} \otimes \hat{G} \quad (\text{with } \hat{G} = -\frac{1}{k^2} \text{ for Poisson's equation })$$

Solving for Gravity

- numerically integrate Poisson's equation *fast fourier transform method*
 - Green's function method

$$\Delta\Phi = \rho \quad \Phi(\vec{x}) = \iiint G(\vec{x} - \vec{x}') \rho(\vec{x}') d^3x' ; \quad G(\vec{x}) = \frac{1}{4\pi x}$$

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FFT \rightarrow convolution becomes multiplication

$$\hat{\Phi} = \hat{\rho} \otimes \hat{G} \quad (\text{with } \hat{G} = -\frac{1}{k^2} \text{ for Poisson's equation })$$



FFT⁻¹

$$\Phi$$

(FFT demands a *regular grid* though!)

Solving for Gravity

- numerically integrate Poisson's equation

$$\Delta\Phi_{k,l,m} = \rho_{k,l,m}$$

- **relaxation technique:** applicable and usable for *any* differential equation
 - FTT technique: only applicable and usable for *linear* differential equation

Solving for Gravity

- numerically integrate Poisson's equation *relaxation technique*

obtain iterative solver by discretizing differential equation

$$\Delta\Phi_{k,l,m} = \rho_{k,l,m}$$

Solving for Gravity

- numerically integrate Poisson's equation relaxation technique

obtain iterative solver by discretizing differential equation

$$\Delta\Phi_{k,l,m} = \rho_{k,l,m}$$

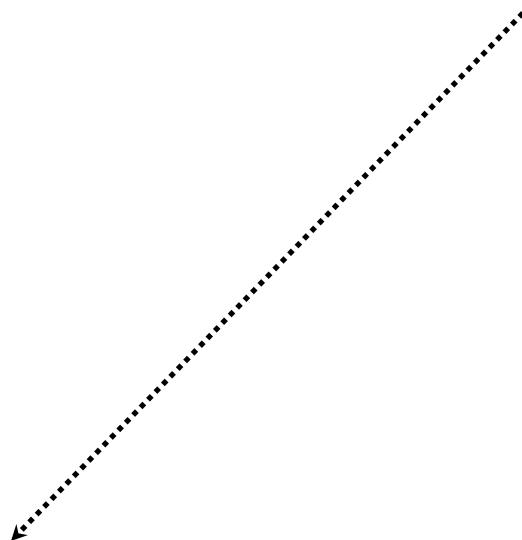
$$\begin{aligned}
 \Delta\Phi_{k,l,m} &= \nabla \cdot \nabla\Phi_{k,l,m} \\
 &= \left(\frac{\partial}{\partial x}, \frac{\partial}{\partial y}, \frac{\partial}{\partial z} \right) \cdot \left(\frac{\partial\Phi_{k,l,m}}{\partial x}, \frac{\partial\Phi_{k,l,m}}{\partial y}, \frac{\partial\Phi_{k,l,m}}{\partial z} \right) \\
 &= \frac{1}{H} \left(\frac{\partial}{\partial x}, \frac{\partial}{\partial y}, \frac{\partial}{\partial z} \right) \cdot \left(\Phi_{k+\frac{1}{2},l,m} - \Phi_{k-\frac{1}{2},l,m}, \Phi_{k,l+\frac{1}{2},m} - \Phi_{k,l-\frac{1}{2},m}, \Phi_{k,l,m+\frac{1}{2}} - \Phi_{k,l,m-\frac{1}{2}} \right) \\
 &= \frac{1}{H} \left(\frac{\partial\Phi_{k+\frac{1}{2},l,m}}{\partial x} - \frac{\partial\Phi_{k-\frac{1}{2},l,m}}{\partial x} + \frac{\partial\Phi_{k,l+\frac{1}{2},m}}{\partial y} - \frac{\partial\Phi_{k,l-\frac{1}{2},m}}{\partial y} + \frac{\partial\Phi_{k,l,m+\frac{1}{2}}}{\partial z} - \frac{\partial\Phi_{k,l,m-\frac{1}{2}}}{\partial z} \right) \\
 &= \frac{1}{H^2} (\Phi_{k+1,l,m} - 2\Phi_{k,l,m} + \Phi_{k-1,l,m} + \Phi_{k,l+1,m} - 2\Phi_{k,l,m} + \Phi_{k,l-1,m} + \Phi_{k,l,m+1} - 2\Phi_{k,l,m} + \Phi_{k,l,m-1})
 \end{aligned}$$

Solving for Gravity

- numerically integrate Poisson's equation relaxation technique

obtain iterative solver by discretizing differential equation

$$\Delta\Phi_{k,l,m} = \rho_{k,l,m}$$



discretized Poisson's equation

$$\Phi_{k,l,m} = \frac{1}{6}(\Phi_{k+1,l,m} + \Phi_{k-1,l,m} + \Phi_{k,l+1,m} + \Phi_{k,l-1,m} + \Phi_{k,l,m+1} + \Phi_{k,l,m-1} - \rho_{k,l,m} H^2)$$

