

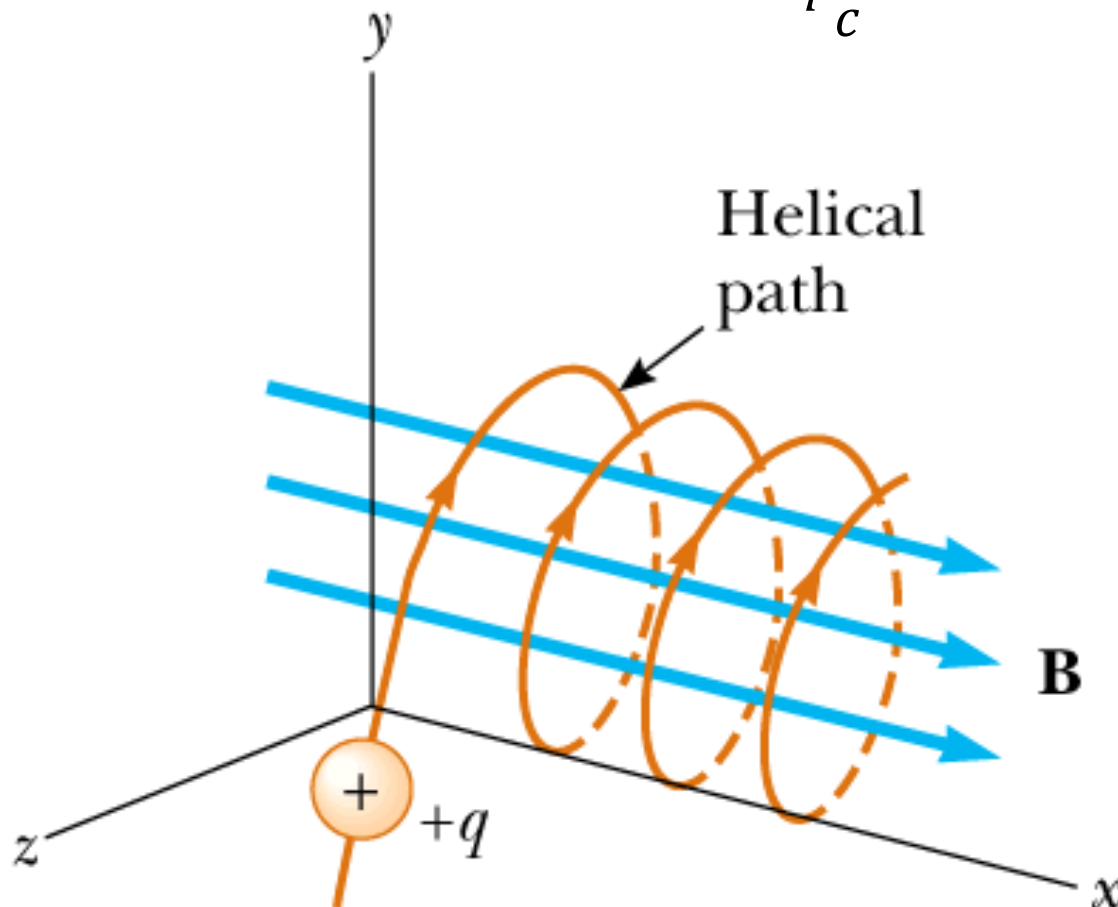
pure non-relativistic treatment!

- Lorentz force
- Maxwell equations
- Poynting vector
- electromagnetic waves
- radiation spectrum
- electromagnetic potentials

- **Lorentz force**
- Maxwell equations
- Poynting vector
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- Lorentz force

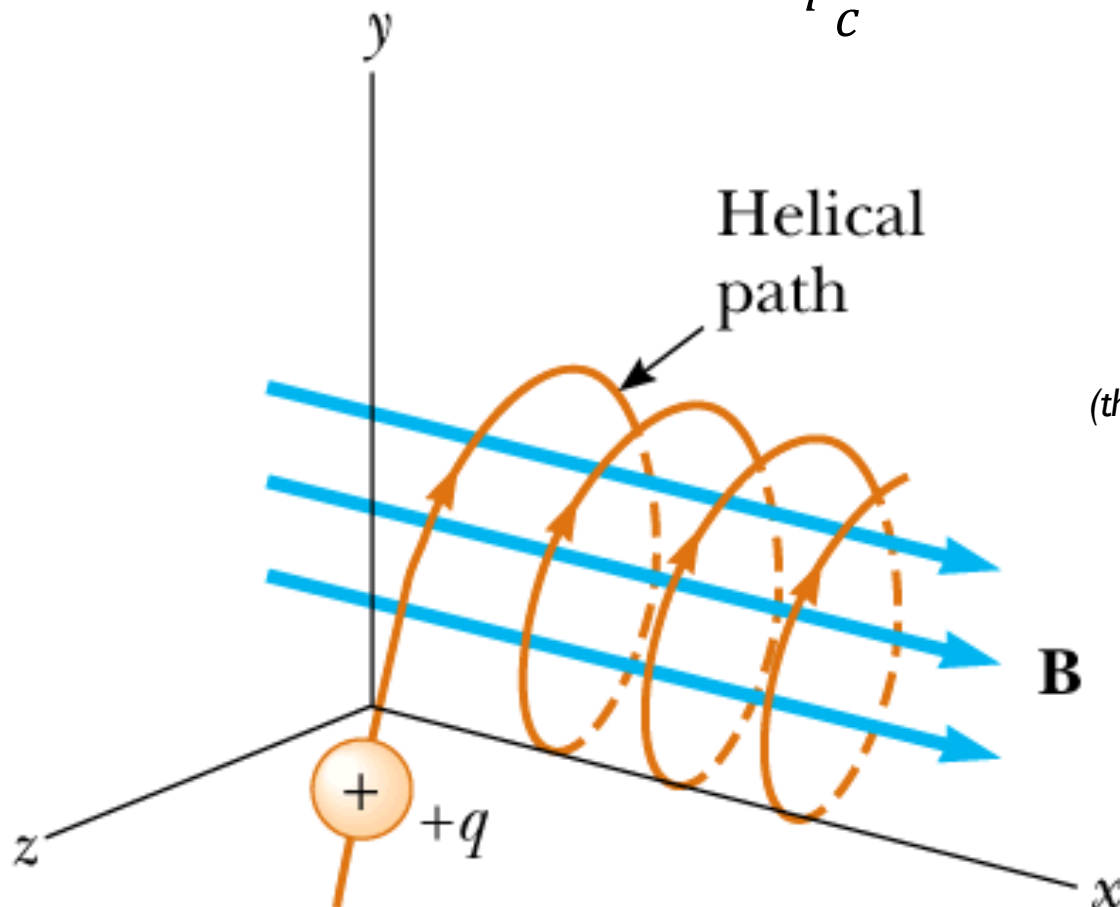
$$\vec{F} = q \frac{1}{c} \vec{v} \times \vec{B}$$



charged particles are forced to spiral around magnetic field lines

- Lorentz force

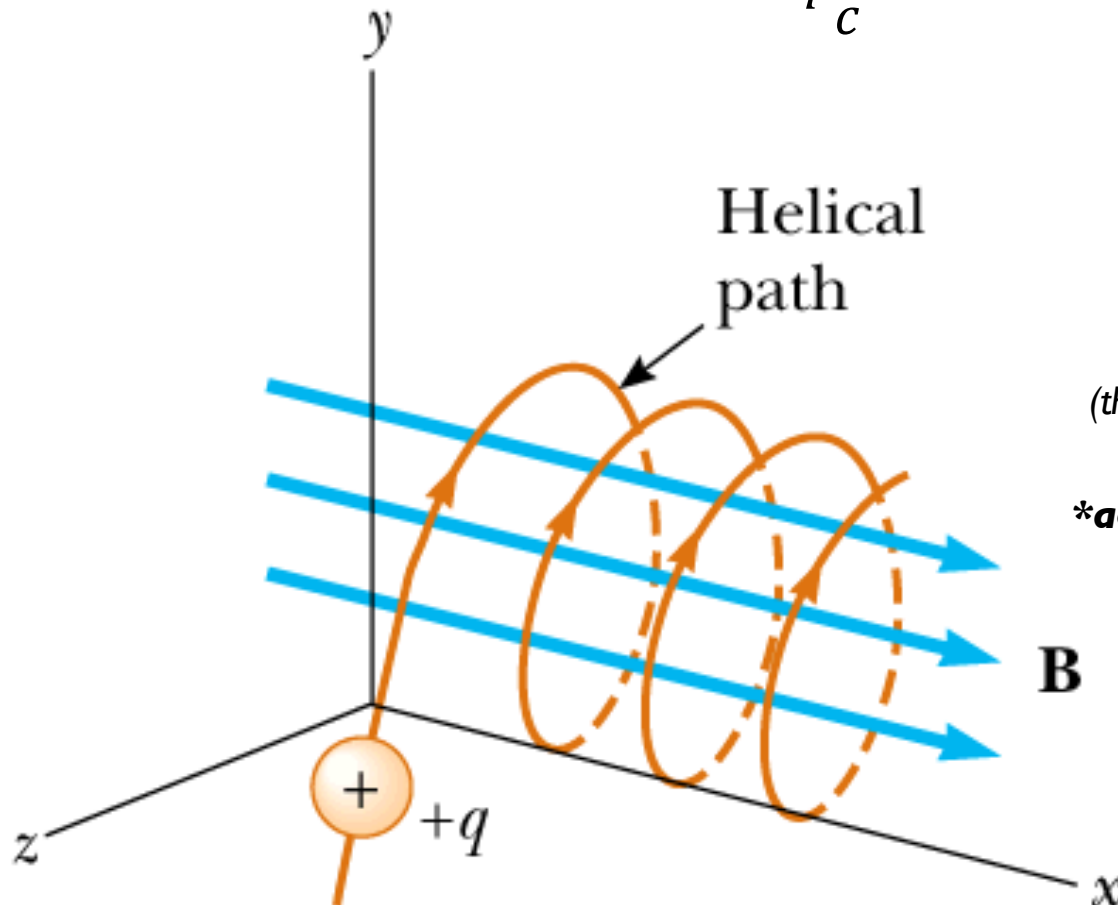
$$\vec{F} = q \frac{1}{c} \vec{v} \times \vec{B}$$



charged particles are forced to spiral around magnetic field lines (they are hence trapped by closed field lines...)

- Lorentz force

$$\vec{F} = q \frac{1}{c} \vec{v} \times \vec{B}$$



charged particles are forced to spiral around magnetic field lines (they are hence trapped by closed field lines...)*

***accelerated charged particles radiate!**

- Lorentz force

$$\vec{F} = q \left(\vec{E} + \frac{1}{c} \vec{v} \times \vec{B} \right)$$

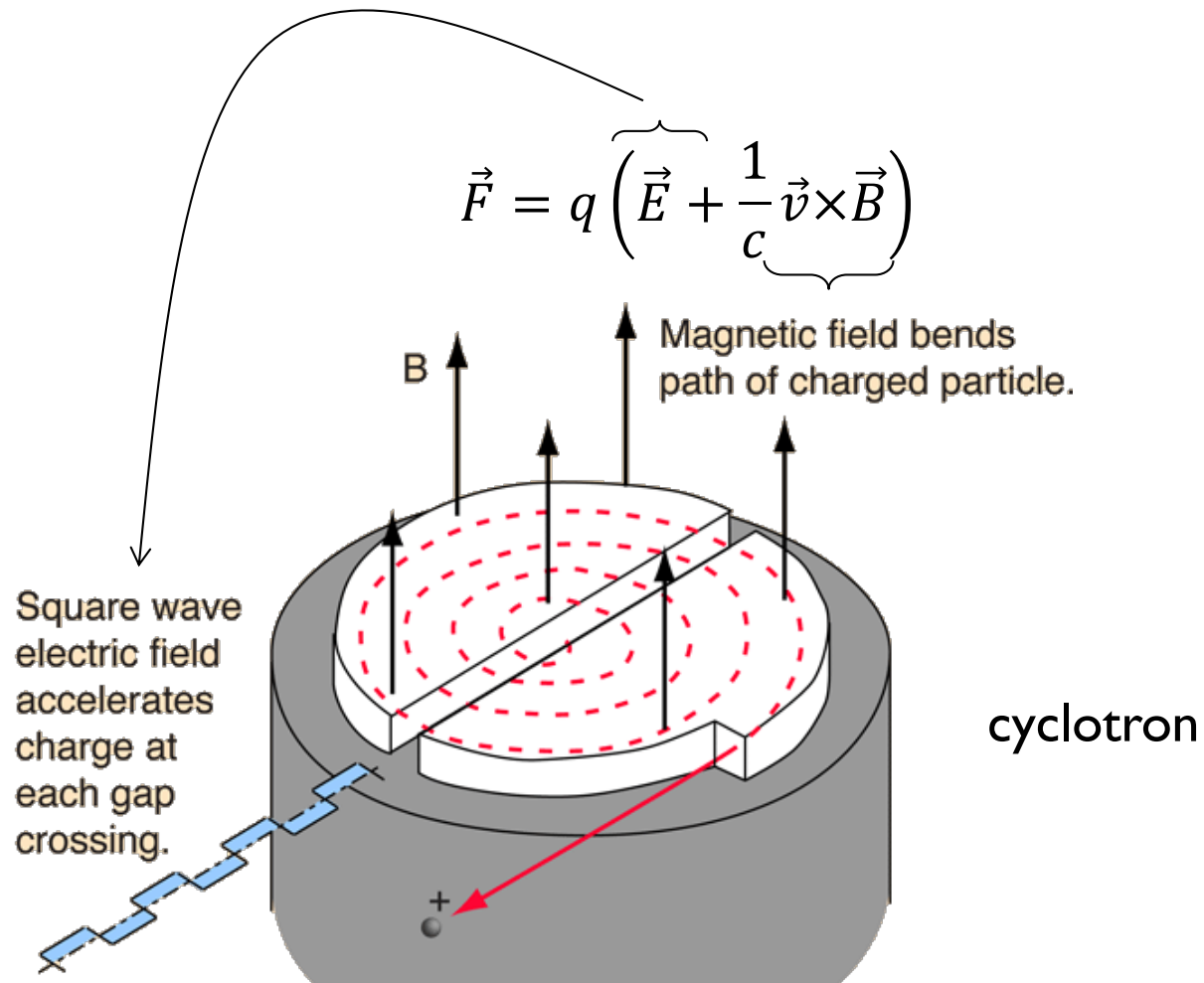
- Lorentz force

$$\vec{F} = q \left(\vec{E} + \frac{1}{c} \vec{v} \times \vec{B} \right)$$

example?

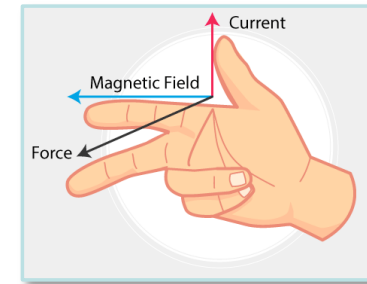
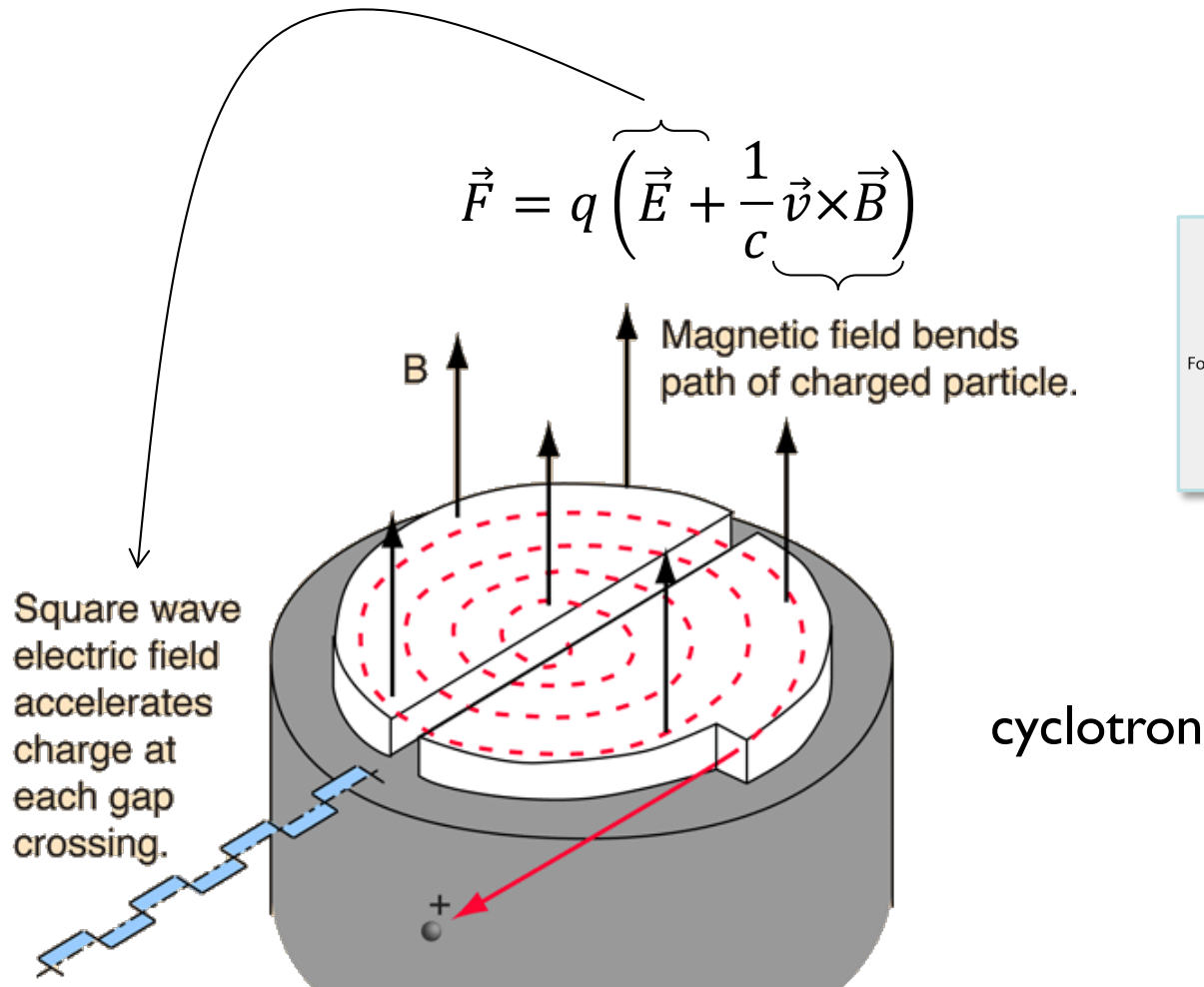
▪ Lorentz force

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- Lorentz force

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- Lorentz force

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- rate of work done by fields?

- Lorentz force

$$\vec{F} = q \left(\vec{E} + \frac{1}{c} \vec{v} \times \vec{B} \right)$$

- rate of work done by fields

$$\frac{dQ}{dt} = \frac{\vec{F} \cdot \overrightarrow{ds}}{dt} = \vec{v} \cdot \vec{F} = \vec{v} \cdot q \left(\vec{E} + \frac{1}{c} \vec{v} \times \vec{B} \right)$$

- Lorentz force

$$\vec{F} = q \left(\vec{E} + \frac{1}{c} \vec{v} \times \vec{B} \right)$$

- rate of work done by fields

$$\vec{v} \cdot \vec{F} = q \vec{v} \cdot \vec{E}$$

- non-relativistic particle

$$\vec{F} = m \frac{d\vec{v}}{dt}$$

- Lorentz force

$$\vec{F} = q \left(\vec{E} + \frac{1}{c} \vec{v} \times \vec{B} \right)$$

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$$\left. \begin{array}{l} \vec{v} \cdot \vec{F} = q \vec{v} \cdot \vec{E} \\ \vec{F} = m \frac{d\vec{v}}{dt} \end{array} \right\} m \vec{v} \cdot \frac{d\vec{v}}{dt} = q \vec{v} \cdot \vec{E}$$

- Lorentz force

$$\vec{F} = q \left(\vec{E} + \frac{1}{c} \vec{v} \times \vec{B} \right)$$

- rate of work done by fields

$$\vec{v} \cdot \vec{F} = q \vec{v} \cdot \vec{E}$$

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$$\left. \begin{array}{l} \vec{v} \cdot \vec{F} = q \vec{v} \cdot \vec{E} \\ \vec{F} = m \frac{d\vec{v}}{dt} \end{array} \right\} \begin{array}{l} m \vec{v} \cdot \frac{d\vec{v}}{dt} = q \vec{v} \cdot \vec{E} \\ \frac{d}{dt} \left(\frac{1}{2} m v^2 \right) \end{array}$$

- Lorentz force

$$\vec{F} = q \left(\vec{E} + \frac{1}{c} \vec{v} \times \vec{B} \right)$$

- rate of work done by fields (on non-relativistic particles)

$$\frac{dE_{kin}}{dt} = q \vec{v} \cdot \vec{E}$$

- Lorentz force

$$\vec{F} = q \left(\vec{E} + \frac{1}{c} \vec{v} \times \vec{B} \right)$$

- rate of work done by fields (on non-relativistic particles)

$$\frac{dE_{kin}}{dt} = \underbrace{q\vec{v}}_{\text{current}} \cdot \vec{E}$$

current $\vec{j} = \rho_q \vec{v}$; ρ_q = charge density

- Lorentz force

$$\vec{F} = q \left(\vec{E} + \frac{1}{c} \vec{v} \times \vec{B} \right)$$

- rate of work done by fields (on non-relativistic particles)

$$\frac{de_{kin}}{dt} = \vec{j} \cdot \vec{E}$$

e_{kin} = energy density

- Lorentz force
- **Maxwell equations**
- Poynting vector
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- Maxwell equations

- Maxwell equations

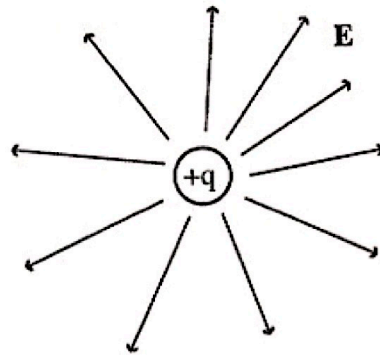
*they describe the behaviour of both electric and magnetic fields
and
their interaction with matter*

- Maxwell equations

$$\vec{D} = \epsilon \vec{E}$$

ϵ : dielectric constant

Gauss' law
 $\nabla \cdot \mathbf{D} = 4\pi\rho$

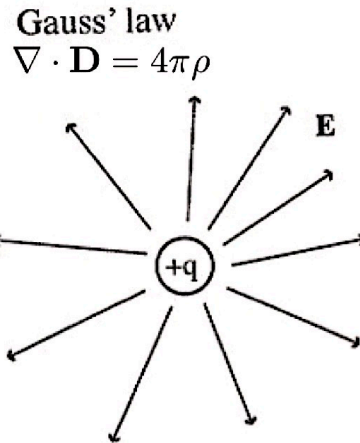


they describe the behaviour of both electric and magnetic fields
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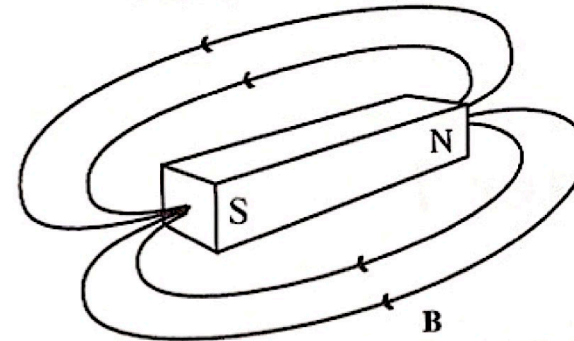
■ Maxwell equations

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No magnetic monopoles
 $\nabla \cdot \mathbf{B} = 0$

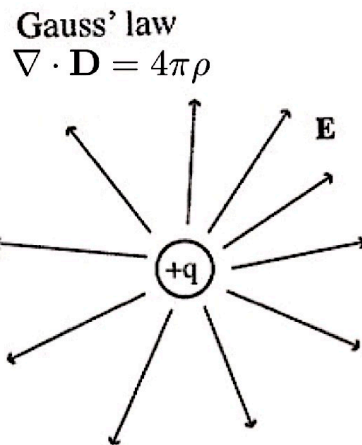


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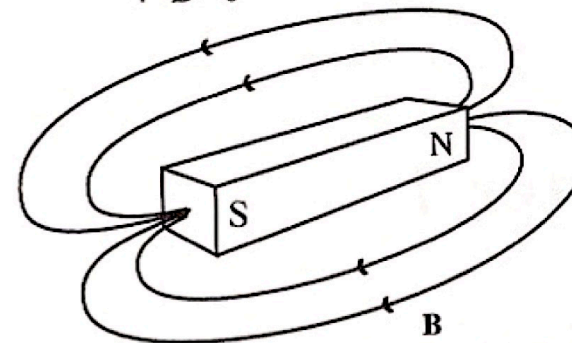
Maxwell equations

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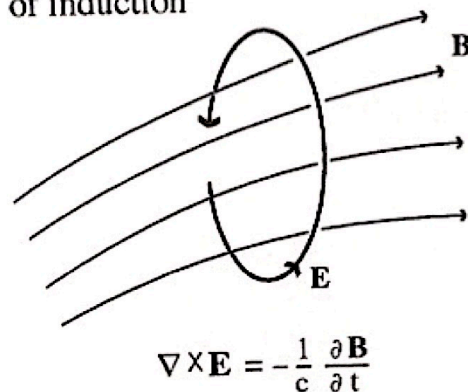
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Faraday's law
of induction

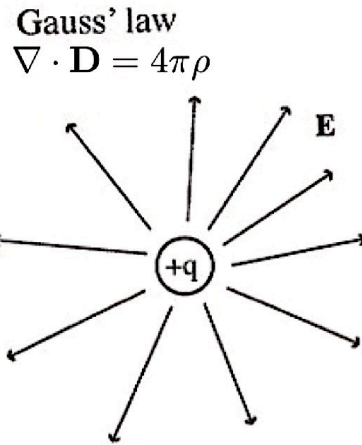


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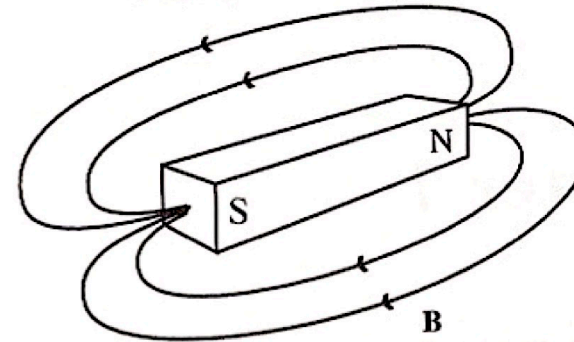
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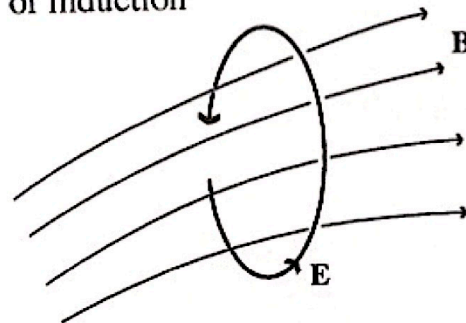
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Faraday's law
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$$\nabla \times \mathbf{E} = -\frac{1}{c} \frac{\partial \mathbf{B}}{\partial t} \longrightarrow$$

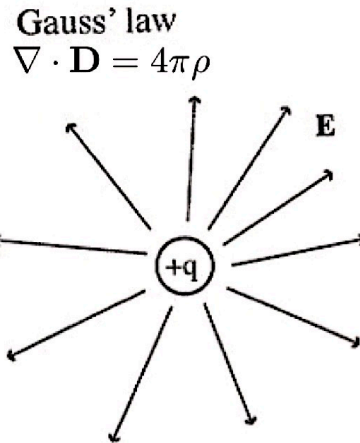
a time-varying B-field induces a (non-conservative) E-field
(which in turn can accelerate charges, leading to a current)

*they describe the behaviour of both electric and magnetic fields
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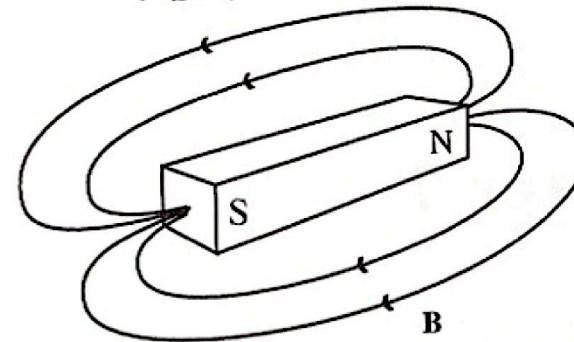
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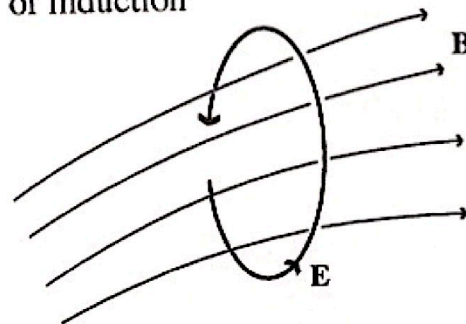
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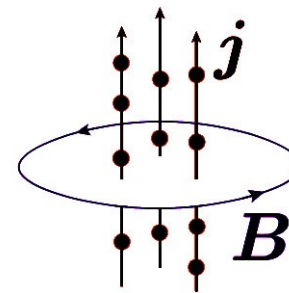
No magnetic monopoles
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Faraday's law
of induction



$$\nabla \times \mathbf{E} = -\frac{1}{c} \frac{\partial \mathbf{B}}{\partial t}$$



$$\nabla \times \mathbf{H} = \frac{4\pi}{c} \mathbf{j} + \frac{1}{c} \frac{\partial \mathbf{D}}{\partial t}$$

$$\vec{H} = \frac{1}{\mu} \vec{B}$$

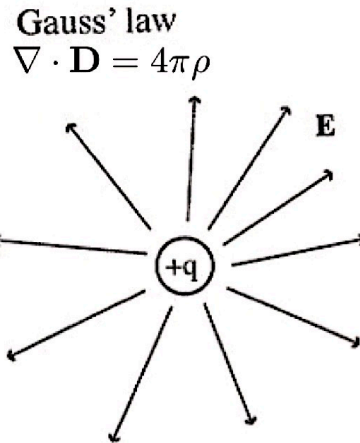
μ : magnetic permeability

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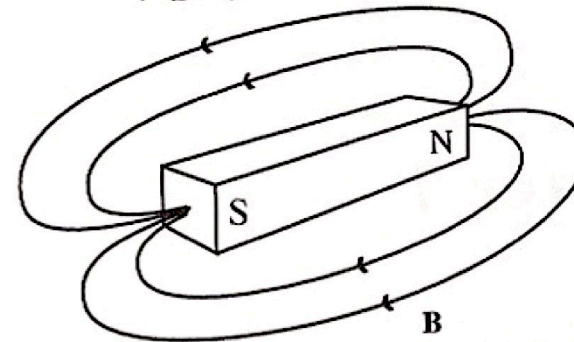
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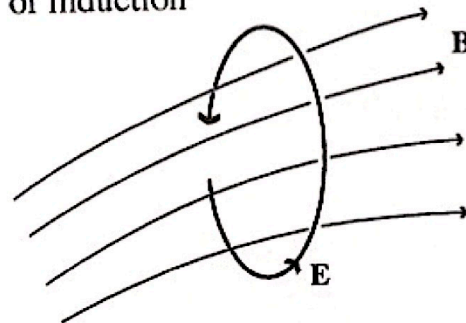
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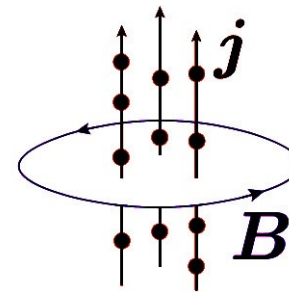