Solving for Gravity

- numerically integrate Poisson's equation relaxation technique

obtain iterative solver by discretizing differential equation

$$\Delta\Phi_{k,l,m} = \rho_{k,l,m}$$

iterative solution: $\Phi_{k,l,m}^i \rightarrow \Phi_{k,l,m}^{i+1}$

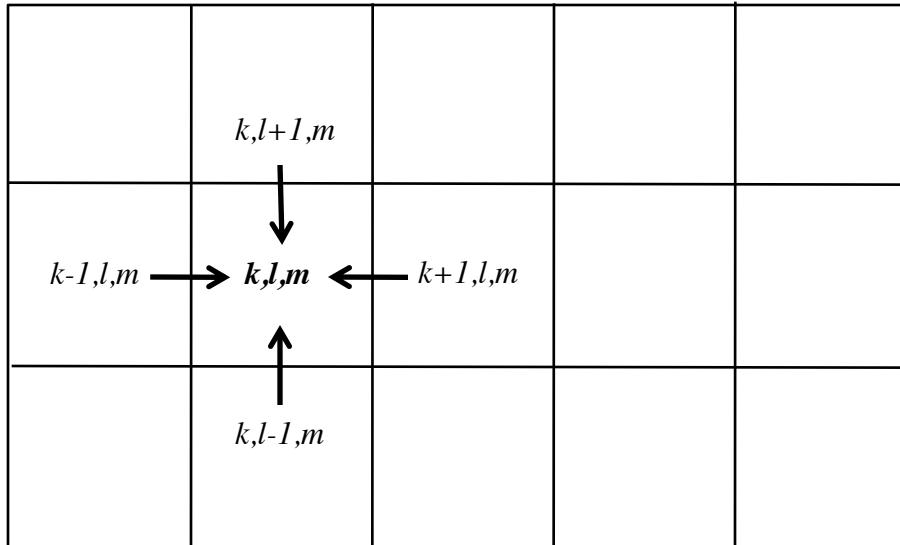
discretized Poisson's equation

$$\Phi_{k,l,m}^{i+1} = \frac{1}{6}(\Phi_{k+1,l,m}^i + \Phi_{k-1,l,m}^i + \Phi_{k,l+1,m}^i + \Phi_{k,l-1,m}^i + \Phi_{k,l,m+1}^i + \Phi_{k,l,m-1}^i - \rho_{k,l,m} H^2)$$

Solving for Gravity

- numerically integrate Poisson's equation relaxation technique

obtain iterative solver by discretizing differential equation



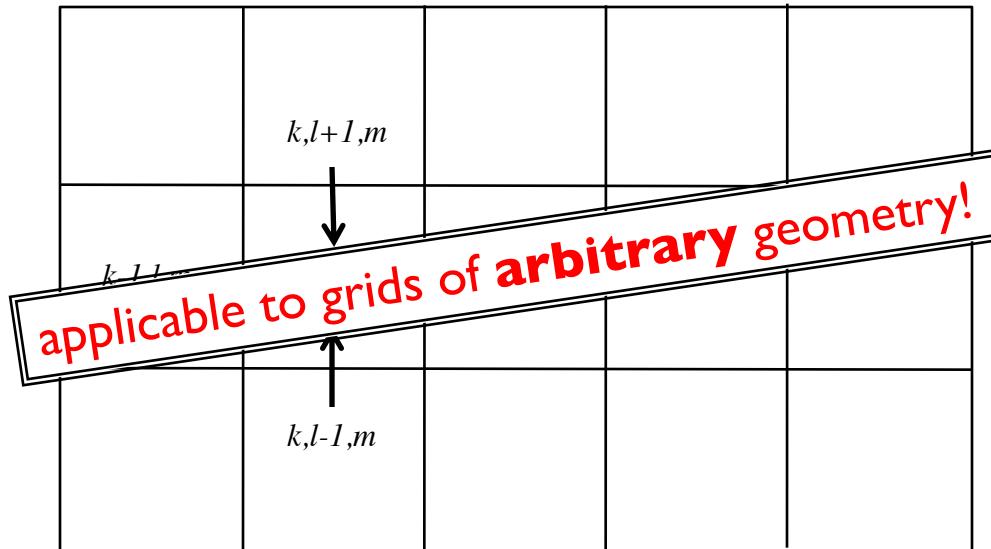
discretized Poisson's equation

$$\Phi_{k,l,m}^{i+1} = \frac{1}{6} (\Phi_{k+1,l,m}^i + \Phi_{k-1,l,m}^i + \Phi_{k,l+1,m}^i + \Phi_{k,l-1,m}^i + \Phi_{k,l,m+1}^i + \Phi_{k,l,m-1}^i - \rho_{k,l,m} H^2)$$

Solving for Gravity

- numerically integrate Poisson's equation relaxation technique

obtain iterative solver by discretizing differential equation



discretized Poisson's equation

$$\Phi_{k,l,m}^{i+1} = \frac{1}{6} (\Phi_{k+1,l,m}^i + \Phi_{k-1,l,m}^i + \Phi_{k,l+1,m}^i + \Phi_{k,l-1,m}^i + \Phi_{k,l,m+1}^i + \Phi_{k,l,m-1}^i - \rho_{k,l,m} H^2)$$