Faraday's law
of induction



$$\nabla \times \mathbf{E} = -\frac{1}{c} \frac{\partial \mathbf{B}}{\partial t}$$

a current and/or time varying E-field induces a B-field



$$\nabla \times \mathbf{H} = \frac{4\pi}{c} \mathbf{j} + \frac{1}{c} \frac{\partial \mathbf{D}}{\partial t}$$

$$\vec{H} = \frac{1}{\mu} \vec{B}$$

μ : magnetic permeability

they describe the behaviour of both electric and magnetic fields
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▪ Maxwell equations

$$\nabla \cdot \vec{D} = 4\pi\rho$$

$$\nabla \cdot \vec{B} = 0$$

$$\nabla \times \vec{E} = -\frac{1}{c} \frac{\partial \vec{B}}{\partial t}$$

$$\nabla \times \vec{H} = \frac{4\pi}{c} \vec{j} + \frac{1}{c} \frac{\partial \vec{D}}{\partial t}$$

$$\vec{j} = \rho \vec{v}$$

$$\vec{D} = \epsilon \vec{E}$$

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charge conservation:
$$\nabla \times \vec{H} = \frac{4\pi}{c} \vec{j} + \frac{1}{c} \frac{\partial \vec{D}}{\partial t}$$

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$$\vec{j} = \rho \vec{v}$$

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- rate of work done per unit volume:

$$\frac{de_{kin}}{dt} = \vec{j} \cdot \vec{E}$$

$$\vec{j} = \rho \vec{v}$$

$$\vec{D} = \epsilon \vec{E}$$

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- rate of work done per unit volume:

$$\frac{de_{kin}}{dt} = \vec{J} \cdot \vec{E} \quad \rightarrow \text{eventually leading to the Poynting vector...}$$

$$\vec{J} = \rho \vec{v}$$

$$\vec{D} = \epsilon \vec{E}$$

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- rate of work done per unit volume:

$$\frac{de_{kin}}{dt} = \vec{J} \cdot \vec{E} \longrightarrow \vec{E} \cdot (\nabla \times \vec{H}) = \frac{4\pi}{c} \vec{E} \cdot \vec{J} + \frac{1}{c} \vec{E} \cdot \frac{\partial \vec{D}}{\partial t}$$

$$\vec{J} = \rho \vec{v}$$

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$$\nabla \cdot \vec{D} = 4\pi\rho$$

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$$\frac{de_{kin}}{dt} = \vec{j} \cdot \vec{E} \longrightarrow \vec{E} \cdot (\nabla \times \vec{H}) = \frac{4\pi}{c} \vec{E} \cdot \vec{j} + \frac{1}{c} \vec{E} \cdot \frac{\partial \vec{D}}{\partial t}$$

$$\begin{aligned} \vec{j} &= \rho \vec{v} \\ \vec{D} &= \epsilon \vec{E} \\ \vec{H} &= \frac{1}{\mu} \vec{B} \end{aligned}$$

$$\vec{j} \cdot \vec{E} = \frac{1}{4\pi} \left(c \vec{E} \cdot (\nabla \times \vec{H}) - \frac{1}{c} \vec{E} \cdot \frac{\partial \vec{D}}{\partial t} \right)$$

- Maxwell equations

$$\nabla \cdot \vec{D} = 4\pi\rho$$

$$\nabla \cdot \vec{B} = 0$$

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- rate of work done per unit volume:

$$\frac{de_{kin}}{dt} = \vec{J} \cdot \vec{E}$$

$$\vec{E} \cdot (\nabla \times \vec{H}) = \vec{H} \cdot (\nabla \times \vec{E}) - \nabla \cdot (\vec{E} \times \vec{H})$$

$$\vec{J} \cdot \vec{E} = \frac{1}{4\pi} \left(\overbrace{c \vec{E} \cdot (\nabla \times \vec{H})} - \frac{1}{c} \vec{E} \cdot \frac{\partial \vec{D}}{\partial t} \right)$$

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$$\frac{de_{kin}}{dt} = \vec{j} \cdot \vec{E} \longrightarrow \vec{j} \cdot \vec{E} = \frac{1}{4\pi} \left(-\vec{H} \cdot \frac{\partial \vec{B}}{\partial t} - c \nabla \cdot (\vec{E} \times \vec{H}) - \vec{E} \cdot \frac{\partial \vec{D}}{\partial t} \right)$$

$$\vec{j} = \rho \vec{v}$$

$$\vec{D} = \epsilon \vec{E}$$

$$\vec{H} = \frac{1}{\mu} \vec{B}$$

- Maxwell equations

$$\nabla \cdot \vec{D} = 4\pi\rho$$

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Poynting vector...

- Lorentz force
- Maxwell equations
- **Poynting vector**
- electromagnetic waves
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- Maxwell equations

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$$\nabla \cdot \vec{B} = 0$$

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$$\nabla \times \vec{H} = \frac{4\pi}{c} \vec{j} + \frac{1}{c} \frac{\partial \vec{D}}{\partial t} \xrightarrow{\nabla \cdot} 0 = \vec{j} + \frac{\partial \vec{\rho}}{\partial t} \xrightarrow{\vec{j} = \rho \vec{v}} \frac{\partial \vec{\rho}}{\partial t} + \rho \vec{v} = 0$$

charge conservation:

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$$\nabla \cdot \vec{D} = 4\pi\rho$$

$$\nabla \cdot \vec{B} = 0$$

$$\nabla \times \vec{E} = -\frac{1}{c} \frac{\partial \vec{B}}{\partial t}$$

$$\nabla \times \vec{H} = \frac{4\pi}{c} \vec{j} + \frac{1}{c} \frac{\partial \vec{D}}{\partial t} \xrightarrow{\nabla \cdot} 0 = \vec{j} + \frac{\partial \vec{\rho}}{\partial t} \xrightarrow{\vec{j} = \rho \vec{v}} \frac{\partial \vec{\rho}}{\partial t} + \rho \vec{v} = 0$$

charge conservation:

- Poynting theorem:

$$\frac{de_{kin}}{dt} = \frac{1}{4\pi} \left(-\vec{H} \cdot \frac{\partial \vec{B}}{\partial t} - c \nabla \cdot (\vec{E} \times \vec{H}) - \vec{E} \cdot \frac{\partial \vec{D}}{\partial t} \right)$$

$$\vec{j} = \rho \vec{v}$$

$$\vec{D} = \epsilon \vec{E}$$

$$\vec{H} = \frac{1}{\mu} \vec{B}$$

- Maxwell equations

$$\nabla \cdot \vec{D} = 4\pi\rho$$

$$\nabla \cdot \vec{B} = 0$$

$$\nabla \times \vec{E} = -\frac{1}{c} \frac{\partial \vec{B}}{\partial t}$$

$$\nabla \times \vec{H} = \frac{4\pi}{c} \vec{j} + \frac{1}{c} \frac{\partial \vec{D}}{\partial t} \xrightarrow{\nabla \cdot} 0 = \vec{j} + \frac{\partial \rho}{\partial t} \xrightarrow{\vec{j} = \rho \vec{v}} \frac{\partial \rho}{\partial t} + \rho \vec{v} = 0$$

charge conservation:

- Poynting theorem:

$$\frac{de_{kin}}{dt} = \frac{1}{4\pi} \left(-\vec{H} \cdot \frac{\partial \vec{B}}{\partial t} - c \nabla \cdot (\vec{E} \times \vec{H}) - \vec{E} \cdot \frac{\partial \vec{D}}{\partial t} \right)$$

$$\begin{aligned} \vec{j} &= \rho \vec{v} \\ \vec{D} &= \epsilon \vec{E} \\ \vec{H} &= \frac{1}{\mu} \vec{B} \end{aligned}$$

the energy density changes at a rate given by the work done on the charges

- Maxwell equations

$$\nabla \cdot \vec{D} = 4\pi\rho$$

$$\nabla \cdot \vec{B} = 0$$

$$\nabla \times \vec{E} = -\frac{1}{c} \frac{\partial \vec{B}}{\partial t}$$

$$\nabla \times \vec{H} = \frac{4\pi}{c} \vec{j} + \frac{1}{c} \frac{\partial \vec{D}}{\partial t} \xrightarrow{\nabla \cdot} 0 = \vec{j} + \frac{\partial \vec{\rho}}{\partial t} \xrightarrow{\vec{j} = \rho \vec{v}} \frac{\partial \vec{\rho}}{\partial t} + \rho \vec{v} = 0$$

charge conservation:

- Poynting theorem:

$$\frac{de_{kin}}{dt} = \frac{1}{4\pi} \left(-\vec{H} \cdot \frac{\partial \vec{B}}{\partial t} - c \nabla \cdot (\vec{E} \times \vec{H}) - \vec{E} \cdot \frac{\partial \vec{D}}{\partial t} \right)$$

$$\begin{aligned} \vec{j} &= \rho \vec{v} \\ \vec{D} &= \epsilon \vec{E} \\ \vec{H} &= \frac{1}{\mu} \vec{B} \end{aligned}$$

*the energy density changes at a rate given by the work done on the charges,
minus the rate at which energy is radiated away*

- Maxwell equations

$$\nabla \cdot \vec{D} = 4\pi\rho$$

$$\nabla \cdot \vec{B} = 0$$

$$\nabla \times \vec{E} = -\frac{1}{c} \frac{\partial \vec{B}}{\partial t}$$

$$\nabla \times \vec{H} = \frac{4\pi}{c} \vec{j} + \frac{1}{c} \frac{\partial \vec{D}}{\partial t} \xrightarrow{\nabla \cdot} 0 = \vec{j} + \frac{\partial \rho}{\partial t} \xrightarrow{\vec{j} = \rho \vec{v}} \frac{\partial \rho}{\partial t} + \rho \vec{v} = 0$$

charge conservation:

- Poynting theorem:

$$\begin{aligned} \frac{de_{kin}}{dt} &= \frac{1}{4\pi} \left(-\vec{H} \cdot \frac{\partial \vec{B}}{\partial t} - c \nabla \cdot (\vec{E} \times \vec{H}) - \vec{E} \cdot \frac{\partial \vec{D}}{\partial t} \right) \\ &= -\frac{1}{8\pi} \frac{\partial}{\partial t} \left(\epsilon E^2 + \frac{1}{\mu} B^2 \right) - \frac{c}{4\pi} \nabla \cdot (\vec{E} \times \vec{H}) \end{aligned}$$

$\epsilon = const., \mu = const.$

$$\begin{aligned} \vec{j} &= \rho \vec{v} \\ \vec{D} &= \epsilon \vec{E} \\ \vec{H} &= \frac{1}{\mu} \vec{B} \end{aligned}$$

- Maxwell equations

$$\nabla \cdot \vec{D} = 4\pi\rho$$

$$\nabla \cdot \vec{B} = 0$$

$$\nabla \times \vec{E} = -\frac{1}{c} \frac{\partial \vec{B}}{\partial t}$$

$$\nabla \times \vec{H} = \frac{4\pi}{c} \vec{j} + \frac{1}{c} \frac{\partial \vec{D}}{\partial t} \xrightarrow{\nabla \cdot} 0 = \vec{j} + \frac{\partial \vec{\rho}}{\partial t} \xrightarrow{\vec{j} = \rho \vec{v}} \frac{\partial \vec{\rho}}{\partial t} + \rho \vec{v} = 0$$

charge conservation:

- Poynting theorem:

$$\frac{de_{kin}}{dt} = -\frac{1}{8\pi} \frac{\partial}{\partial t} \left(\epsilon E^2 + \frac{1}{\mu} B^2 \right) - \nabla \cdot \frac{c}{4\pi} (\vec{E} \times \vec{H})$$

$$\vec{j} = \rho \vec{v}$$

$$\vec{D} = \epsilon \vec{E}$$

$$\vec{H} = \frac{1}{\mu} \vec{B}$$

- Maxwell equations

$$\nabla \cdot \vec{D} = 4\pi\rho$$

$$\nabla \cdot \vec{B} = 0$$

$$\nabla \times \vec{E} = -\frac{1}{c} \frac{\partial \vec{B}}{\partial t}$$

$$\nabla \times \vec{H} = \frac{4\pi}{c} \vec{j} + \frac{1}{c} \frac{\partial \vec{D}}{\partial t} \xrightarrow{\nabla \cdot} 0 = \vec{j} + \frac{\partial \vec{\rho}}{\partial t} \xrightarrow{\vec{j} = \rho \vec{v}} \frac{\partial \vec{\rho}}{\partial t} + \rho \vec{v} = 0$$

charge conservation:

- Poynting theorem:

$$\frac{de_{kin}}{dt} = -\frac{1}{8\pi} \frac{\partial}{\partial t} \left(\epsilon E^2 + \frac{1}{\mu} B^2 \right) - \nabla \cdot \underbrace{\frac{c}{4\pi} (\vec{E} \times \vec{H})}_{\vec{S}}$$

$$\vec{j} = \rho \vec{v}$$

$$\vec{D} = \epsilon \vec{E}$$

$$\vec{H} = \frac{1}{\mu} \vec{B}$$

Poynting vector: $\vec{S} = \frac{c}{4\pi} (\vec{E} \times \vec{H})$ directional energy flux

- Maxwell equations

$$\nabla \cdot \vec{D} = 4\pi\rho$$

$$\nabla \cdot \vec{B} = 0$$

$$\nabla \times \vec{E} = -\frac{1}{c} \frac{\partial \vec{B}}{\partial t}$$

$$\nabla \times \vec{H} = \frac{4\pi}{c} \vec{j} + \frac{1}{c} \frac{\partial \vec{D}}{\partial t} \xrightarrow{\nabla \cdot} 0 = \vec{j} + \frac{\partial \vec{\rho}}{\partial t} \xrightarrow{\vec{j} = \rho \vec{v}} \frac{\partial \vec{\rho}}{\partial t} + \rho \vec{v} = 0$$

charge conservation:

- Poynting theorem:

$$\frac{de_{kin}}{dt} = -\frac{1}{8\pi} \frac{\partial}{\partial t} \left(\epsilon E^2 + \frac{1}{\mu} B^2 \right) - \nabla \cdot \vec{S} \qquad \vec{S} = \frac{c}{4\pi} (\vec{E} \times \vec{H})$$

$$\vec{j} = \rho \vec{v}$$

$$\vec{D} = \epsilon \vec{E}$$

$$\vec{H} = \frac{1}{\mu} \vec{B}$$

- Maxwell equations

$$\nabla \cdot \vec{D} = 4\pi\rho$$

$$\nabla \cdot \vec{B} = 0$$

$$\nabla \times \vec{E} = -\frac{1}{c} \frac{\partial \vec{B}}{\partial t}$$

$$\nabla \times \vec{H} = \frac{4\pi}{c} \vec{j} + \frac{1}{c} \frac{\partial \vec{D}}{\partial t} \xrightarrow{\nabla \cdot} 0 = \vec{j} + \frac{\partial \vec{\rho}}{\partial t} \xrightarrow{\vec{j} = \rho \vec{v}} \frac{\partial \vec{\rho}}{\partial t} + \rho \vec{v} = 0$$

charge conservation:

- Poynting theorem:

$$\frac{de_{kin}}{dt} = -\frac{1}{8\pi} \frac{\partial}{\partial t} \left(\epsilon E^2 + \frac{1}{\mu} B^2 \right) - \nabla \cdot \vec{S} \quad \vec{S} = \frac{c}{4\pi} (\vec{E} \times \vec{H})$$

$$\begin{aligned} \vec{j} &= \rho \vec{v} \\ \vec{D} &= \epsilon \vec{E} \\ \vec{H} &= \frac{1}{\mu} \vec{B} \end{aligned}$$

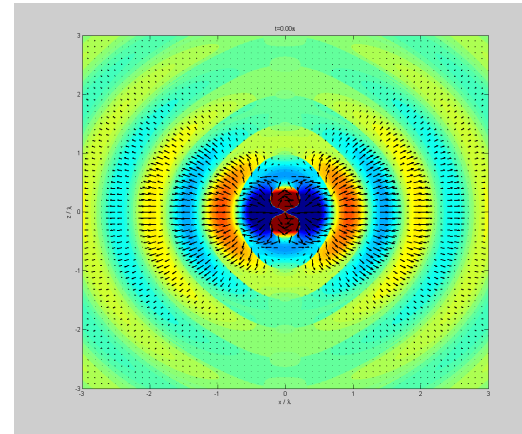
flux is orthogonal to both \vec{E} and \vec{B} !