Solving for Gravity

- numerically integrate Poisson's equation

accuracy of either

relaxation

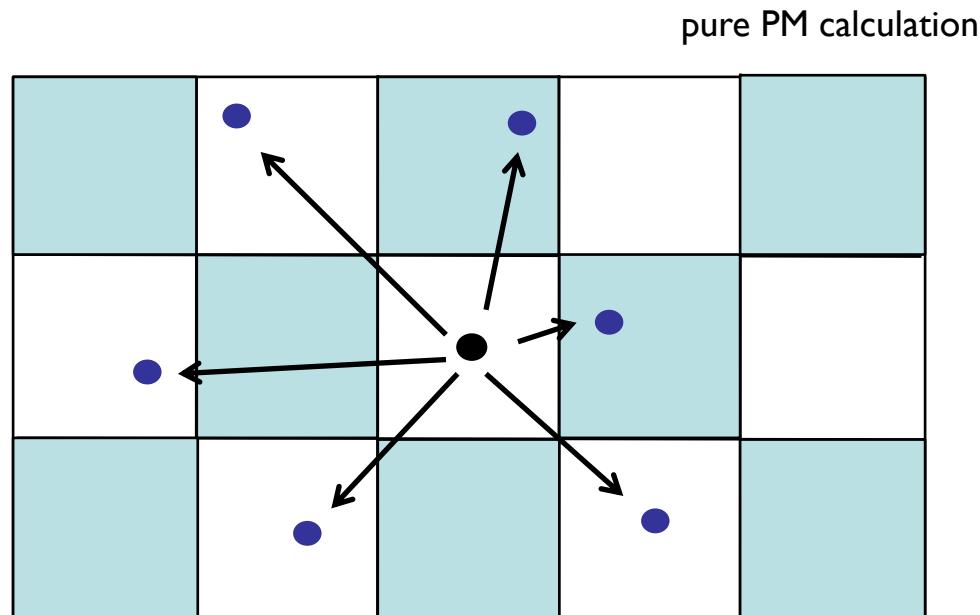
or

FFT method

to solve Poisson's equation?