- Maxwell equations

$$\nabla \cdot \vec{D} = 4\pi\rho$$

$$\nabla \cdot \vec{B} = 0$$

$$\nabla \times \vec{E} = -\frac{1}{c} \frac{\partial \vec{B}}{\partial t}$$



Dipole radiation of a dipole vertically in the page, showing electric field strength (color) and Poynting vector (arrows) in the plane of the page.

charge conservation:

$$\nabla \times \vec{H} = \frac{4\pi}{c} \vec{j} + \frac{1}{c} \frac{\partial \vec{D}}{\partial t} \quad \xrightarrow{\nabla \cdot} \quad 0 = \vec{j} + \frac{\partial \vec{\rho}}{\partial t} \quad \xrightarrow{\vec{j} = \rho \vec{v}} \quad \frac{\partial \vec{\rho}}{\partial t} + \rho \vec{v} = 0$$

- Poynting theorem:

$$\frac{de_{kin}}{dt} = -\frac{1}{8\pi} \frac{\partial}{\partial t} \left(\epsilon E^2 + \frac{1}{\mu} B^2 \right) - \nabla \cdot \vec{S} \quad \vec{S} = \frac{c}{4\pi} (\vec{E} \times \vec{H})$$

flux is orthogonal to both \vec{E} and \vec{B} !

$$\begin{aligned} \vec{j} &= \rho \vec{v} \\ \vec{D} &= \epsilon \vec{E} \\ \vec{H} &= \frac{1}{\mu} \vec{B} \end{aligned}$$

- Maxwell equations

$$\nabla \cdot \vec{D} = 4\pi\rho$$

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charge conservation:

- Poynting theorem: *integral form!?*

$$\frac{de_{kin}}{dt} = -\frac{1}{8\pi} \frac{\partial}{\partial t} \left(\epsilon E^2 + \frac{1}{\mu} B^2 \right) - \nabla \cdot \vec{S} \qquad \vec{S} = \frac{c}{4\pi} (\vec{E} \times \vec{H})$$

$$\begin{aligned} \vec{j} &= \rho \vec{v} \\ \vec{D} &= \epsilon \vec{E} \\ \vec{H} &= \frac{1}{\mu} \vec{B} \end{aligned}$$

- Maxwell equations

$$\nabla \cdot \vec{D} = 4\pi\rho$$

$$\nabla \cdot \vec{B} = 0$$

$$\nabla \times \vec{E} = -\frac{1}{c} \frac{\partial \vec{B}}{\partial t}$$

$$\nabla \times \vec{H} = \frac{4\pi}{c} \vec{j} + \frac{1}{c} \frac{\partial \vec{D}}{\partial t} \xrightarrow{\nabla \cdot} 0 = \vec{j} + \frac{\partial \vec{\rho}}{\partial t} \xrightarrow{\vec{j} = \rho \vec{v}} \frac{\partial \vec{\rho}}{\partial t} + \rho \vec{v} = 0$$

charge conservation:

- Poynting theorem: *integral form*

$$\int_V \frac{de_{kin}}{dt} dV = -\frac{d}{dt} \int_V \frac{1}{8\pi} \left(\epsilon E^2 + \frac{1}{\mu} B^2 \right) dV - \int_{\Sigma} \vec{S} \cdot d\vec{A} \quad \vec{S} = \frac{c}{4\pi} (\vec{E} \times \vec{H})$$

$$\vec{j} = \rho \vec{v}$$

$$\vec{D} = \epsilon \vec{E}$$

$$\vec{H} = \frac{1}{\mu} \vec{B}$$

- Maxwell equations

$$\nabla \cdot \vec{D} = 4\pi\rho$$

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$$\nabla \times \vec{E} = -\frac{1}{c} \frac{\partial \vec{B}}{\partial t}$$

$$\nabla \times \vec{H} = \frac{4\pi}{c} \vec{j} + \frac{1}{c} \frac{\partial \vec{D}}{\partial t} \xrightarrow{\nabla \cdot} 0 = \vec{j} + \frac{\partial \vec{\rho}}{\partial t} \xrightarrow{\vec{j} = \rho \vec{v}} \frac{\partial \vec{\rho}}{\partial t} + \rho \vec{v} = 0$$

charge conservation:

- Poynting theorem: *integral form*

$$\underbrace{\int_V \frac{de_{kin}}{dt} dV}_{\text{rate of change of mechanical energy}} = - \underbrace{\frac{d}{dt} \int_V \frac{1}{8\pi} \left(\epsilon E^2 + \frac{1}{\mu} B^2 \right) dV}_{\text{rate of change of E- and B-field energy}} - \int_{\Sigma} \vec{S} \cdot d\vec{A} \quad \vec{S} = \frac{c}{4\pi} (\vec{E} \times \vec{H})$$

$$\vec{j} = \rho \vec{v}$$

$$\vec{D} = \epsilon \vec{E} \quad \text{rate of change of mechanical energy}$$

$$\vec{H} = \frac{1}{\mu} \vec{B}$$

- Maxwell equations

$$\nabla \cdot \vec{D} = 4\pi\rho$$

$$\nabla \cdot \vec{B} = 0$$

$$\nabla \times \vec{E} = -\frac{1}{c} \frac{\partial \vec{B}}{\partial t}$$

$$\nabla \times \vec{H} = \frac{4\pi}{c} \vec{j} + \frac{1}{c} \frac{\partial \vec{D}}{\partial t} \xrightarrow{\nabla \cdot} 0 = \vec{j} + \frac{\partial \vec{\rho}}{\partial t} \xrightarrow{\vec{j} = \rho \vec{v}} \frac{\partial \vec{\rho}}{\partial t} + \rho \vec{v} = 0$$

charge conservation:

- Poynting theorem: *integral form*

$$\frac{d}{dt} (U_{mech} + U_{field}) = - \int_{\Sigma} \vec{S} \cdot d\vec{A} \qquad \vec{S} = \frac{c}{4\pi} (\vec{E} \times \vec{H})$$

$$\vec{j} = \rho \vec{v}$$

$$\vec{D} = \epsilon \vec{E}$$

$$\vec{H} = \frac{1}{\mu} \vec{B}$$

- Maxwell equations

$$\nabla \cdot \vec{D} = 4\pi\rho$$

$$\nabla \cdot \vec{B} = 0$$

$$\nabla \times \vec{E} = -\frac{1}{c} \frac{\partial \vec{B}}{\partial t}$$

$$\nabla \times \vec{H} = \frac{4\pi}{c} \vec{J} + \frac{1}{c} \frac{\partial \vec{D}}{\partial t} \xrightarrow{\nabla \cdot} 0 = \vec{J} + \frac{\partial \vec{\rho}}{\partial t} \xrightarrow{\vec{J} = \rho \vec{v}} \frac{\partial \vec{\rho}}{\partial t} + \rho \vec{v} = 0$$

charge conservation:

- Poynting theorem: *integral form*

$$\frac{d}{dt} (U_{mech} + U_{field}) = - \int_{\Sigma} \vec{S} \cdot d\vec{A} \qquad \vec{S} = \frac{c}{4\pi} (\vec{E} \times \vec{H})$$

$$\begin{aligned} \vec{J} &= \rho \vec{v} \\ \vec{D} &= \epsilon \vec{E} \\ \vec{H} &= \frac{1}{\mu} \vec{B} \end{aligned}$$

*the rate of change of total energy within a volume
is equal to the net inward flow of energy through the bounding surface*

- Radiation?

- Poynting theorem:

$$\frac{d}{dt} (U_{mech} + U_{field}) = - \int_{\Sigma} \vec{S} \cdot d\vec{A} \qquad \vec{S} = \frac{c}{4\pi} (\vec{E} \times \vec{H})$$

*the rate of change of total energy within a volume
is equal to the net inward flow of energy through the bounding surface*

- Radiation

In **electrostatics** both \vec{E} and \vec{B} decrease like r^{-2}

→ \vec{S} decreases like r^{-4} and thus the integral goes to zero since the surface area increases only as r^{-2} .

- Poynting theorem:

$$\frac{d}{dt} (U_{mech} + U_{field}) = - \int_{\Sigma} \vec{S} \cdot d\vec{A} \qquad \vec{S} = \frac{c}{4\pi} (\vec{E} \times \vec{H})$$

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■ Radiation

In **electrostatics** both \vec{E} and \vec{B} decrease like r^{-2}

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For **time varying** fields \vec{E} and \vec{B} decrease like r^{-1}

→ the integral can contribute a finite amount to the rate of change of energy of the system.

■ Poynting theorem:

$$\frac{d}{dt} (U_{mech} + U_{field}) = - \int_{\Sigma} \vec{S} \cdot d\vec{A} \qquad \vec{S} = \frac{c}{4\pi} (\vec{E} \times \vec{H})$$

*the rate of change of total energy within a volume
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■ Radiation

In electrostatics both \vec{E} and \vec{B} decrease like r^{-2}

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For **time varying** fields \vec{E} and \vec{B} decrease like r^{-1}

→ the integral can contribute a finite amount to the rate of change of energy of the system.

This energy flowing in (or out) at large distances is called **radiation**.

■ Poynting theorem:

$$\frac{d}{dt} (U_{mech} + U_{field}) = - \int_{\Sigma} \vec{S} \cdot d\vec{A} \qquad \vec{S} = \frac{c}{4\pi} (\vec{E} \times \vec{H})$$

*the rate of change of total energy within a volume
is equal to the net inward flow of energy through the bounding surface*

- Lorentz force
- Maxwell equations
- Poynting vector
- **electromagnetic waves**
- radiation spectrum
- electromagnetic potentials

- Electromagnetic Waves

- Electromagnetic Waves – Maxwell equations

$$\nabla \cdot \vec{D} = 4\pi\rho$$

$$\nabla \cdot \vec{B} = 0$$

$$\nabla \times \vec{E} = -\frac{1}{c} \frac{\partial \vec{B}}{\partial t}$$

$$\nabla \times \vec{H} = \frac{4\pi}{c} \vec{J} + \frac{1}{c} \frac{\partial \vec{D}}{\partial t}$$

- Electromagnetic Waves – Maxwell equations in vacuum

$$\nabla \cdot \vec{E} = 0$$

$$\nabla \cdot \vec{B} = 0$$

$$\nabla \times \vec{E} = -\frac{1}{c} \frac{\partial \vec{B}}{\partial t}$$

$$\nabla \times \vec{B} = \frac{1}{c} \frac{\partial \vec{E}}{\partial t}$$

- Electromagnetic Waves – Maxwell equations in vacuum

$$\nabla \cdot \vec{E} = 0$$

$$\nabla \cdot \vec{B} = 0$$

$$\nabla \times \vec{E} = -\frac{1}{c} \frac{\partial \vec{B}}{\partial t} \quad \xrightarrow{\nabla \times} \quad \nabla \times (\nabla \times \vec{E}) = -\frac{1}{c} \frac{\partial \nabla \times \vec{B}}{\partial t}$$

$$\nabla \times \vec{B} = \frac{1}{c} \frac{\partial \vec{E}}{\partial t}$$

- Electromagnetic Waves – Maxwell equations in vacuum

$$\nabla \cdot \vec{E} = 0$$

$$\nabla \cdot \vec{B} = 0$$

$$\nabla \times \vec{E} = -\frac{1}{c} \frac{\partial \vec{B}}{\partial t} \quad \xrightarrow{\nabla \times} \quad \nabla \times (\nabla \times \vec{E}) = -\frac{1}{c} \frac{\partial \nabla \times \vec{B}}{\partial t}$$

$$\nabla \times \vec{B} = \frac{1}{c} \frac{\partial \vec{E}}{\partial t}$$

- Electromagnetic Waves – Maxwell equations in vacuum

$$\nabla \cdot \vec{E} = 0$$

$$\nabla \cdot \vec{B} = 0$$

$$\nabla \times \vec{E} = -\frac{1}{c} \frac{\partial \vec{B}}{\partial t} \xrightarrow{\nabla \times} \nabla \times (\nabla \times \vec{E}) = -\frac{1}{c} \frac{\partial \nabla \times \vec{B}}{\partial t}$$

$$\nabla \times \vec{B} = \frac{1}{c} \frac{\partial \vec{E}}{\partial t} \longrightarrow = -\frac{1}{c} \frac{\partial \nabla \times \vec{B}}{\partial t} = -\frac{1}{c^2} \frac{\partial^2 \vec{E}}{\partial t^2}$$

- Electromagnetic Waves – Maxwell equations in vacuum

$$\nabla \cdot \vec{E} = 0$$

$$\nabla \cdot \vec{B} = 0$$

$$\begin{array}{l} \nabla \times \vec{E} = -\frac{1}{c} \frac{\partial \vec{B}}{\partial t} \xrightarrow{\nabla \times} \nabla \times (\nabla \times \vec{E}) = -\frac{1}{c} \frac{\partial \nabla \times \vec{B}}{\partial t} \\ \nabla \times \vec{B} = \frac{1}{c} \frac{\partial \vec{E}}{\partial t} \xrightarrow{\quad} \qquad \qquad \qquad = -\frac{1}{c} \frac{\partial \nabla \times \vec{B}}{\partial t} = -\frac{1}{c^2} \frac{\partial^2 \vec{E}}{\partial t^2} \\ \qquad \qquad \qquad \qquad \qquad \qquad \qquad \qquad \qquad \qquad \qquad \qquad \qquad \qquad \qquad \qquad \qquad \qquad \downarrow \\ \qquad \qquad \qquad \qquad \qquad \qquad \qquad \qquad \qquad \qquad \qquad \qquad \qquad \qquad \qquad \qquad \qquad \qquad \nabla \times (\nabla \times \vec{E}) = \nabla(\nabla \cdot \vec{E}) - \nabla^2 \vec{E} \end{array}$$

- Electromagnetic Waves – Maxwell equations in vacuum

$$\nabla \cdot \vec{E} = 0$$

$$\nabla \cdot \vec{B} = 0$$

$$\begin{array}{l} \nabla \times \vec{E} = -\frac{1}{c} \frac{\partial \vec{B}}{\partial t} \xrightarrow{\nabla \times} \nabla \times (\nabla \times \vec{E}) = -\frac{1}{c} \frac{\partial \nabla \times \vec{B}}{\partial t} \\ \nabla \times \vec{B} = \frac{1}{c} \frac{\partial \vec{E}}{\partial t} \xrightarrow{\quad} \quad \quad \quad = -\frac{1}{c} \frac{\partial \nabla \times \vec{B}}{\partial t} = -\frac{1}{c^2} \frac{\partial^2 \vec{E}}{\partial t^2} \end{array}$$

$$\nabla \times (\nabla \times \vec{E}) = \nabla(\nabla \cdot \vec{E}) - \nabla^2 \vec{E}$$

reminder: $\nabla^2 \vec{A} = \begin{pmatrix} \nabla^2 A_x \\ \nabla^2 A_y \\ \nabla^2 A_z \end{pmatrix}$

(it is a vectorial Laplace operator...)

- Electromagnetic Waves – Maxwell equations in vacuum

$$\underline{\underline{\nabla \cdot \vec{E} = 0}}$$

$$\nabla \cdot \vec{B} = 0$$

$$\nabla \times \vec{E} = -\frac{1}{c} \frac{\partial \vec{B}}{\partial t} \xrightarrow{\nabla \times} \nabla \times (\nabla \times \vec{E}) = -\frac{1}{c} \frac{\partial \nabla \times \vec{B}}{\partial t}$$

$$\nabla \times \vec{B} = \frac{1}{c} \frac{\partial \vec{E}}{\partial t} \xrightarrow{\quad} = -\frac{1}{c} \frac{\partial \nabla \times \vec{B}}{\partial t} = -\frac{1}{c^2} \frac{\partial^2 \vec{E}}{\partial t^2}$$

$$\nabla \times (\nabla \times \vec{E}) = \underbrace{\nabla (\nabla \cdot \vec{E})}_{= 0} - \nabla^2 \vec{E}$$

- Electromagnetic Waves – Maxwell equations in vacuum

$$\underline{\underline{\nabla \cdot \vec{E} = 0}}$$

$$\nabla \cdot \vec{B} = 0$$

$$\begin{array}{l} \nabla \times \vec{E} = -\frac{1}{c} \frac{\partial \vec{B}}{\partial t} \xrightarrow{\nabla \times} \nabla \times (\nabla \times \vec{E}) = -\frac{1}{c} \frac{\partial \nabla \times \vec{B}}{\partial t} \\ \nabla \times \vec{B} = \frac{1}{c} \frac{\partial \vec{E}}{\partial t} \xrightarrow{\quad} \quad \quad \quad = -\frac{1}{c} \frac{\partial \nabla \times \vec{B}}{\partial t} = -\frac{1}{c^2} \frac{\partial^2 \vec{E}}{\partial t^2} \\ \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \downarrow \\ \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \nabla \times (\nabla \times \vec{E}) = -\nabla^2 \vec{E} \end{array}$$

- Electromagnetic Waves – Maxwell equations in vacuum

$$\underline{\underline{\nabla \cdot \vec{E} = 0}}$$

$$\nabla \cdot \vec{B} = 0$$

$$\nabla \times \vec{E} = -\frac{1}{c} \frac{\partial \vec{B}}{\partial t} \quad \xrightarrow{\nabla \times} \quad \nabla \times (\nabla \times \vec{E}) = -\frac{1}{c} \frac{\partial \nabla \times \vec{B}}{\partial t}$$

$$\nabla \times \vec{B} = \frac{1}{c} \frac{\partial \vec{E}}{\partial t} \quad \longrightarrow \quad = -\frac{1}{c} \frac{\partial \nabla \times \vec{B}}{\partial t} = \underline{\underline{-\frac{1}{c^2} \frac{\partial^2 \vec{E}}{\partial t^2}}}$$

$$\nabla \times (\nabla \times \vec{E}) = \underline{\underline{-\nabla^2 \vec{E}}}$$

- Electromagnetic Waves – Maxwell equations in vacuum

$$\nabla \cdot \vec{E} = 0$$

$$\nabla \cdot \vec{B} = 0$$

$$\nabla \times \vec{E} = -\frac{1}{c} \frac{\partial \vec{B}}{\partial t}$$

$$\nabla \times \vec{B} = \frac{1}{c} \frac{\partial \vec{E}}{\partial t}$$

$$\nabla^2 \vec{E} - \frac{1}{c^2} \frac{\partial^2 \vec{E}}{\partial t^2} = 0$$

- Electromagnetic Waves – Maxwell equations in vacuum

$$\nabla \cdot \vec{E} = 0$$

$$\nabla \cdot \vec{B} = 0$$

$$\nabla \times \vec{E} = -\frac{1}{c} \frac{\partial \vec{B}}{\partial t}$$

$$\nabla \times \vec{B} = \frac{1}{c} \frac{\partial \vec{E}}{\partial t}$$

$$\nabla^2 \vec{E} - \frac{1}{c^2} \frac{\partial^2 \vec{E}}{\partial t^2} = 0$$

analogy for \vec{B} ...

- Electromagnetic Waves – Maxwell equations in vacuum

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well known wave-equations for $\vec{E}(\vec{r}, t)$ and $\vec{B}(\vec{r}, t)$

- Electromagnetic Waves – propagation in vacuum

Ansatz:

$$\vec{E}(\vec{r}, t) = \vec{a}_1 E_0 e^{i(\vec{k} \cdot \vec{r} - \omega t)}$$

$$\vec{B}(\vec{r}, t) = \vec{a}_2 B_0 e^{i(\vec{k} \cdot \vec{r} - \omega t)}$$

$$\nabla^2 \vec{E} - \frac{1}{c^2} \frac{\partial^2 \vec{E}}{\partial t^2} = 0$$

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- Electromagnetic Waves – propagation in vacuum

Ansatz:

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$\vec{k} = k\vec{n}$, \vec{n} : unit vector in direction of wave propagation

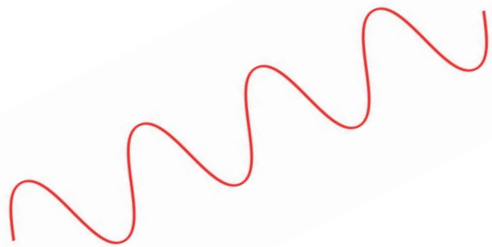
- Electromagnetic Waves – propagation in vacuum

Ansatz:

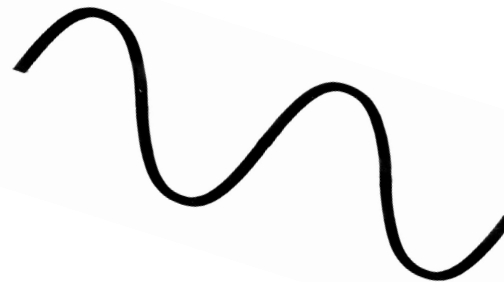
$$\vec{E}(\vec{r}, t) = \vec{a}_1 E_0 e^{i(\vec{k} \cdot \vec{r} - \omega t)}$$

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?



- Electromagnetic Waves – propagation in vacuum

Ansatz:

$$\vec{E}(\vec{r}, t) = \vec{a}_1 E_0 e^{i(\vec{k} \cdot \vec{r} - \omega t)}$$

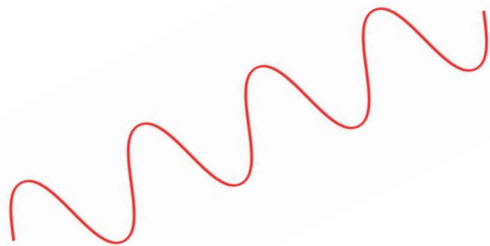
$$\vec{B}(\vec{r}, t) = \vec{a}_2 B_0 e^{i(\vec{k} \cdot \vec{r} - \omega t)}$$

\vec{a}_1 : unit vector, direction **to be determined**
 E_0 : complex constant, value **to be determined**

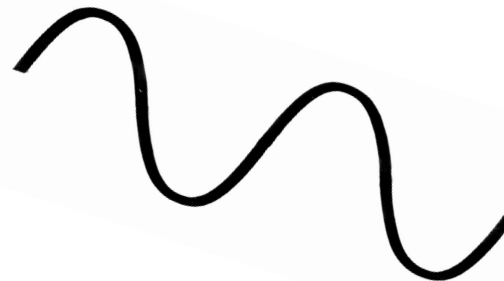
?

\vec{a}_2 : unit vector, direction **to be determined**
 B : complex constant, value **to be determined**

$\vec{k} = k\vec{n}$, \vec{n} : unit vector in direction of wave propagation



?



- Electromagnetic Waves – propagation in vacuum

Ansatz:

$$\vec{E}(\vec{r}, t) = \vec{a}_1 E_0 e^{i(\vec{k} \cdot \vec{r} - \omega t)}$$

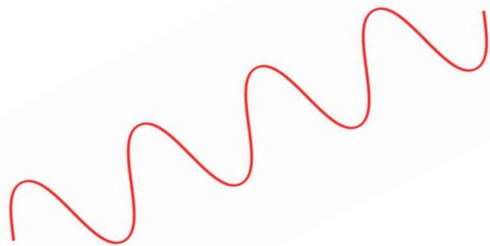
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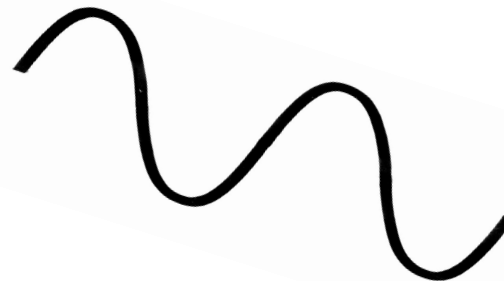
?

\vec{a}_2 : unit vector, direction **to be determined**
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$\vec{k} = k\vec{n}$, \vec{n} : unit vector in direction of wave propagation



?



insert into Maxwell equations →

- Electromagnetic Waves – propagation in vacuum

$$\nabla \cdot \vec{E} = 0$$

$$\nabla \cdot \vec{B} = 0$$

$$\nabla \times \vec{E} = -\frac{1}{c} \frac{\partial \vec{B}}{\partial t}$$

$$\nabla \times \vec{B} = \frac{1}{c} \frac{\partial \vec{E}}{\partial t}$$

- Electromagnetic Waves – propagation in vacuum

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- Electromagnetic Waves – propagation in vacuum

$$\nabla \cdot \vec{E} = 0$$

$$i\vec{k} \cdot \vec{a}_1 E_0 = 0$$

$$\nabla \cdot \vec{B} = 0$$

$$i\vec{k} \cdot \vec{a}_2 B_0 = 0$$

$$\nabla \times \vec{E} = -\frac{1}{c} \frac{\partial \vec{B}}{\partial t}$$

$$\vec{E}(\vec{r}, t) = \vec{a}_1 E_0 e^{i(\vec{k} \cdot \vec{r} - \omega t)}$$

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$$i\vec{k} \times \vec{a}_1 E_0 = \frac{i\omega}{c} \vec{a}_2 B_0$$

$$\nabla \times \vec{B} = \frac{1}{c} \frac{\partial \vec{E}}{\partial t}$$

$$i\vec{k} \times \vec{a}_2 B_0 = -\frac{i\omega}{c} \vec{a}_1 E_0$$

- Electromagnetic Waves – propagation in vacuum

both \vec{a}_1 and \vec{a}_2 are orthogonal to the direction of the wave propagation \vec{k}

$$\left\{ \begin{array}{l} i\vec{k} \cdot \vec{a}_1 E_0 = 0 \\ i\vec{k} \cdot \vec{a}_2 B_0 = 0 \end{array} \right.$$
$$i\vec{k} \times \vec{a}_1 E_0 = \frac{i\omega}{c} \vec{a}_2 B_0$$
$$i\vec{k} \times \vec{a}_2 B_0 = -\frac{i\omega}{c} \vec{a}_1 E_0$$

- Electromagnetic Waves – propagation in vacuum

both \vec{a}_1 and \vec{a}_2 are orthogonal to the direction of the wave propagation \vec{k}



\vec{a}_1 and \vec{a}_2 are orthogonal to each other

$$\left\{ \begin{array}{l} i\vec{k} \cdot \vec{a}_1 E_0 = 0 \\ i\vec{k} \cdot \vec{a}_2 B_0 = 0 \end{array} \right.$$

$$\left\{ \begin{array}{l} i\vec{k} \times \vec{a}_1 E_0 = \frac{i\omega}{c} \vec{a}_2 B_0 \\ i\vec{k} \times \vec{a}_2 B_0 = -\frac{i\omega}{c} \vec{a}_1 E_0 \end{array} \right.$$

- Electromagnetic Waves – propagation in vacuum

\vec{k} , \vec{a}_1 and \vec{a}_2 form a right-handed set

$$i\vec{k} \cdot \vec{a}_1 E_0 = 0$$

$$i\vec{k} \cdot \vec{a}_2 B_0 = 0$$

$$i\vec{k} \times \vec{a}_1 E_0 = \frac{i\omega}{c} \vec{a}_2 B_0$$

$$i\vec{k} \times \vec{a}_2 B_0 = -\frac{i\omega}{c} \vec{a}_1 E_0$$

- Electromagnetic Waves – propagation in vacuum

\vec{k}, \vec{a}_1 and \vec{a}_2 form a right-handed set:

$$\begin{aligned}\vec{a}_1 &= \vec{k} \times \vec{a}_2 \\ \vec{a}_2 &= -\vec{k} \times \vec{a}_1\end{aligned}$$

$$i\vec{k} \cdot \vec{a}_1 E_0 = 0$$

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$$i\vec{k} \times \vec{a}_1 E_0 = \frac{i\omega}{c} \vec{a}_2 B_0$$

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- Electromagnetic Waves – propagation in vacuum

\vec{k} , \vec{a}_1 and \vec{a}_2 form a right-handed set:

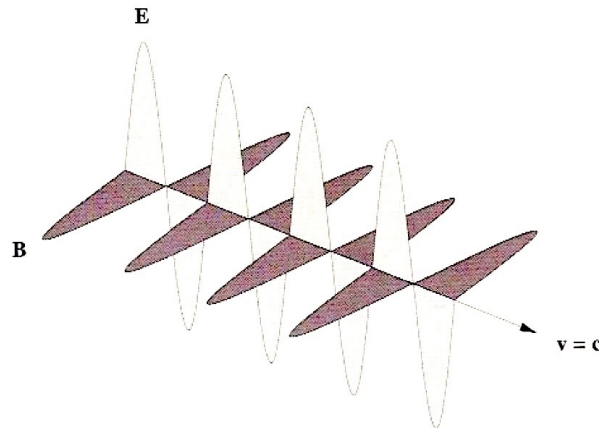
$$\begin{aligned}\vec{a}_1 &= \vec{k} \times \vec{a}_2 \\ \vec{a}_2 &= -\vec{k} \times \vec{a}_1\end{aligned}$$

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$$\vec{E}(\vec{r}, t) = \vec{a}_1 E_0 e^{i(\vec{k} \cdot \vec{r} - \omega t)}$$

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- Electromagnetic Waves – propagation in vacuum

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$$i\vec{k} \cdot \vec{a}_1 E_0 = 0$$

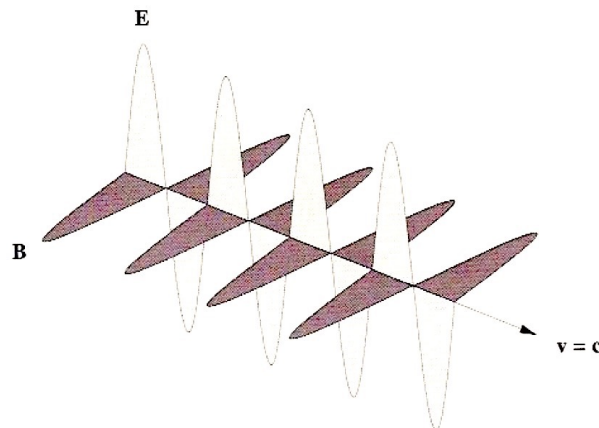
$$i\vec{k} \cdot \vec{a}_2 B_0 = 0$$

$$E_0 = \frac{\omega}{kc} B_0$$

$$B_0 = \frac{\omega}{kc} E_0$$

$$i\vec{k} \times \vec{a}_1 E_0 = \frac{i\omega}{c} \vec{a}_2 B_0$$

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- Electromagnetic Waves – propagation in vacuum

\vec{k} , \vec{a}_1 and \vec{a}_2 form a right-handed set:

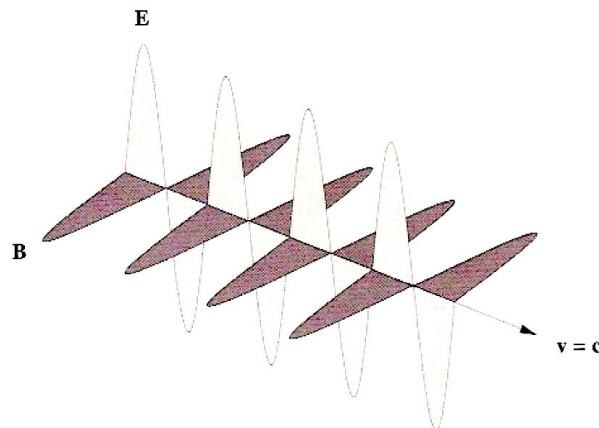
$$\begin{aligned}\vec{a}_1 &= \vec{k} \times \vec{a}_2 \\ \vec{a}_2 &= -\vec{k} \times \vec{a}_1\end{aligned}$$

$$i\vec{k} \cdot \vec{a}_1 E_0 = 0$$

$$i\vec{k} \cdot \vec{a}_2 B_0 = 0$$

$$\begin{cases} E_0 = B_0 \\ \omega = kc \end{cases} \begin{cases} E_0 = \frac{\omega}{kc} B_0 \\ B_0 = \frac{\omega}{kc} E_0 \end{cases}$$

$$\begin{cases} i\vec{k} \times \vec{a}_1 E_0 = \frac{i\omega}{c} \vec{a}_2 B_0 \\ i\vec{k} \times \vec{a}_2 B_0 = -\frac{i\omega}{c} \vec{a}_1 E_0 \end{cases}$$



$$\vec{E}(\vec{r}, t) = \vec{a}_1 E_0 e^{i(\vec{k} \cdot \vec{r} - \omega t)}$$

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- Electromagnetic Waves – propagation in vacuum