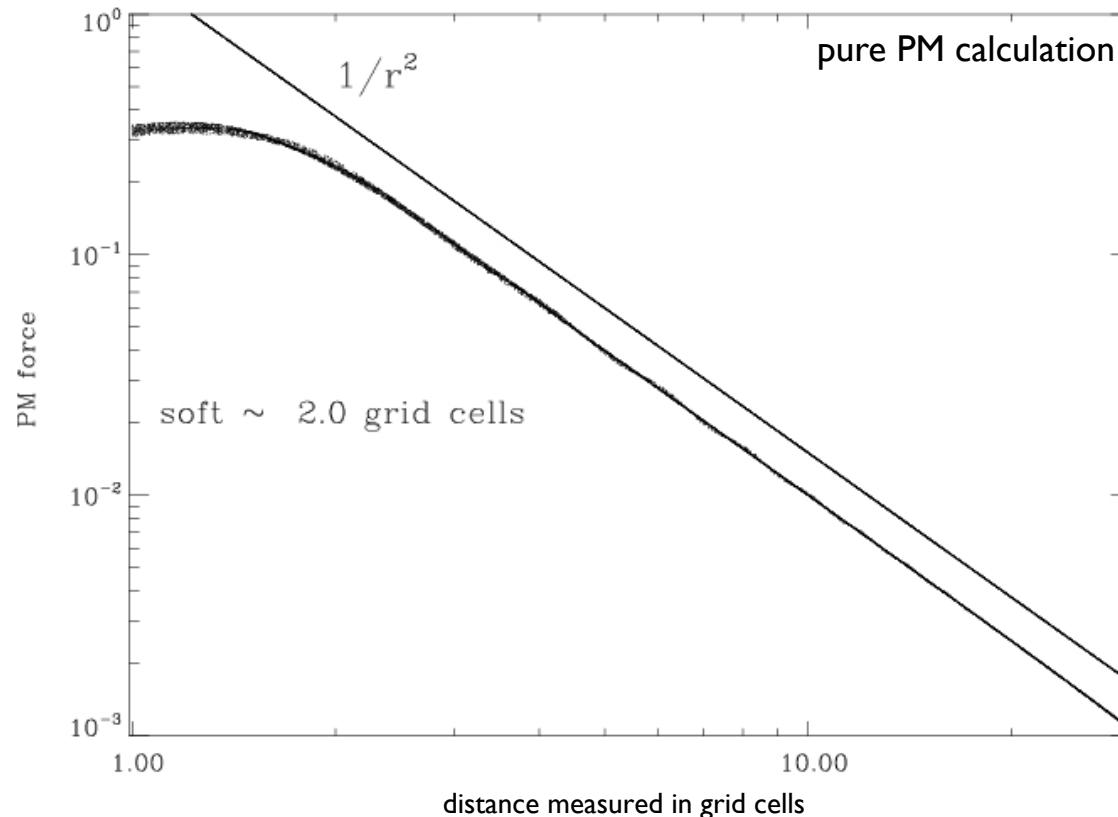
Solving for Gravity

- numerically integrate Poisson's equation



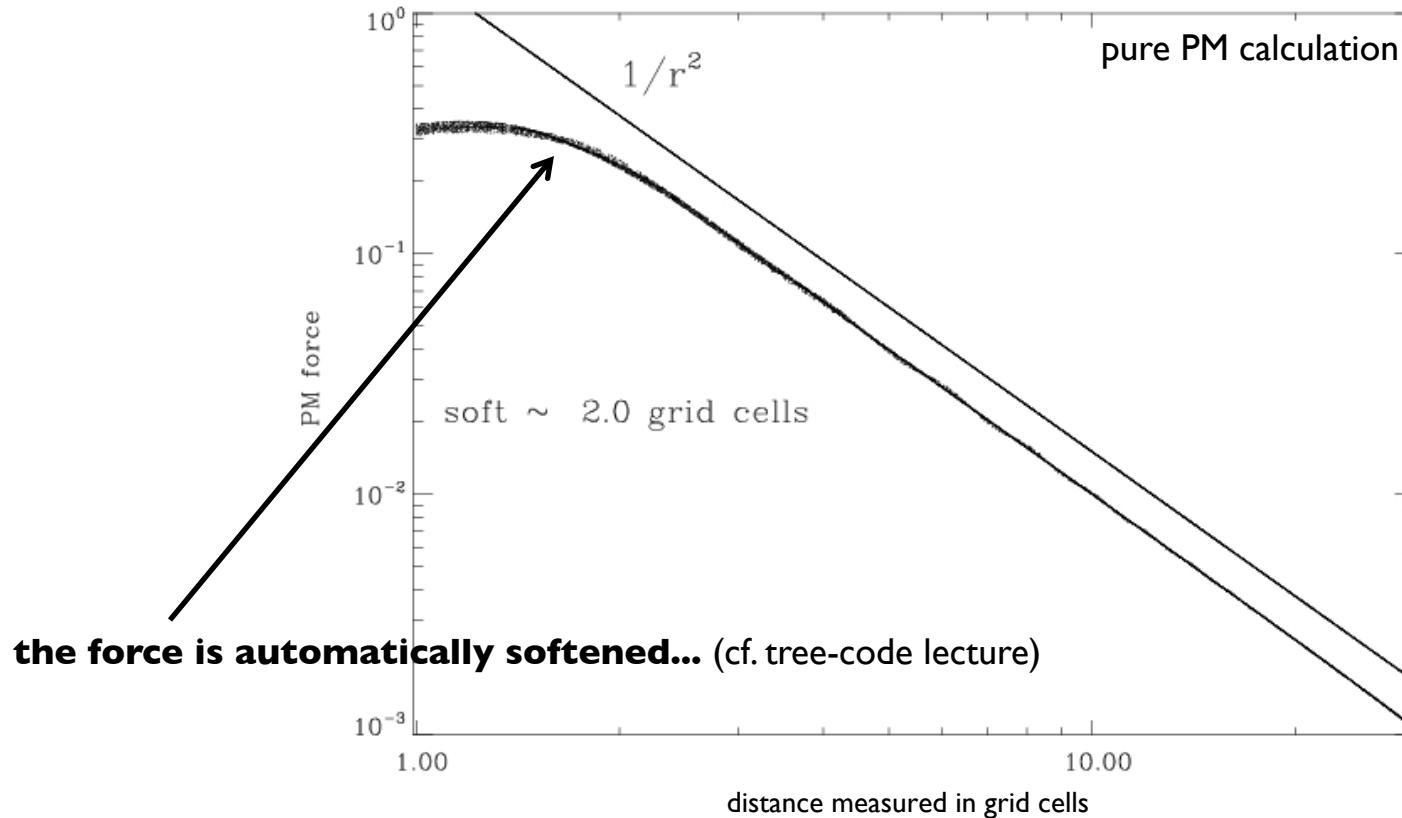
Solving for Gravity

- numerically integrate Poisson's equation



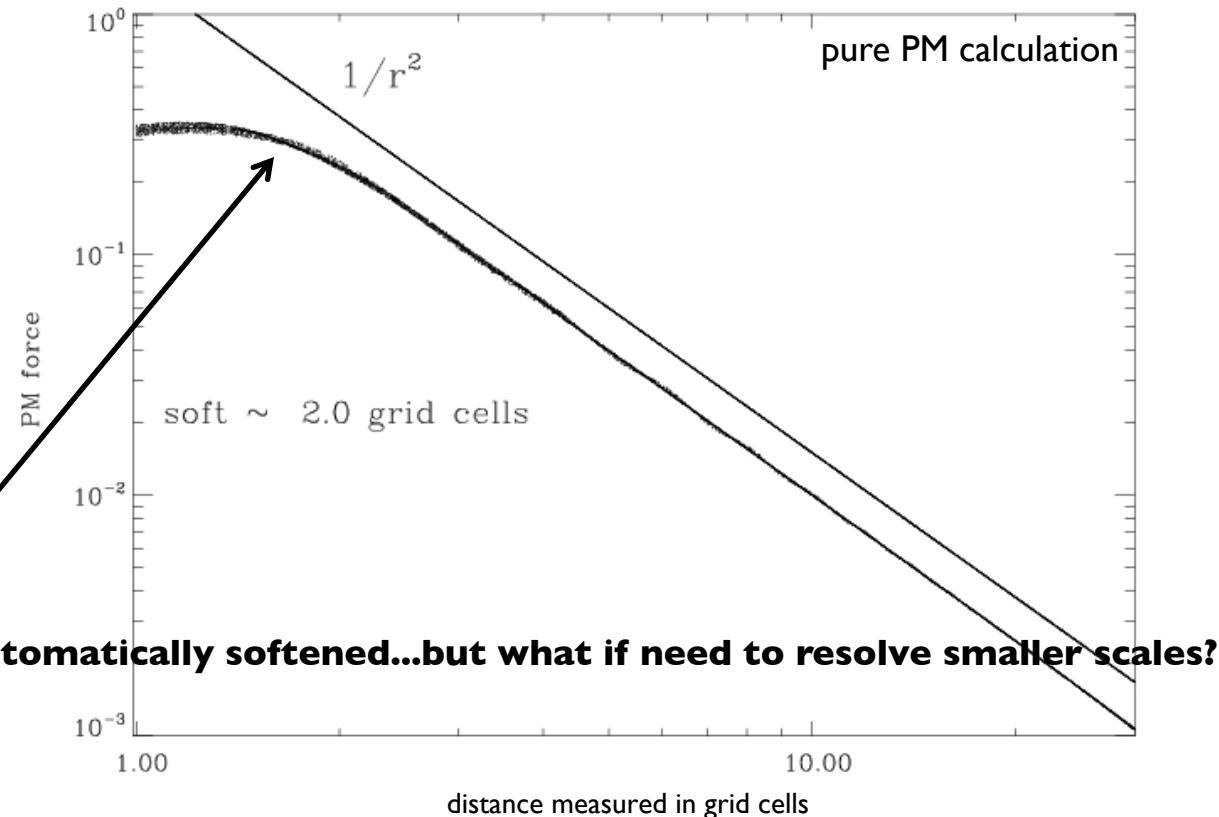
Solving for Gravity

- numerically integrate Poisson's equation



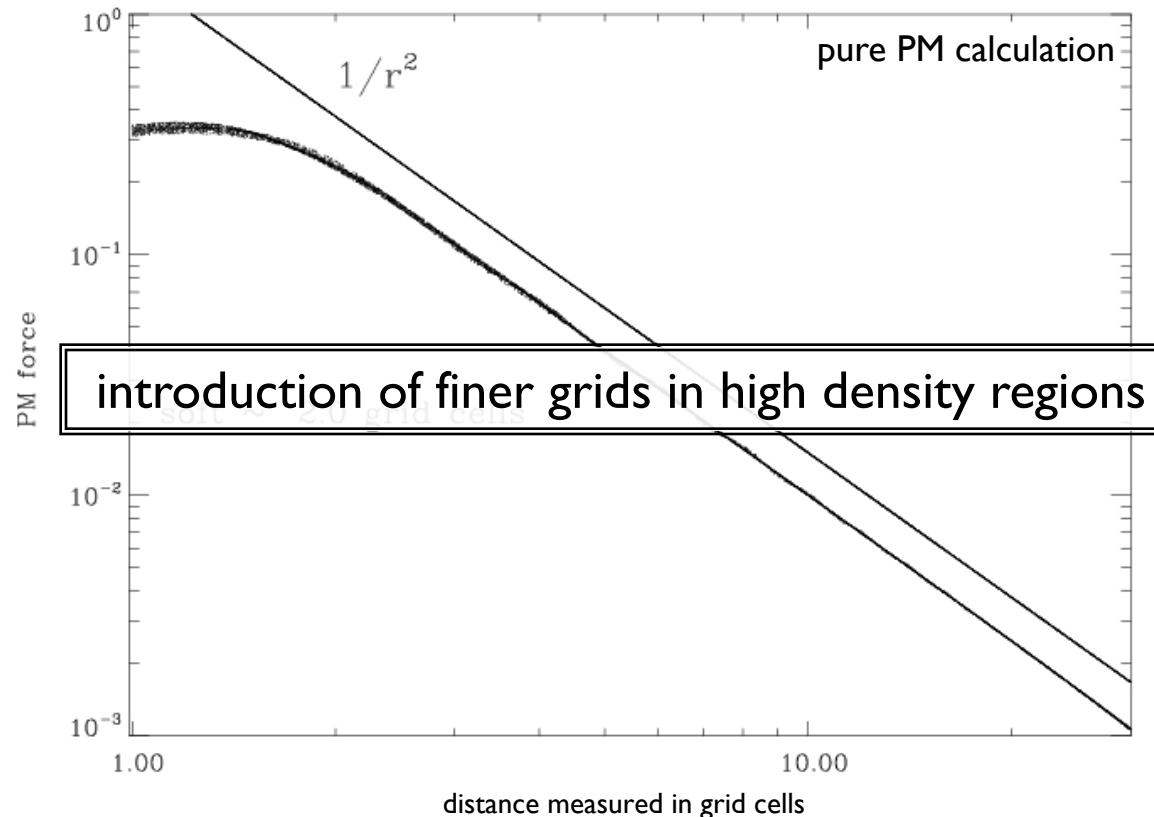