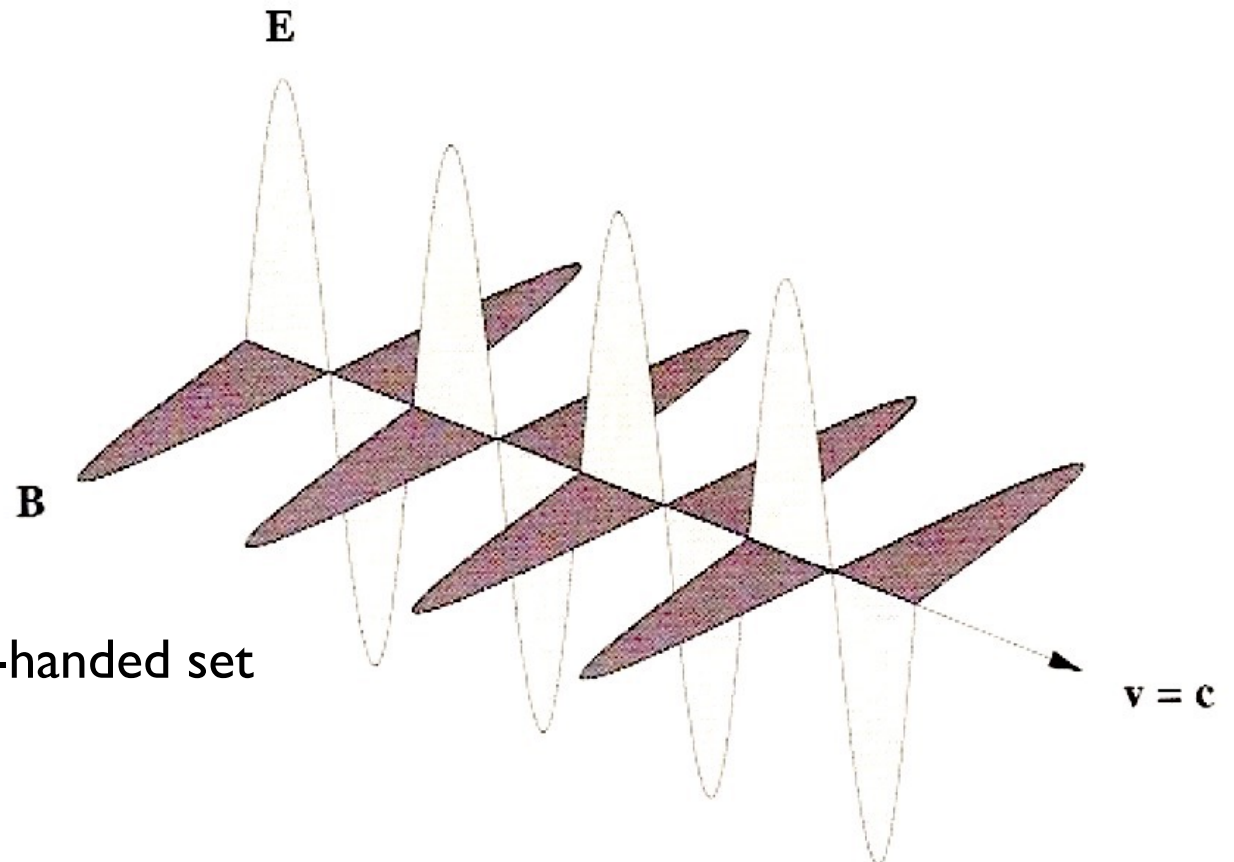
$$\vec{E}(\vec{r}, t) = \vec{a}_1 E_0 e^{i(\vec{k} \cdot \vec{r} - \omega t)}$$

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$$E_0 = B_0$$

$$\omega = kc$$

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- Electromagnetic Waves – energy flux and density

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$$\vec{S} = \frac{c}{4\pi} (\vec{E} \times \vec{B})$$

Poynting vector = directional energy flux

- Electromagnetic Waves – energy flux and density

$$\vec{E}(\vec{r}, t) = \vec{a}_1 E_0 e^{i(\vec{k} \cdot \vec{r} - \omega t)}$$

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$$\langle \vec{S} \rangle = \frac{c}{4\pi} \langle \vec{E} \times \vec{B} \rangle$$

time-averaging to eliminate ωt part

- Electromagnetic Waves – energy flux and density

$$\vec{E}(\vec{r}, t) = \vec{a}_1 E_0 e^{i(\vec{k} \cdot \vec{r} - \omega t)}$$

$$\vec{B}(\vec{r}, t) = \vec{a}_2 B_0 e^{i(\vec{k} \cdot \vec{r} - \omega t)}$$

$$\langle \vec{S} \rangle = \frac{c}{8\pi} |E_0|^2 = \frac{c}{8\pi} |B_0|^2 \quad \text{time-averaging to eliminate } \omega t \text{ part}$$

- Electromagnetic Waves – energy flux and density

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$$\langle \vec{S} \rangle = \frac{c}{8\pi} |E_0|^2 = \frac{c}{8\pi} |B_0|^2 \quad \text{time-averaging to eliminate } \omega t \text{ part}$$

$$U_{field} = \frac{1}{8\pi} (E^2 + B^2)$$

- Electromagnetic Waves – energy flux and density

$$\vec{E}(\vec{r}, t) = \vec{a}_1 E_0 e^{i(\vec{k} \cdot \vec{r} - \omega t)}$$

$$\vec{B}(\vec{r}, t) = \vec{a}_2 B_0 e^{i(\vec{k} \cdot \vec{r} - \omega t)}$$

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$$\langle U_{field} \rangle = \frac{1}{8\pi} |E_0|^2 = \frac{1}{8\pi} |B_0|^2$$

$$\rightarrow \text{velocity of energy flow: } \frac{\langle \vec{S} \rangle}{\langle U_{field} \rangle} = c$$

- Lorentz force
- Maxwell equations
- Poynting vector
- Electromagnetic waves
- **radiation spectrum**
- electromagnetic potentials

- electric and magnetic field

$$\vec{E}(\vec{r}, t) = \vec{a}_1 E_0 e^{i(\vec{k} \cdot \vec{r} - \omega t)}$$

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$$\vec{B}(\vec{r}, t) = \vec{a}_2 B_0 e^{i(\vec{k} \cdot \vec{r} - \omega t)}$$

remember: $E_0 = B_0$

- electric and magnetic field

$$\vec{E}(\vec{r}, t) = \vec{a}_1 E_0 e^{i(\vec{k} \cdot \vec{r} - \omega t)}$$

$$\vec{B}(\vec{r}, t) = \vec{a}_2 E_0 e^{i(\vec{k} \cdot \vec{r} - \omega t)}$$

- electric and magnetic field

$$\vec{E}(\vec{r}, t) = \vec{a}_1 E_0 e^{i(\vec{k} \cdot \vec{r} - \omega t)}$$

sufficient to focus on electric field alone...

$$\vec{B}(\vec{r}, t) = \vec{a}_2 E_0 e^{i(\vec{k} \cdot \vec{r} - \omega t)}$$

- electric and magnetic field

$$\vec{E}(\vec{r}, t) = \vec{a}_1 E_0 e^{i(\vec{k} \cdot \vec{r} - \omega t)}$$

- electric and magnetic field

$$\vec{E}(\vec{r}, t) = \vec{a}_1 E_0 e^{i(\vec{k} \cdot \vec{r} - \omega t)}$$

distribution of all possible ω ?

- electric and magnetic field

$$\vec{E}(\vec{r}, t) = \vec{a}_1 E_0 e^{i(\vec{k} \cdot \vec{r} - \omega t)}$$

distribution of all possible ω ?

- frequency distribution (Fourier transformation of $E(t)$)

$$\hat{E}(\omega) = \frac{1}{2\pi} \int_{-\infty}^{\infty} E(t) e^{i\omega t} dt$$

- electric and magnetic field

$$\vec{E}(\vec{r}, t) = \vec{a}_1 E_0 e^{i(\vec{k} \cdot \vec{r} - \omega t)}$$

distribution of all possible ω ?

- frequency distribution

$$\hat{E}(\omega) = \frac{1}{2\pi} \int_{-\infty}^{\infty} E(t) e^{i\omega t} dt$$

- radiation spectrum

- energy per unit area per unit time

$$\frac{dW}{dA dt} = ?$$

- electric and magnetic field

$$\vec{E}(\vec{r}, t) = \vec{a}_1 E_0 e^{i(\vec{k} \cdot \vec{r} - \omega t)}$$

distribution of all possible ω ?

- frequency distribution

$$\hat{E}(\omega) = \frac{1}{2\pi} \int_{-\infty}^{\infty} E(t) e^{i\omega t} dt$$

- radiation spectrum

- energy per unit area per unit time

$$\frac{dW}{dA dt} = \text{Poynting vector}$$

- electric and magnetic field

$$\vec{E}(\vec{r}, t) = \vec{a}_1 E_0 e^{i(\vec{k} \cdot \vec{r} - \omega t)}$$

distribution of all possible ω ?

- frequency distribution

$$\hat{E}(\omega) = \frac{1}{2\pi} \int_{-\infty}^{\infty} E(t) e^{i\omega t} dt$$

- radiation spectrum

- energy per unit area per unit time

$$\frac{dW}{dA dt} = \frac{c}{4\pi} |E(t)|^2$$

- electric and magnetic field

$$\vec{E}(\vec{r}, t) = \vec{a}_1 E_0 e^{i(\vec{k} \cdot \vec{r} - \omega t)}$$

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- energy per unit area per unit *frequency*

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$$\vec{E}(\vec{r}, t) = \vec{a}_1 E_0 e^{i(\vec{k} \cdot \vec{r} - \omega t)}$$

distribution of all possible ω ?

- frequency distribution

$$\hat{E}(\omega) = \frac{1}{2\pi} \int_{-\infty}^{\infty} E(t) e^{i\omega t} dt$$

- radiation spectrum

- energy per unit area per unit time

$$\frac{dW}{dA dt} = \frac{c}{4\pi} |E(t)|^2 \longrightarrow \frac{dW}{dA} = \frac{c}{4\pi} \int_{-\infty}^{\infty} |E(t)|^2 dt$$

- energy per unit area per unit *frequency*

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distribution of all possible ω ?

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$$\frac{dW}{dA dt} = \frac{c}{4\pi} |E(t)|^2 \longrightarrow \frac{dW}{dA} = \frac{c}{4\pi} \int_{-\infty}^{\infty} |E(t)|^2 dt$$

- energy per unit area per unit frequency

$$= \frac{c}{4\pi} 2\pi \int_{-\infty}^{\infty} |E(\omega)|^2 d\omega$$

$$\frac{dW}{dA d\omega} = ?$$

- electric and magnetic field

$$\vec{E}(\vec{r}, t) = \vec{a}_1 E_0 e^{i(\vec{k} \cdot \vec{r} - \omega t)}$$

distribution of all possible ω ?

- frequency distribution

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$$\frac{dW}{dA dt} = \frac{c}{4\pi} |E(t)|^2 \longrightarrow \frac{dW}{dA} = \frac{c}{4\pi} \int_{-\infty}^{\infty} |E(t)|^2 dt$$

- energy per unit area per unit frequency

$$\frac{dW}{dA d\omega} = ?$$

$$= \frac{c}{4\pi} 2\pi \int_{-\infty}^{\infty} |E(\omega)|^2 d\omega$$

$$= \frac{c}{4\pi} 2\pi \int_0^{\infty} |E(\omega)|^2 d\omega$$

- electric and magnetic field

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distribution of all possible ω ?

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$$\hat{E}(\omega) = \frac{1}{2\pi} \int_{-\infty}^{\infty} E(t) e^{i\omega t} dt$$

- radiation spectrum

- energy per unit area per unit time

$$\frac{dW}{dA dt} = \frac{c}{4\pi} |E(t)|^2 \longrightarrow \frac{dW}{dA} = c \int_0^{\infty} |E(\omega)|^2 d\omega$$

- energy per unit area per unit *frequency*

$$\frac{dW}{dA d\omega} = ?$$

- electric and magnetic field

$$\vec{E}(\vec{r}, t) = \vec{a}_1 E_0 e^{i(\vec{k} \cdot \vec{r} - \omega t)}$$

distribution of all possible ω ?

- frequency distribution

$$\hat{E}(\omega) = \frac{1}{2\pi} \int_{-\infty}^{\infty} E(t) e^{i\omega t} dt$$

- radiation spectrum

- energy per unit area per unit time

$$\frac{dW}{dA dt} = \frac{c}{4\pi} |E(t)|^2 \longrightarrow \frac{dW}{dA} = c \int_0^{\infty} |E(\omega)|^2 d\omega$$

- energy per unit area per unit *frequency*

$$\frac{dW}{dA d\omega} = c |E(\omega)|^2$$

- radiation spectrum – energy per unit area and unit frequency

$$\frac{dW}{dA d\omega} = c|E(\omega)|^2$$

- radiation spectrum – energy per unit area and unit frequency

$$\frac{dW}{dA d\omega} = c|E(\omega)|^2$$

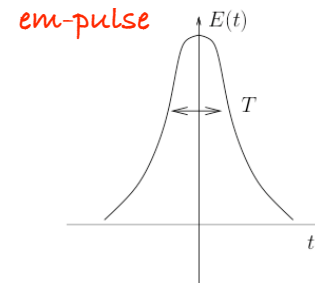
- examples

- radiation spectrum – energy per unit area and unit frequency

$$\frac{dW}{dA d\omega} = c|E(\omega)|^2$$

- examples

- finite pulse

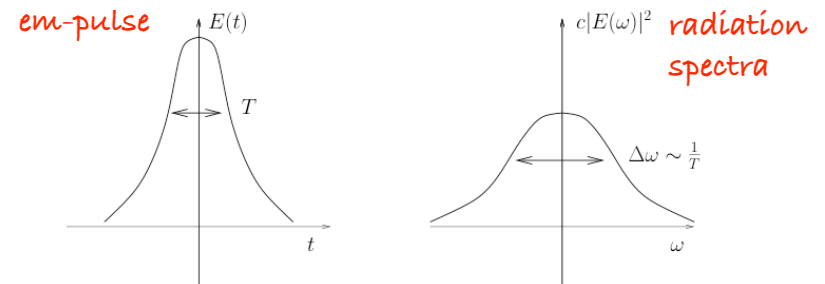


- radiation spectrum – energy per unit area and unit frequency

$$\frac{dW}{dA d\omega} = c|E(\omega)|^2$$

- examples

- finite pulse



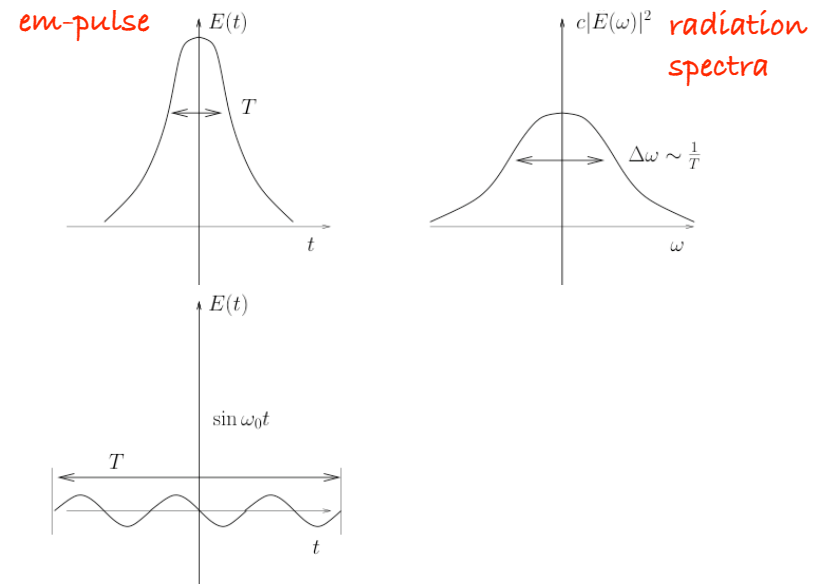
- radiation spectrum – energy per unit area and unit frequency

$$\frac{dW}{dA d\omega} = c|E(\omega)|^2$$

- examples

- finite pulse

- periodic signal of finite time



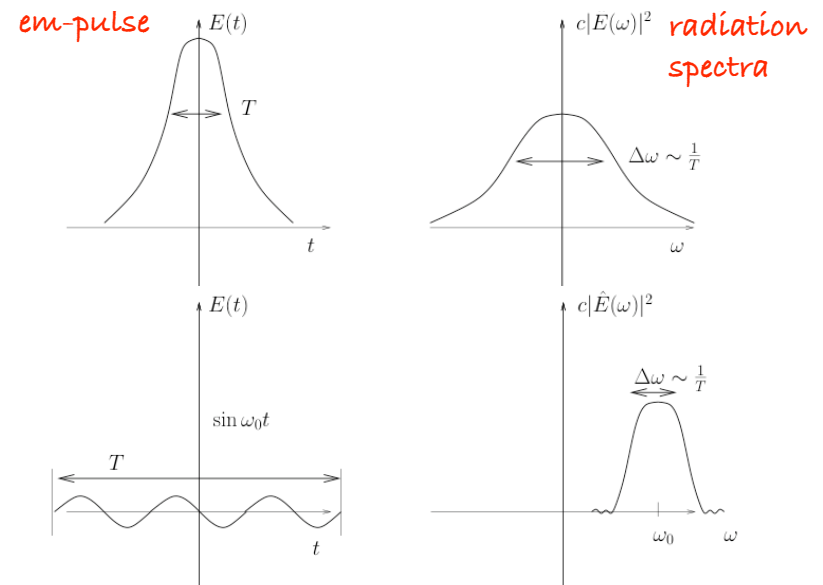
- radiation spectrum – energy per unit area and unit frequency

$$\frac{dW}{dA d\omega} = c|E(\omega)|^2$$

- examples

- finite pulse

- periodic signal of finite time



- radiation spectrum – energy per unit area and unit frequency

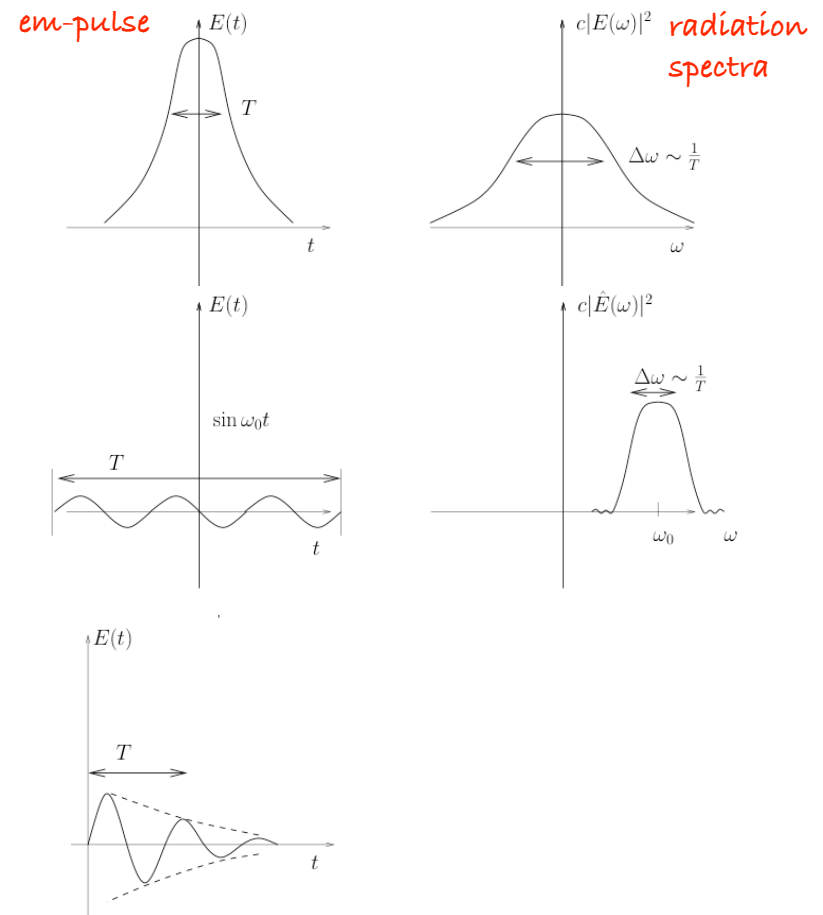
$$\frac{dW}{dA d\omega} = c|E(\omega)|^2$$

- examples

- finite pulse

- periodic signal of finite time

- decaying periodic signal of finite time

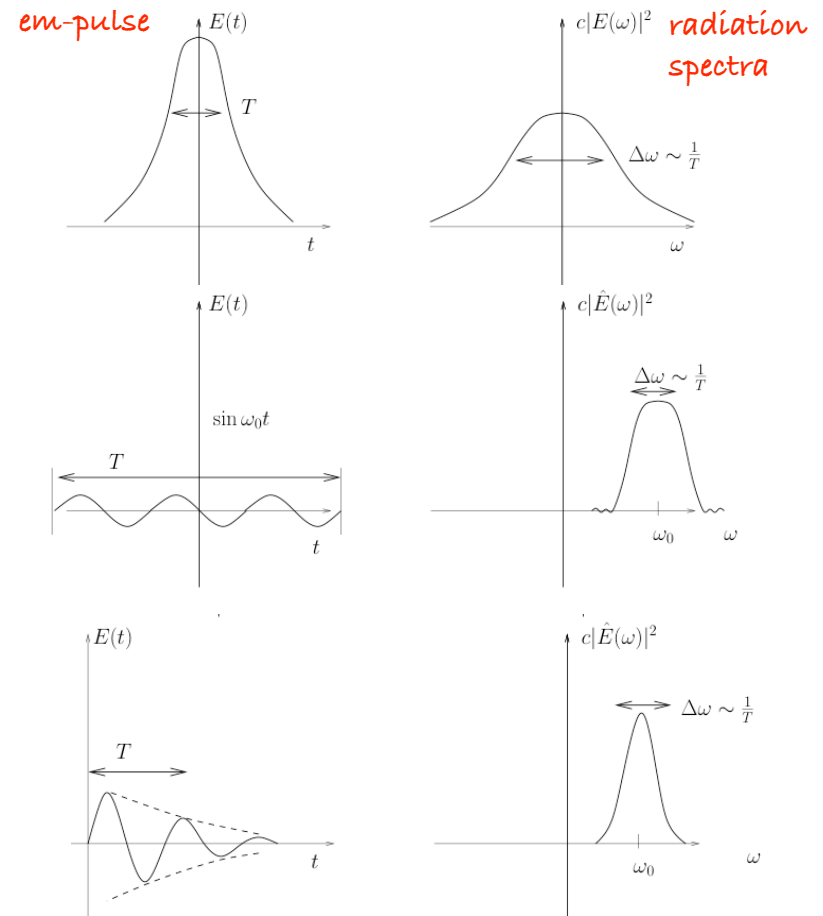


- radiation spectrum – energy per unit area and unit frequency

$$\frac{dW}{dA d\omega} = c|E(\omega)|^2$$

examples

- finite pulse
- periodic signal of finite time
- decaying periodic signal of finite time



- Lorentz force
- Maxwell equations
- Poynting vector
- electromagnetic waves
- radiation spectrum
- **electromagnetic potentials**

- Electromagnetic Potentials

- Maxwell equations

$$\nabla \cdot \vec{E} = 4\pi\rho$$

$$\nabla \cdot \vec{B} = 0$$

$$\nabla \times \vec{E} = -\frac{1}{c} \frac{\partial \vec{B}}{\partial t}$$

$$\nabla \times \vec{B} = \frac{4\pi}{c} \vec{J} + \frac{1}{c} \frac{\partial \vec{E}}{\partial t}$$

- Electromagnetic Potentials

- Maxwell equations

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- vector fields

$$\vec{E}(\vec{r}, t) \quad \vec{B}(\vec{r}, t)$$

▪ Electromagnetic Potentials

- Maxwell equations

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- potential \vec{A}

—————→ a divergence-free field can be written as a rotation of a vector field

- vector fields

$$\vec{E}(\vec{r}, t) \quad \vec{B}(\vec{r}, t)$$

▪ Electromagnetic Potentials

- Maxwell equations

$$\nabla \cdot \vec{E} = 4\pi\rho$$

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- potential \vec{A}

$$\vec{B} = \nabla \times \vec{A}$$

- vector fields

$$\vec{E}(\vec{r}, t) \quad \vec{B}(\vec{r}, t)$$

- Electromagnetic Potentials

- Maxwell equations

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$$\vec{B} = \nabla \times \vec{A}$$

$$\nabla \times \vec{E} = -\frac{1}{c} \frac{\partial \nabla \times \vec{A}}{\partial t} = -\nabla \times \frac{1}{c} \frac{\partial \vec{A}}{\partial t}$$

- vector fields

$$\vec{E}(\vec{r}, t) \quad \vec{B}(\vec{r}, t)$$

- Electromagnetic Potentials

- Maxwell equations

$$\nabla \cdot \vec{E} = 4\pi\rho$$

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$$\nabla \times \vec{B} = \frac{4\pi}{c} \vec{J} + \frac{1}{c} \frac{\partial \vec{E}}{\partial t}$$

- potential \vec{A}

$$\vec{B} = \nabla \times \vec{A}$$

$$\underline{\underline{\nabla \times \vec{E}}} = -\frac{1}{c} \frac{\partial \nabla \times \vec{A}}{\partial t} = \underline{\underline{-\nabla \times \frac{1}{c} \frac{\partial \vec{A}}{\partial t}}}$$

- vector fields

$$\vec{E}(\vec{r}, t) \quad \vec{B}(\vec{r}, t)$$

- Electromagnetic Potentials

- Maxwell equations

$$\nabla \cdot \vec{E} = 4\pi\rho$$

$$\nabla \cdot \vec{B} = 0$$

$$\nabla \times \vec{E} = -\frac{1}{c} \frac{\partial \vec{B}}{\partial t}$$

$$\nabla \times \vec{B} = \frac{4\pi}{c} \vec{J} + \frac{1}{c} \frac{\partial \vec{E}}{\partial t}$$

- potential \vec{A}

$$\vec{B} = \nabla \times \vec{A}$$

$$0 = \nabla \times \left(\vec{E} + \frac{1}{c} \frac{\partial \vec{A}}{\partial t} \right)$$

- vector fields

$$\vec{E}(\vec{r}, t) \quad \vec{B}(\vec{r}, t)$$

- Electromagnetic Potentials

- Maxwell equations

$$\nabla \cdot \vec{E} = 4\pi\rho$$

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$$\nabla \times \vec{E} = -\frac{1}{c} \frac{\partial \vec{B}}{\partial t}$$

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- vector fields

$$\vec{E}(\vec{r}, t) \quad \vec{B}(\vec{r}, t)$$

a curl-free field can be written as a gradient of a potential

- Electromagnetic Potentials

- Maxwell equations

$$\nabla \cdot \vec{E} = 4\pi\rho$$

$$\nabla \cdot \vec{B} = 0$$

$$\nabla \times \vec{E} = -\frac{1}{c} \frac{\partial \vec{B}}{\partial t}$$

$$\nabla \times \vec{B} = \frac{4\pi}{c} \vec{J} + \frac{1}{c} \frac{\partial \vec{E}}{\partial t}$$

- vector fields

$$\vec{E}(\vec{r}, t) \quad \vec{B}(\vec{r}, t)$$

- potential \vec{A}

$$\vec{B} = \nabla \times \vec{A}$$

$$0 = \nabla \times \left(\vec{E} + \frac{1}{c} \frac{\partial \vec{A}}{\partial t} \right)$$

$$\left(\vec{E} + \frac{1}{c} \frac{\partial \vec{A}}{\partial t} \right) = -\nabla\phi$$

a curl-free field can be written as a gradient of a potential

▪ Electromagnetic Potentials

- Maxwell equations

$$\nabla \cdot \vec{E} = 4\pi\rho$$

$$\nabla \cdot \vec{B} = 0$$

$$\nabla \times \vec{E} = -\frac{1}{c} \frac{\partial \vec{B}}{\partial t}$$

$$\nabla \times \vec{B} = \frac{4\pi}{c} \vec{J} + \frac{1}{c} \frac{\partial \vec{E}}{\partial t}$$

- potentials
- \vec{A}
- and
- ϕ

$$\vec{B} = \nabla \times \vec{A}$$

$$\vec{E} = -\nabla\phi - \frac{1}{c} \frac{\partial \vec{A}}{\partial t}$$

▪ Electromagnetic Potentials

- Maxwell equations

$$\nabla \cdot \vec{E} = 4\pi\rho$$

$$\nabla \cdot \vec{B} = 0$$

satisfied by construction

$$\nabla \times \vec{E} = -\frac{1}{c} \frac{\partial \vec{B}}{\partial t}$$

$$\nabla \times \vec{B} = \frac{4\pi}{c} \vec{J} + \frac{1}{c} \frac{\partial \vec{E}}{\partial t}$$

- potentials \vec{A} and ϕ

$$\vec{B} = \nabla \times \vec{A}$$

$$\vec{E} = -\nabla\phi - \frac{1}{c} \frac{\partial \vec{A}}{\partial t}$$

- Electromagnetic Potentials

- Maxwell equations

$$\nabla \cdot \vec{E} = 4\pi\rho$$

$$\nabla \times \vec{B} = \frac{4\pi}{c} \vec{J} + \frac{1}{c} \frac{\partial \vec{E}}{\partial t}$$

- potentials \vec{A} and ϕ

$$\vec{B} = \nabla \times \vec{A}$$

$$\vec{E} = -\nabla\phi - \frac{1}{c} \frac{\partial \vec{A}}{\partial t}$$

- Electromagnetic Potentials

- Maxwell equations

$$\nabla \cdot \vec{E} = 4\pi\rho$$

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- potentials \vec{A} and ϕ

$$\vec{B} = \nabla \times \vec{A}$$
$$\vec{E} = -\nabla\phi - \frac{1}{c} \frac{\partial \vec{A}}{\partial t}$$

turn equation around!?