Solving for Gravity

- numerically integrate Poisson's equation



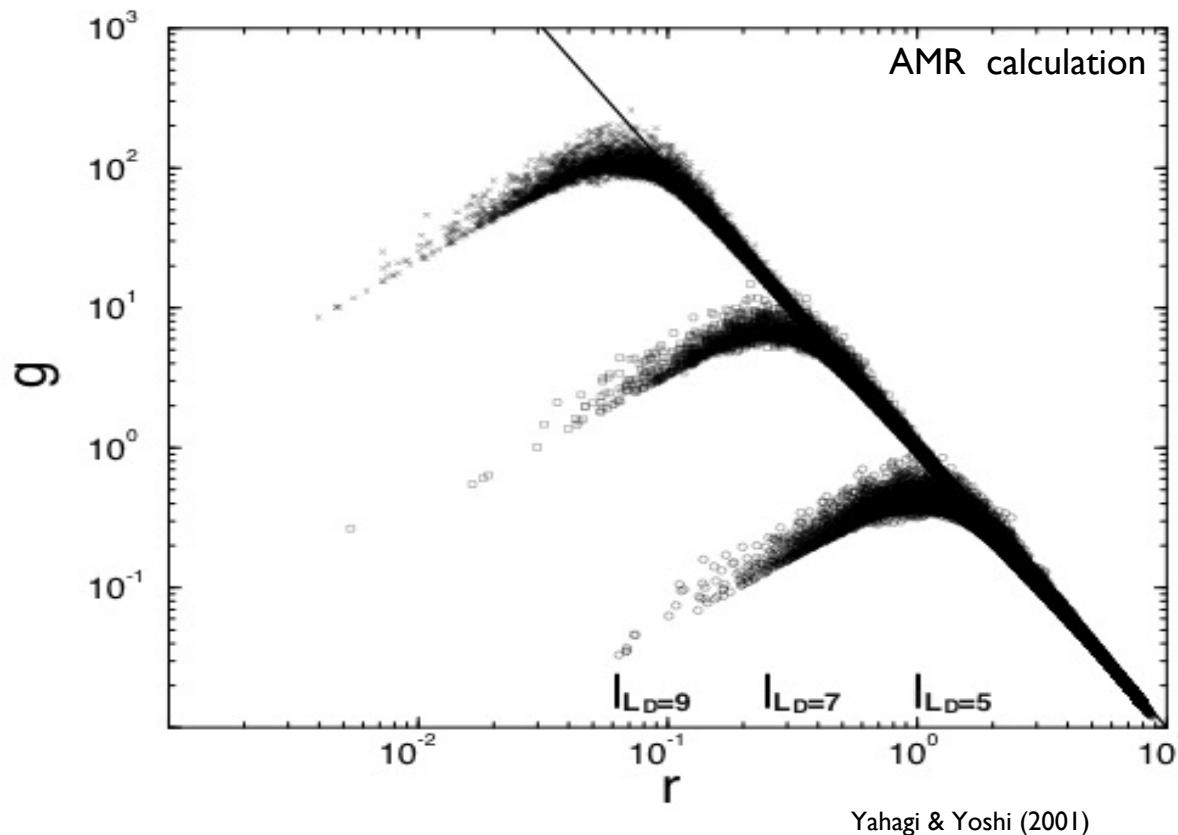
Solving for Gravity

- numerically integrate Poisson's equation



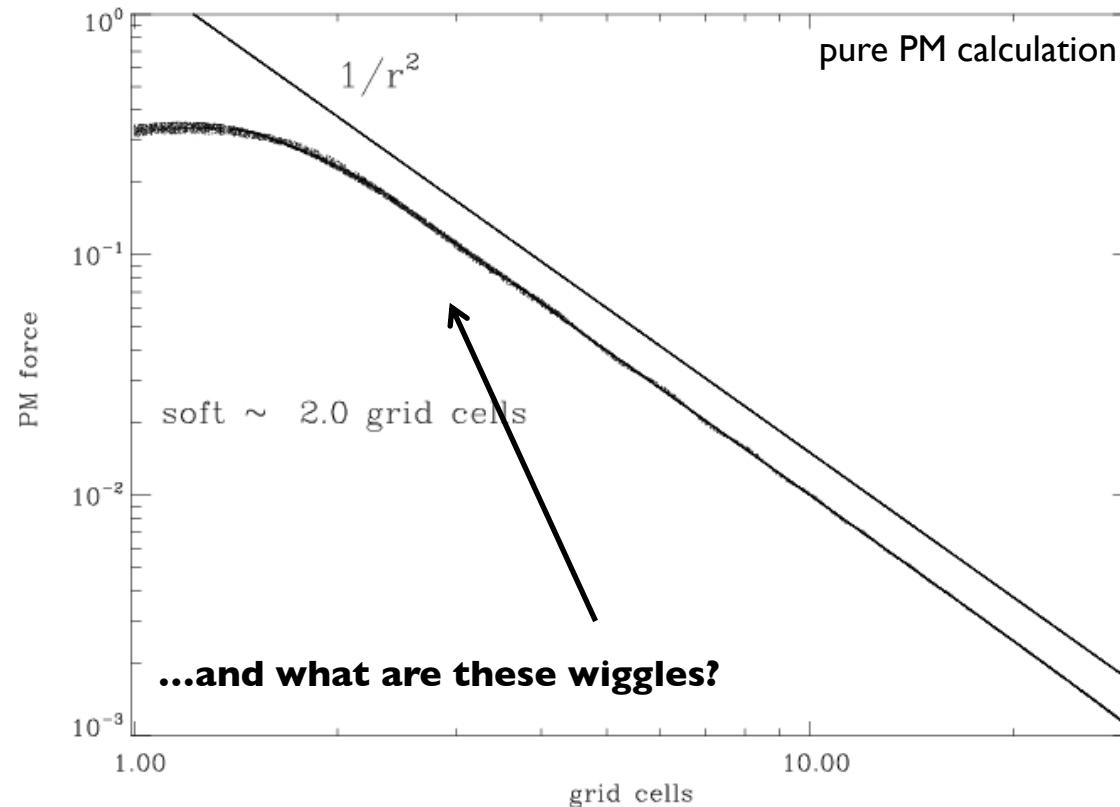
Solving for Gravity

- numerically integrate Poisson's equation



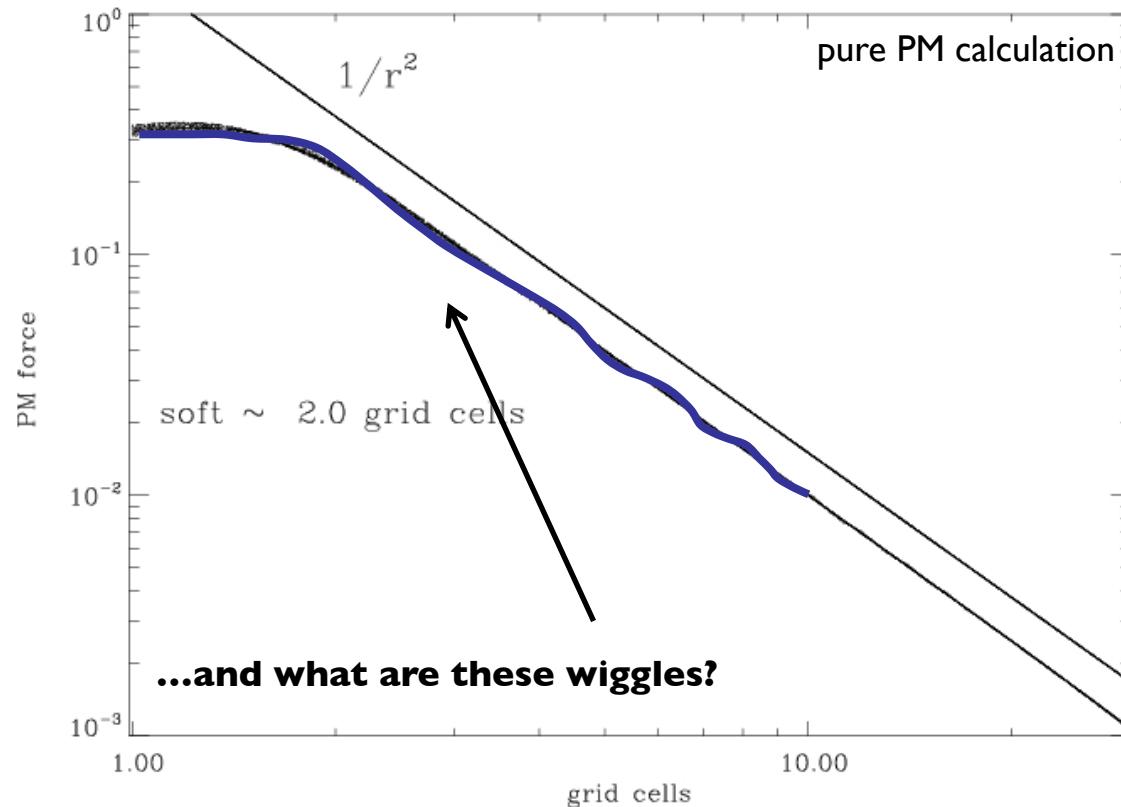
Solving for Gravity

- numerically integrate Poisson's equation