- Electromagnetic Potentials

- Maxwell equations

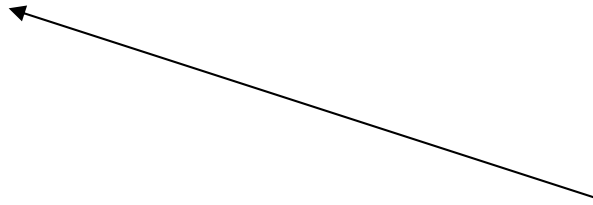
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$$\nabla \times \vec{B} = \frac{4\pi}{c} \vec{J} + \frac{1}{c} \frac{\partial \vec{E}}{\partial t}$$

- potentials \vec{A} and ϕ

$$\vec{B} = \nabla \times \vec{A}$$

$$\vec{E} = -\nabla\phi - \frac{1}{c} \frac{\partial \vec{A}}{\partial t}$$



- Electromagnetic Potentials

- Maxwell equations

$$\nabla \cdot \vec{E} = 4\pi\rho$$

$$\nabla \times \vec{B} = \frac{4\pi}{c} \vec{J} + \frac{1}{c} \frac{\partial \vec{E}}{\partial t}$$

- potentials \vec{A} and ϕ

$$\vec{B} = \nabla \times \vec{A}$$

$$\vec{E} = -\nabla\phi - \frac{1}{c} \frac{\partial \vec{A}}{\partial t}$$

$$4\pi\rho = \nabla \cdot \vec{E} = -\nabla \cdot \nabla\phi - \frac{1}{c} \frac{\partial \nabla \cdot \vec{A}}{\partial t} = -\nabla \cdot \nabla\phi - \frac{1}{c} \frac{\partial}{\partial t} \left(\nabla \cdot \vec{A} + \frac{1}{c} \frac{\partial\phi}{\partial t} - \frac{1}{c} \frac{\partial\phi}{\partial t} \right) = -\nabla \cdot \nabla\phi - \frac{1}{c^2} \frac{\partial^2\phi}{\partial t^2} - \frac{1}{c} \frac{\partial}{\partial t} \left(\nabla \cdot \vec{A} + \frac{1}{c} \frac{\partial\phi}{\partial t} \right)$$

- Electromagnetic Potentials

- Maxwell equations

$$\nabla \cdot \vec{E} = 4\pi\rho$$

$$\nabla \times \vec{B} = \frac{4\pi}{c} \vec{J} + \frac{1}{c} \frac{\partial \vec{E}}{\partial t}$$

- potentials \vec{A} and ϕ

$$\vec{B} = \nabla \times \vec{A}$$

$$\vec{E} = -\nabla\phi - \frac{1}{c} \frac{\partial \vec{A}}{\partial t}$$

$$4\pi\rho = \nabla \cdot \vec{E} = -\nabla \cdot \nabla\phi - \frac{1}{c} \frac{\partial \nabla \cdot \vec{A}}{\partial t} = -\nabla \cdot \nabla\phi - \frac{1}{c} \frac{\partial}{\partial t} \left(\nabla \cdot \vec{A} + \frac{1}{c} \frac{\partial\phi}{\partial t} - \frac{1}{c} \frac{\partial\phi}{\partial t} \right) = -\nabla \cdot \nabla\phi - \frac{1}{c^2} \frac{\partial^2\phi}{\partial t^2} - \frac{1}{c} \frac{\partial}{\partial t} \left(\nabla \cdot \vec{A} + \frac{1}{c} \frac{\partial\phi}{\partial t} \right)$$

$$\longrightarrow \nabla^2\phi - \frac{1}{c^2} \frac{\partial^2\phi}{\partial t^2} = 4\pi\rho - \frac{1}{c} \frac{\partial}{\partial t} \left(\nabla \cdot \vec{A} + \frac{1}{c} \frac{\partial\phi}{\partial t} \right)$$

- Electromagnetic Potentials

- Maxwell equations

$$\nabla \cdot \vec{E} = 4\pi\rho$$

$$\nabla \times \vec{B} = \frac{4\pi}{c} \vec{J} + \frac{1}{c} \frac{\partial \vec{E}}{\partial t}$$

- potentials \vec{A} and ϕ

$$\vec{B} = \nabla \times \vec{A}$$

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$$4\pi\rho = \nabla \cdot \vec{E} = -\nabla \cdot \nabla\phi - \frac{1}{c} \frac{\partial \nabla \cdot \vec{A}}{\partial t} = -\nabla \cdot \nabla\phi - \frac{1}{c} \frac{\partial}{\partial t} \left(\nabla \cdot \vec{A} + \frac{1}{c} \frac{\partial\phi}{\partial t} - \frac{1}{c} \frac{\partial\phi}{\partial t} \right) = -\nabla \cdot \nabla\phi - \frac{1}{c^2} \frac{\partial^2\phi}{\partial t^2} - \frac{1}{c} \frac{\partial}{\partial t} \left(\nabla \cdot \vec{A} + \frac{1}{c} \frac{\partial\phi}{\partial t} \right)$$

$$\longrightarrow \nabla^2\phi - \frac{1}{c^2} \frac{\partial^2\phi}{\partial t^2} = 4\pi\rho - \frac{1}{c} \frac{\partial}{\partial t} \left(\nabla \cdot \vec{A} + \frac{1}{c} \frac{\partial\phi}{\partial t} \right) \quad \text{inhomogeneous wave-equation}$$

▪ Electromagnetic Potentials

- Maxwell equations

$$\nabla \times \vec{B} = \frac{4\pi}{c} \vec{J} + \frac{1}{c} \frac{\partial \vec{E}}{\partial t}$$

- potentials \vec{A} and ϕ

$$\vec{B} = \nabla \times \vec{A}$$

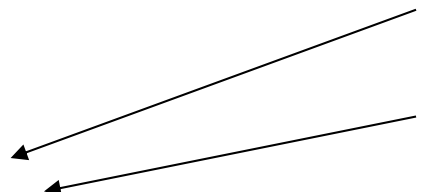
$$\vec{E} = -\nabla\phi - \frac{1}{c} \frac{\partial \vec{A}}{\partial t}$$

$$\nabla^2 \phi - \frac{1}{c^2} \frac{\partial^2 \phi}{\partial t^2} = 4\pi\rho - \frac{1}{c} \frac{\partial}{\partial t} \left(\nabla \cdot \vec{A} + \frac{1}{c} \frac{\partial \phi}{\partial t} \right)$$

▪ Electromagnetic Potentials

• Maxwell equations

• potentials \vec{A} and ϕ

$$\nabla \times \vec{B} = \frac{4\pi}{c} \vec{J} + \frac{1}{c} \frac{\partial \vec{E}}{\partial t}$$
$$\vec{B} = \nabla \times \vec{A}$$
$$\vec{E} = -\nabla \phi - \frac{1}{c} \frac{\partial \vec{A}}{\partial t}$$
$$\nabla^2 \phi - \frac{1}{c^2} \frac{\partial^2 \phi}{\partial t^2} = 4\pi \rho - \frac{1}{c} \frac{\partial}{\partial t} \left(\nabla \cdot \vec{A} + \frac{1}{c} \frac{\partial \phi}{\partial t} \right)$$


- Electromagnetic Potentials

- Maxwell equations

- potentials \vec{A} and ϕ

$$\nabla \times \vec{B} = \frac{4\pi}{c} \vec{j} + \frac{1}{c} \frac{\partial \vec{E}}{\partial t}$$

$$\vec{B} = \nabla \times \vec{A}$$

$$\vec{E} = -\nabla \phi - \frac{1}{c} \frac{\partial \vec{A}}{\partial t}$$

$$\nabla^2 \phi - \frac{1}{c^2} \frac{\partial^2 \phi}{\partial t^2} = 4\pi \rho - \frac{1}{c} \frac{\partial}{\partial t} \left(\nabla \cdot \vec{A} + \frac{1}{c} \frac{\partial \phi}{\partial t} \right)$$

$$\nabla \times \vec{B} = \nabla \times \nabla \times \vec{A} = \frac{4\pi}{c} \vec{j} + \frac{1}{c} \frac{\partial}{\partial t} \left(-\nabla \phi - \frac{1}{c} \frac{\partial \vec{A}}{\partial t} \right)$$

- Electromagnetic Potentials

- Maxwell equations

- potentials \vec{A} and ϕ

$$\nabla \times \vec{B} = \frac{4\pi}{c} \vec{j} + \frac{1}{c} \frac{\partial \vec{E}}{\partial t}$$

$$\vec{B} = \nabla \times \vec{A}$$

$$\vec{E} = -\nabla \phi - \frac{1}{c} \frac{\partial \vec{A}}{\partial t}$$

$$\nabla^2 \phi - \frac{1}{c^2} \frac{\partial^2 \phi}{\partial t^2} = 4\pi \rho - \frac{1}{c} \frac{\partial}{\partial t} \left(\nabla \cdot \vec{A} + \frac{1}{c} \frac{\partial \phi}{\partial t} \right)$$

$$\nabla \times \vec{B} = \nabla \times \nabla \times \vec{A} = \frac{4\pi}{c} \vec{j} + \frac{1}{c} \frac{\partial}{\partial t} \left(-\nabla \phi - \frac{1}{c} \frac{\partial \vec{A}}{\partial t} \right)$$

$$\xrightarrow{\nabla \times \nabla \times \vec{A} = -\nabla^2 \vec{A} + \nabla(\nabla \cdot \vec{A})} \nabla^2 \vec{A} - \frac{1}{c^2} \frac{\partial^2 \vec{A}}{\partial t^2} = -\frac{4\pi}{c} \vec{j} + \frac{1}{c} \frac{\partial}{\partial t} \left(\nabla \cdot \vec{A} + \frac{1}{c} \frac{\partial \phi}{\partial t} \right)$$

- Electromagnetic Potentials

- Maxwell equations

- potentials \vec{A} and ϕ

$$\nabla \times \vec{B} = \frac{4\pi}{c} \vec{j} + \frac{1}{c} \frac{\partial \vec{E}}{\partial t}$$

$$\vec{B} = \nabla \times \vec{A}$$

$$\vec{E} = -\nabla \phi - \frac{1}{c} \frac{\partial \vec{A}}{\partial t}$$

$$\nabla^2 \phi - \frac{1}{c^2} \frac{\partial^2 \phi}{\partial t^2} = 4\pi \rho - \frac{1}{c} \frac{\partial}{\partial t} \left(\nabla \cdot \vec{A} + \frac{1}{c} \frac{\partial \phi}{\partial t} \right)$$

$$\nabla \times \vec{B} = \nabla \times \nabla \times \vec{A} = \frac{4\pi}{c} \vec{j} + \frac{1}{c} \frac{\partial}{\partial t} \left(-\nabla \phi - \frac{1}{c} \frac{\partial \vec{A}}{\partial t} \right)$$

$$\xrightarrow{\nabla \times \nabla \times \vec{A} = -\nabla^2 \vec{A} + \nabla(\nabla \cdot \vec{A})}$$

$$\nabla^2 \vec{A} - \frac{1}{c^2} \frac{\partial^2 \vec{A}}{\partial t^2} = -\frac{4\pi}{c} \vec{j} + \frac{1}{c} \frac{\partial}{\partial t} \left(\nabla \cdot \vec{A} + \frac{1}{c} \frac{\partial \phi}{\partial t} \right)$$

inhomogeneous wave-equation

- Electromagnetic Potentials

- electromagnetic fields \vec{E} and \vec{B} equations

$$\vec{B} = \nabla \times \vec{A}$$

$$\vec{E} = -\nabla\phi - \frac{1}{c} \frac{\partial \vec{A}}{\partial t}$$

- expressed via potentials \vec{A} and ϕ

$$\nabla^2 \phi - \frac{1}{c^2} \frac{\partial^2 \phi}{\partial t^2} = 4\pi\rho - \frac{1}{c} \frac{\partial}{\partial t} \left(\nabla \cdot \vec{A} + \frac{1}{c} \frac{\partial \phi}{\partial t} \right)$$

$$\nabla^2 \vec{A} - \frac{1}{c^2} \frac{\partial^2 \vec{A}}{\partial t^2} = -\frac{4\pi}{c} \vec{j} + \frac{1}{c} \frac{\partial}{\partial t} \left(\nabla \cdot \vec{A} + \frac{1}{c} \frac{\partial \phi}{\partial t} \right)$$

- Electromagnetic Potentials

- electromagnetic fields \vec{E} and \vec{B} equations

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$$\nabla^2 \phi - \frac{1}{c^2} \frac{\partial^2 \phi}{\partial t^2} = 4\pi\rho - \frac{1}{c} \frac{\partial}{\partial t} \left(\nabla \cdot \vec{A} + \frac{1}{c} \frac{\partial \phi}{\partial t} \right)$$

$$\nabla^2 \vec{A} - \frac{1}{c^2} \frac{\partial^2 \vec{A}}{\partial t^2} = -\frac{4\pi}{c} \vec{j} + \frac{1}{c} \frac{\partial}{\partial t} \left(\nabla \cdot \vec{A} + \frac{1}{c} \frac{\partial \phi}{\partial t} \right)$$

how to solve them?

- Electromagnetic Potentials

- electromagnetic fields \vec{E} and \vec{B} equations

$$\vec{B} = \nabla \times \vec{A}$$

$$\vec{E} = -\nabla\phi - \frac{1}{c} \frac{\partial \vec{A}}{\partial t}$$

- expressed via potentials \vec{A} and ϕ

$$\nabla^2 \phi - \frac{1}{c^2} \frac{\partial^2 \phi}{\partial t^2} = 4\pi\rho - \frac{1}{c} \frac{\partial}{\partial t} \left(\nabla \cdot \vec{A} + \frac{1}{c} \frac{\partial \phi}{\partial t} \right)$$

$$\nabla^2 \vec{A} - \frac{1}{c^2} \frac{\partial^2 \vec{A}}{\partial t^2} = -\frac{4\pi}{c} \vec{j} + \frac{1}{c} \frac{\partial}{\partial t} \left(\nabla \cdot \vec{A} + \frac{1}{c} \frac{\partial \phi}{\partial t} \right)$$

how to solve them?

\vec{A} and ϕ are not uniquely determined!

- Electromagnetic Potentials

- electromagnetic fields \vec{E} and \vec{B} equations

$$\vec{B} = \nabla \times \vec{A} \qquad \text{Gauge transformation} \qquad = \nabla \times (\vec{A} + \nabla \psi)$$

$$\vec{E} = -\nabla \phi - \frac{1}{c} \frac{\partial \vec{A}}{\partial t}$$

- expressed via potentials \vec{A} and ϕ

$$\nabla^2 \phi - \frac{1}{c^2} \frac{\partial^2 \phi}{\partial t^2} = 4\pi\rho - \frac{1}{c} \frac{\partial}{\partial t} \left(\nabla \cdot \vec{A} + \frac{1}{c} \frac{\partial \phi}{\partial t} \right)$$

how to solve them?

$$\nabla^2 \vec{A} - \frac{1}{c^2} \frac{\partial^2 \vec{A}}{\partial t^2} = -\frac{4\pi}{c} \vec{j} + \frac{1}{c} \frac{\partial}{\partial t} \left(\nabla \cdot \vec{A} + \frac{1}{c} \frac{\partial \phi}{\partial t} \right)$$

\vec{A} and ϕ are not uniquely determined!

- Electromagnetic Potentials

- electromagnetic fields \vec{E} and \vec{B} equations

$$\begin{aligned}\vec{B} &= \nabla \times \vec{A} && \stackrel{\text{Gauge transformation}}{=} \nabla \times (\vec{A} + \nabla \psi) \\ \vec{E} &= -\nabla \phi - \frac{1}{c} \frac{\partial \vec{A}}{\partial t} && = -\nabla \phi - \frac{1}{c} \frac{\partial (\vec{A} + \nabla \psi)}{\partial t}\end{aligned}$$

- expressed via potentials \vec{A} and ϕ

$$\nabla^2 \phi - \frac{1}{c^2} \frac{\partial^2 \phi}{\partial t^2} = 4\pi\rho - \frac{1}{c} \frac{\partial}{\partial t} \left(\nabla \cdot \vec{A} + \frac{1}{c} \frac{\partial \phi}{\partial t} \right)$$

how to solve them?

$$\nabla^2 \vec{A} - \frac{1}{c^2} \frac{\partial^2 \vec{A}}{\partial t^2} = -\frac{4\pi}{c} \vec{j} + \frac{1}{c} \frac{\partial}{\partial t} \left(\nabla \cdot \vec{A} + \frac{1}{c} \frac{\partial \phi}{\partial t} \right)$$

\vec{A} and ϕ are not uniquely determined!

- Electromagnetic Potentials

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$$\begin{aligned}\vec{B} &= \nabla \times \vec{A} && \stackrel{\text{Gauge transformation}}{=} \nabla \times (\vec{A} + \nabla \psi) \\ \vec{E} &= -\nabla \phi - \frac{1}{c} \frac{\partial \vec{A}}{\partial t} && = -\nabla \left(\phi - \frac{1}{c} \frac{\partial \psi}{\partial t} \right) - \frac{1}{c} \frac{\partial \vec{A}}{\partial t}\end{aligned}$$

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can be used as constraint...

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$$\phi(\vec{r}, t) = \int \frac{\rho(\vec{r}', t - \frac{1}{c} |\vec{r} - \vec{r}'|)}{|\vec{r} - \vec{r}'|} d^3\vec{r}'$$

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- retarded time

$$t' = t - \frac{1}{c} |\vec{r} - \vec{r}'|$$

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→ time duration taken for a signal to reach its destination
(propagation delay)

- Lorentz force
- Maxwell equations
- Poynting vector
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summary!?

- **summary**

a) we have known charge and current distributions

$$\rho(\vec{r}, t)$$

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c) determine E and B

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