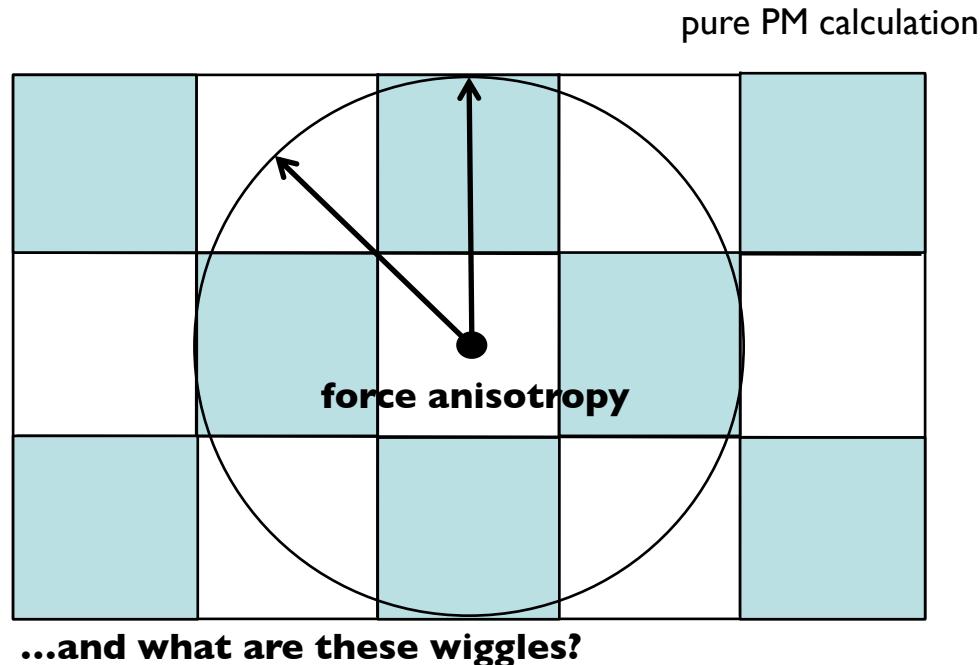
Solving for Gravity

- numerically integrate Poisson's equation



Solving for Gravity

- numerically integrate Poisson's equation

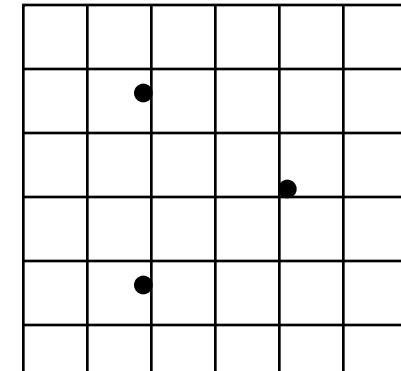


Solving for Gravity

- numerically integrate Poisson's equation

$$\Delta\Phi(\vec{g}_{k,l,m}) = 4\pi G \rho(\vec{g}_{k,l,m})$$

$$\vec{F}(\vec{g}_{k,l,m}) = -m \nabla \Phi(\vec{g}_{k,l,m})$$



1. calculate mass density on grid $\vec{x}_i \rightarrow \rho(\vec{g}_{k,l,m})$
2. solve Poisson's equation on grid $\Phi(\vec{g}_{k,l,m})$
3. differentiate potential to get forces $\vec{F}(\vec{g}_{k,l,m})$
4. interpolate forces back to particles $\vec{F}(\vec{g}_{k,l,m}) \rightarrow \vec{F}(\vec{x}_i)$

Solving for Gravity

- obtaining the forces

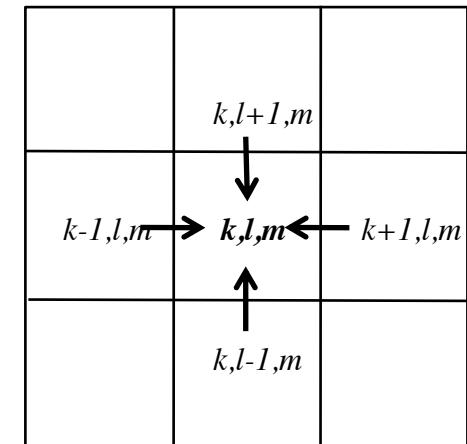
$$\vec{F}(\vec{g}_{k,l,m}) = -m \nabla \Phi(\vec{g}_{k,l,m})$$



$$F_x(\vec{g}_{k,l,m}) = -m \frac{\Phi(\vec{g}_{k+1,l,m}) - \Phi(\vec{g}_{k-1,l,m})}{2H}$$

$$F_y(\vec{g}_{k,l,m}) = -m \frac{\Phi(\vec{g}_{k,l+1,m}) - \Phi(\vec{g}_{k,l-1,m})}{2H}$$

$$F_z(\vec{g}_{k,l,m}) = -m \frac{\Phi(\vec{g}_{k,l,m+1}) - \Phi(\vec{g}_{k,l,m-1})}{2H}$$



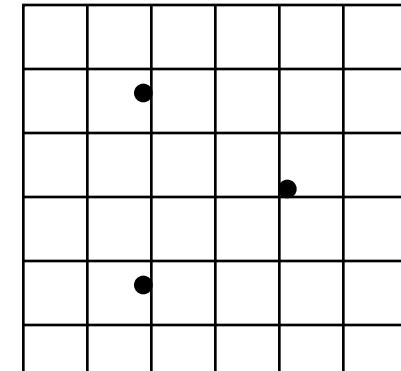
H = (current) grid spacing

Solving for Gravity

- numerically integrate Poisson's equation

$$\Delta\Phi(\vec{g}_{k,l,m}) = 4\pi G \rho(\vec{g}_{k,l,m})$$

$$\vec{F}(\vec{g}_{k,l,m}) = -m \nabla \Phi(\vec{g}_{k,l,m})$$

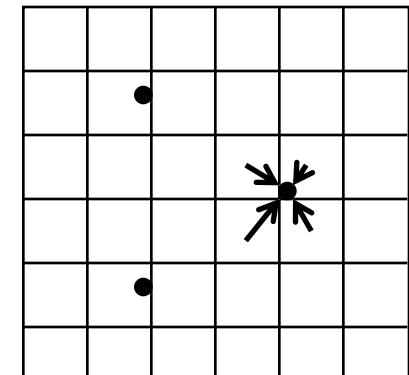


- calculate mass density on grid $\vec{x}_i \rightarrow \rho(\vec{g}_{k,l,m})$
- solve Poisson's equation on grid $\Phi(\vec{g}_{k,l,m})$
- differentiate potential to get forces $\vec{F}(\vec{g}_{k,l,m})$
- interpolate forces back to particles $\vec{F}(\vec{g}_{k,l,m}) \rightarrow \vec{F}(\vec{x}_i)$

Solving for Gravity

- interpolating the forces

$$\vec{F}(\vec{g}_{k,l,m}) \rightarrow \vec{F}(\vec{r}_i)$$

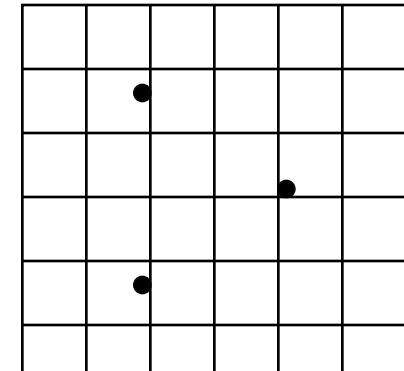


Solving for Gravity

- Particle-Mesh (PM) method

$$\Delta\Phi(\vec{g}_{k,l,m}) = 4\pi G \rho(\vec{g}_{k,l,m})$$

$$\vec{F}(\vec{g}_{k,l,m}) = -m \nabla \Phi(\vec{g}_{k,l,m})$$



1. calculate mass density on grid

$$\vec{x}_i \rightarrow \rho(\vec{g}_{k,l,m})$$

2. solve Poisson's equation on grid

$$\Phi(\vec{g}_{k,l,m})$$

3. differentiate potential to get forces

$$\vec{F}(\vec{g}_{k,l,m})$$

4. interpolate forces back to particles

$$\vec{F}(\vec{g}_{k,l,m}) \rightarrow \vec{F}(\vec{x}_i)$$

anyone fancies to write a PM code as the project?